

A conservative well-balanced hybrid SPH scheme for the shallow-water model

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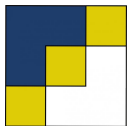
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HYDROCEAN

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History of the SPH method

- 1977: Monaghan, Gingold, Lucy - particle method for astrophysics, coined the term SPH: “smoothed particle hydrodynamics”
- 1994: Monaghan - SPH for free-surface hydrodynamics
- 1998: Vila - SPH formulation using Riemann problems
- recent & ongoing work:
 - multi-fluid SPH
 - variable mesh
 - viscous terms

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The regularizing kernel

General kernel expression

$$W(r, h) = \frac{C_\theta}{h} \theta \left(\frac{|r|}{h} \right), \text{ with}$$

- θ cut-off function
- C_θ normalization constant

properties of this kernel:

- 1 bell-shaped even function
of class C^∞
- 2 compact support K
- 3 bell parameters:
 r (position) and h (width)

$$4 \int_{\mathbb{R}} W(r, h) dr = 1$$

$$5 \int_{\mathbb{R}} W'(r, h) dr = 0$$

The particle approximation

- $f(x) = (f * \delta)(x)$, with $f : \mathbb{R} \mapsto \mathbb{R}$ and δ the Dirac distribution

$$= \int_{\mathbb{R}} f(y) \delta(x - y) dy$$

- $\Pi^h(f)(x) = (f * W)(x)$

$$= \int_K f(y) W(x - y, h) dy \simeq f(x)$$

- $\Pi^h(f')(x) = \int_K f'(y) W(x - y, h) dy$

$$= [f(y) W(x - y, h)]_{\partial K} - \int_K f(y) (W(x - y, h))' dy$$

$$= \int_K f(y) W'(x - y, h) dy \simeq f'(x)$$

The particle approximation

Accuracy of the continuous approximation

second-order accuracy requires properties 4 and 5
(Mas-Gallic - Raviart, 1987; Monaghan, 1992):

- $\int_{\mathbb{R}} W(r, h) dr = 1$, i.e. $\Pi^h(1) = 1$
- $\int_{\mathbb{R}} W'(r, h) dr = 0$, i.e. $\Pi^h(1') = 0$

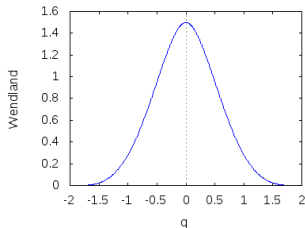
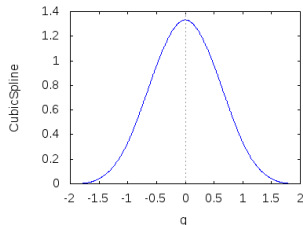
Different kernels

cut-off for the cubic spline
kernel (Monaghan, 1998):

$$\theta(q) = \begin{cases} 4 - 6q^2 + 3q^3 & \text{if } 0 \leq q < 1 \\ (2 - q)^3 & \text{if } 1 \leq q < 2 \\ 0 & \text{otherwise} \end{cases}$$

cut-off for the Wendland kernel
(Wendland, 1995):

$$\theta(q) = \begin{cases} (2 - q)^4(1 + 2q) & \text{if } 0 \leq q < 2 \\ 0 & \text{otherwise} \end{cases}$$



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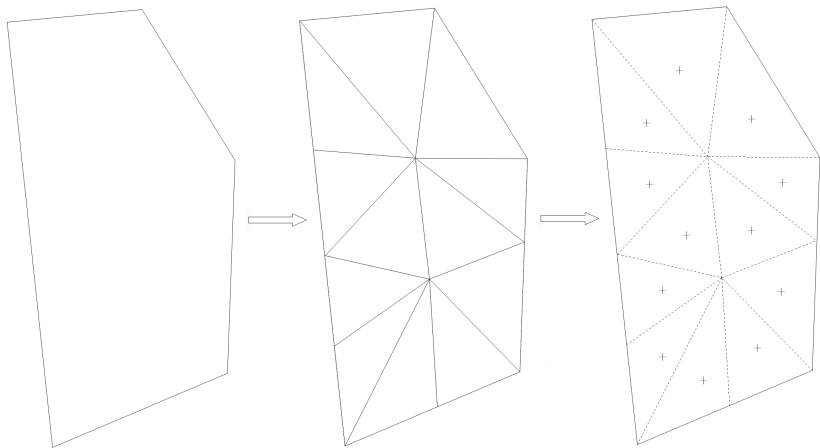
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The “mesh”



Discrete SPH equations

A quadrature formula

$$\int_{\mathbb{R}} f(y) dy \simeq \sum_{j \in \mathbb{Z}} \omega(x_j) f(x_j) = \sum_{j \in \mathbb{Z}} \omega_j f_j, \text{ where:}$$

- x_j are the quadrature points, or *particles*
 - $\omega_j = \omega(x_j)$ are their volumes
 - f_j denotes $f(x_j)$
-
- $W_{ij} = W(x_i - x_j, h)$
 - \mathcal{P} : set of interacting particles x_j close enough to particle x_i

Discrete SPH equations

Approximation of a function

$$\begin{aligned}\Pi^h(f)(x) &= \int_K f(y)W(x-y, h) dy \\ \text{becomes } \Pi^h(f)_i &= \sum_{j \in \mathcal{P}} \omega_j f_j W_{ij} \simeq f_i\end{aligned}$$

Approximation of its derivative

$$\begin{aligned}\Pi^h(f')(x) &= \int_K f(y)W'(x-y, h) dy \\ \text{becomes } \Pi^h(f')_i &= \sum_{j \in \mathcal{P}} \omega_j f_j W'_{ij} \simeq f'_i\end{aligned}$$

Main issue: consistency

Properties not verified in discrete form!

the discrete analogues of 4 and 5 are generally not true:

$$\sum_{j \in \mathcal{P}} \omega_j W_{ij} \neq 1 \quad \text{and} \quad \sum_{j \in \mathcal{P}} \omega_j W'_{ij} \neq 0$$

↪ loss of the consistency

aim of the SPH methods: numerical resolution of PDE's

→ we need a suitable derivation operator

Main issue: consistency

Weak formulation

reinforce the derivation operator:

$$\begin{aligned} D_h(f)_i &= \Pi^h(f')_i - f_i \Pi^h(1')_i \\ &= \sum_{j \in \mathcal{P}} \omega_j (f_j - f_i) W'_{ij} \simeq f'_i \end{aligned}$$

$\rightsquigarrow D_h(f)_i$ is exactly 0 for constant f

yet another issue: this formulation is *not conservative!*

Main issue: consistency

Conservativity

the formulation $D_h(f)_i$ will be *conservative* iff

$$\sum_{i \in \mathbb{Z}} \omega_i D_h(f)_i = 0$$

$$\begin{aligned} \sum_{i \in \mathbb{Z}} \omega_i D_h(f)_i &= \sum_{i \in \mathbb{Z}} \sum_{j \in \mathbb{Z}} \omega_i \omega_j (f_j - f_i) W'_{ij}, \text{ with } W' \text{ odd} \\ &= -2 \sum_{i \in \mathbb{Z}} \left(\omega_i f_i \sum_{j \in \mathbb{Z}} \omega_j W'_{ij} \right) \neq 0 \end{aligned}$$

Main issue: consistency

Strong formulation

D_h^* , adjoint of D_h with respect to $\langle f, g \rangle_h = \sum_{i \in \mathbb{Z}} \omega_i f_i g_i$:

D_h^* such that $\forall (f, g), \langle D_h(f), g \rangle = - \langle f, D_h^*(g) \rangle_h$

$$\rightsquigarrow D_h^*(g)_i = \sum_{j \in \mathcal{P}} \omega_j (g_i + g_j) W'_{ij} \simeq g'_i$$

this strong formulation **is** conservative!

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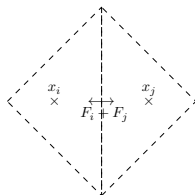
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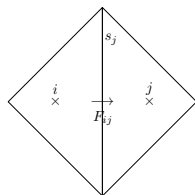
Hybridization SPH - Finite Volumes

SPH



$$\partial_x F_i \simeq \sum_{j \in \mathcal{P}} \omega_j (F_i + F_j) W'_{ij}$$

Finite Volumes



$$\partial_x F_i \simeq \sum_{j \in \gamma(i)} s_j F_{ij}$$

(F_{ij} : **any** conservative FV flux)

Hybrid formulation FV-SPH

$$\partial_x F_i \simeq \sum_{j \in \mathcal{P}} \omega_j 2F_{ij} W'_{ij}$$

Summary

SPH-FV approximation of a PDE

consider a general PDE of the form $\partial_t \Phi + \partial_x F(\Phi) = S(\Phi)$

$$\text{SPH approximation of } \partial_x F(\Phi): \sum_{j \in \mathcal{P}} 2\omega_j F_{ij} W'_{ij}$$

- F_{ij} : any conservative FV flux from particle i to particle j
- choice to make to discretize $\partial_t \Phi$ and $S(\Phi)$
- conservative flux discretization

The shallow-water equations with topography

$$\begin{cases} \partial_t h + \partial_x(hu) & = 0 \\ \partial_t(hu) + \partial_x\left(hu^2 + \frac{1}{2}gh^2\right) & = -gh\partial_x Z \end{cases}$$

where

- $h \geq 0$: water height
- $u \in \mathbb{R}$: water velocity in the x direction
- $g > 0$: gravity constant
- Z : smooth topography

they can be rewritten as $\partial_t \Phi + \partial_x F(\Phi) = S(\Phi)$, with

$$\Phi = \begin{pmatrix} h \\ hu \end{pmatrix}, F(\Phi) = \begin{pmatrix} hu \\ hu^2 + \frac{1}{2}gh^2 \end{pmatrix}, S(\Phi) = \begin{pmatrix} 0 \\ -gh\partial_x Z \end{pmatrix}$$

Discretization

goal: discretize $\partial_t \Phi + \partial_x F(\Phi) = S(\Phi)$

$(x, t) \in \mathbb{R} \times \mathbb{R}_+ \longrightarrow (x_i, t^n)$, with $(i, n) \in \mathbb{Z} \times \mathbb{N}$ and steps Δx and Δt

$\Phi(x, t) \longrightarrow \Phi_i^n$

$\partial_t \Phi \longrightarrow \frac{\Phi_i^{n+1} - \Phi_i^n}{\Delta t}$ (explicit Euler)

$\partial_x F(\Phi) \longrightarrow \sum_{j \in \mathcal{P}} 2\omega_j F_{ij} W'_{ij}$ (SPH discretization)

$S(\Phi) \longrightarrow S_i$ (any discretization)

Discretized shallow-water equations

Hybrid scheme applied to the shallow-water equations

$$\begin{cases} \frac{h_i^{n+1} - h_i^n}{\Delta t} + \sum_{j \in \mathcal{P}} 2\omega_j (hu)_{ij} W'_{ij} & = 0 \\ \frac{h_i^{n+1} u_i^{n+1} - h_i^n u_i^n}{\Delta t} + \sum_{j \in \mathcal{P}} 2\omega_j \left(hu^2 + \frac{1}{2} gh^2 \right)_{ij} W'_{ij} & = s_i \end{cases}$$

- s_i : discretization of $-gh\partial_x Z$
- $F^{\Delta x}(\Phi_i^n, \Phi_j^n)$: numerical flux, such that

$$F^{\Delta x}(\Phi_i^n, \Phi_j^n) = \begin{pmatrix} (hu)_{ij} \\ (hu^2 + \frac{1}{2} gh^2)_{ij} \end{pmatrix}$$

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The lake at rest steady state

Steady states

a solution Φ of a PDE will be a *steady state* iff

$$\partial_t \Phi = 0$$

for the shallow-water equations:

$$\begin{cases} \partial_t h & = 0 \\ \partial_t(hu) & = 0 \end{cases} \Rightarrow \begin{cases} \partial_x(hu) & = 0 \\ \partial_x(hu^2 + \frac{1}{2}gh^2) & = -gh\partial_x Z \end{cases}$$

Lake at rest steady state

$$\begin{cases} u & = 0 \\ h + Z & = cst \end{cases} \quad (\text{lake at } rest)$$

Well-balanced schemes

Well-balancedness

the scheme will be *well-balanced* for a steady state iff

for Φ^n verifying the steady state, $\forall i \in \mathbb{Z}, \Phi_i^{n+1} = \Phi_i^n$

Properties to be verified by the desired hybrid scheme

- well-balancedness (preservation of the lake at rest steady state)
- conservativity
- based on SPH approximations

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Variables H and X

consider the following variables (Berthon-Foucher, 2012):

- $H = h + Z$: the free surface
- $X = \frac{h}{H}$: the water volume fraction

from now on, $V = \begin{pmatrix} H \\ Hu \end{pmatrix}$ is a new set of variables

Case of the lake at rest

for the lake at rest steady state, V is *constant throughout the domain*: indeed,

$$\left. \begin{array}{l} u = 0 \\ H = h + Z = cst \end{array} \right\} \Rightarrow V = \begin{pmatrix} H \\ 0 \end{pmatrix} = cst$$

Variables H and X

rewrite the shallow-water equations for weak solutions:

$$\begin{cases} \partial_t h + \partial_x(X(Hu)) = 0, \\ \partial_t(hu) + \partial_x\left(X(Hu^2 + \frac{1}{2}gH^2)\right) = \frac{g}{2}\partial_x(hZ) - gh\partial_x Z \end{cases}$$

this system can be written as

$$\partial_t \Phi + \partial_x(XF(V)) = \tilde{S}(\Phi)$$

with *the same flux function F* as the original shallow-water equations

Numerical scheme used on for this reformulation

Hybrid scheme applied to the reformulation

$$\left\{ \begin{array}{l} \frac{h_i^{n+1} - h_i^n}{\Delta t} + \sum_{j \in \mathcal{P}} 2\omega_j X_{ij} (Hu)_{ij} W'_{ij} = 0 \\ \frac{h_i^{n+1} u_i^{n+1} - h_i^n u_i^n}{\Delta t} + \sum_{j \in \mathcal{P}} 2\omega_j X_{ij} \left(Hu^2 + \frac{1}{2} g H^2 \right)_{ij} W'_{ij} = \left(\frac{g}{2} \partial_x (hZ) - gh \partial_x Z \right)_i \end{array} \right.$$

still to be defined:

- X_{ij} , an average of X_i and X_j
- $\left(\frac{g}{2} \partial_x (hZ) - gh \partial_x Z \right)_i$, the source term discretization

Source term discretization

What is left to do?

ensuring the well-balancedness by finding a suitable discretization of the source term $\frac{g}{2}\partial_x(hZ) - gh\partial_x Z$

noting that $h = HX$ and $Z = H(1 - X)$ yields

$$\frac{g}{2}\partial_x(hZ) - gh\partial_x Z = \frac{g}{2}\partial_x(XH^2(1 - X)) - gXH\partial_x(H(1 - X))$$

First naïve approach: brutal SPH discretization

$$\left(\frac{g}{2}\partial_x(hZ) - gh\partial_x Z\right)_i = \frac{g}{2} \sum_{j \in \mathcal{P}} 2\omega_j (X_{ij}H_{ij} - 2\bar{X}_i\bar{H}_i) (H_{ij}(1 - X_{ij})) W'_{ij}$$

still to determine: the averages X_{ij} , H_{ij} , \bar{X}_i , \bar{H}_i

however, this discretization leads to a *non-conservative* scheme!

$$\left(\text{i.e. } \sum_{i \in \mathcal{Z}} \sum_{j \in \mathcal{P}} X_{ij} F^{\Delta x}(V_i^n, V_j^n) \neq \sum_{i \in \mathcal{Z}} \left(\frac{g}{2}\partial_x(hZ) - gh\partial_x Z\right)_i\right)$$

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Idea behind the fix

main idea: add a term to make the previous naïve discretization conservative

$$\text{we have } \sum_{j \in \mathcal{P}} \omega_j W'_{ij} \simeq 0$$

\rightsquigarrow to add a factor of $\sum_{j \in \mathcal{P}} \omega_j W'_{ij}$ is to add a representation of zero

Second approach

$$\begin{aligned} \left(\frac{g}{2} \partial_x (hZ) - gh \partial_x Z \right)_i &= \frac{g}{2} \sum_{j \in \mathcal{P}} 2\omega_j (X_{ij} H_{ij} - 2\bar{X}_i \bar{H}_i) (H_{ij} (1 - X_{ij})) W'_{ij} \\ &\quad + Y_i \sum_{j \in \mathcal{P}} \omega_j W'_{ij} \end{aligned}$$

The correction term

How to choose Y_i ?

such that the additional term *corrects* the non-conservativity
 \rightsquigarrow add what was missing for the naïve approach to be conservative

Final form of the discretization

$$\left(\frac{g}{2} \partial_x (hZ) - gh \partial_x Z \right)_i = \frac{g}{2} \sum_{j \in \mathcal{P}} 2\omega_j (X_{ij} H_{ij} - 2\bar{X}_i \bar{H}_i) (H_{ij} (1 - X_{ij})) W'_{ij} \\ + g \sum_{j \in \mathcal{P}} 2\omega_j \bar{H}_i^2 \tilde{X}_i (1 - \tilde{X}_i) W'_{ij}$$

averages \bar{X}_i , \bar{H}_i , \tilde{X}_i , X_{ij} and H_{ij} still to be determined
 \rightsquigarrow whatever the averages, the scheme *will be conservative!*

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The theorem

Well-balancedness of the scheme

Assume both free surface averages to satisfy:

$$H_{ij} = \bar{H}_i = H, \quad \text{as soon as } H_i = H_j = H.$$

Assume \bar{X}_i is defined by

$$\bar{X}_i = \frac{1}{2} \frac{\sum_j \omega_j X_{ij}^2 W'_{ij}}{\sum_j \omega_j (X_{ij} - 1) W'_{ij} + (\bar{X}_i - 1) \sum_j \omega_j W'_{ij}}.$$

Then the scheme defined by the SPH hybridization and source term discretization preserves the lake at rest.

Outline of the proof

What to prove? *Preservation of the lake at rest*

assume Φ^n at rest, i.e. $\forall i \in \mathbb{Z}, \begin{cases} h_i^n + Z_i = H \geq 0 \\ u_i^n = 0 \end{cases}$

\rightsquigarrow prove that $\forall i \in \mathbb{Z}, \Phi_i^{n+1} = \Phi_i^n$

An important property of numerical fluxes

recall $F^{\Delta x}(\Phi_i^n, \Phi_j^n) = \begin{pmatrix} (hu)_{ij} \\ (hu^2 + \frac{1}{2}gh^2)_{ij} \end{pmatrix}$

$\forall \Phi, F^{\Delta x}(\Phi, \Phi) = F(\Phi)$ (*consistency*)

Outline of the proof

use the consistency: we have, $\forall i \in \mathbb{Z}$, $V_i^n = \begin{pmatrix} H \\ 0 \end{pmatrix}$

therefore, $\forall (i, j) \in \mathbb{Z}^2$,

$$F^{\Delta x}(V_i^n, V_j^n) = F^{\Delta x} \left(\begin{pmatrix} H \\ 0 \end{pmatrix}, \begin{pmatrix} H \\ 0 \end{pmatrix} \right) = F \begin{pmatrix} H \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ \frac{1}{2}gH^2 \end{pmatrix}$$

$$\begin{pmatrix} (Hu)_{ij} \\ (Hu^2 + \frac{1}{2}gH^2)_{ij} \end{pmatrix} = \begin{pmatrix} 0 \\ \frac{1}{2}gH^2 \end{pmatrix}$$

Outline of the proof

recall the scheme:

$$\left\{ \begin{array}{l} \frac{h_i^{n+1} - h_i^n}{\Delta t} + \sum_{j \in \mathcal{P}} 2\omega_j X_{ij} (Hu)_{ij} W'_{ij} = 0 \\ \frac{h_i^{n+1} u_i^{n+1} - h_i^n u_i^n}{\Delta t} + \sum_{j \in \mathcal{P}} 2\omega_j X_{ij} \left(Hu^2 + \frac{1}{2} g H^2 \right)_{ij} W'_{ij} = \\ \left(\frac{g}{2} \partial_x (hZ) - gh \partial_x Z \right)_i \end{array} \right.$$

- $(Hu)_{ij} = 0 \implies h_i^{n+1} = h_i^n$ directly
- $\left(Hu^2 + \frac{1}{2} g H^2 \right)_{ij} = \frac{1}{2} g H^2 \implies h_i^{n+1} u_i^{n+1} = h_i^n u_i^n$
(after a few easy steps)

\rightsquigarrow the scheme **is** well-balanced!



An important remark: consistency of the average

$$\text{recall that } \bar{X}_i = \frac{1}{2} \frac{\sum_j \omega_j X_{ij}^2 W'_{ij}}{\sum_j \omega_j (X_{ij} - 1) W'_{ij} + (\tilde{X}_i - 1) \sum_j \omega_j W'_{ij}}$$

is \bar{X}_i consistent with X ?

- $\sum_j \omega_j X_{ij}^2 W'_{ij}$ is consistent with $\partial_x X^2$
- $\sum_j \omega_j (X_{ij} - 1) W'_{ij}$ is consistent with $\partial_x (X - 1)$
- $\sum_j \omega_j W'_{ij}$ is consistent with 0

\rightsquigarrow therefore \bar{X}_i is consistent with $\frac{1}{2} \frac{\partial_x X^2}{\partial_x (X - 1)} = X$

Definitions of the other averages

the well-balancedness is *independent from the definitions of the 4 other averages*: some usual expressions can be used

- $\bar{H}_i = H_i^n = h_i^n + Z_i$
- $\tilde{X}_i = X_i^n = \frac{h_i^n}{h_i^n + Z_i}$
- $H_{ij} = \begin{cases} H_i^n & \text{if } (Hu)_{ij} > 0 \\ H_j^n & \text{otherwise} \end{cases} \quad (\text{upwind expression})$
- $X_{ij} = \begin{cases} X_i^n & \text{if } (Hu)_{ij} > 0 \\ X_j^n & \text{otherwise} \end{cases}$

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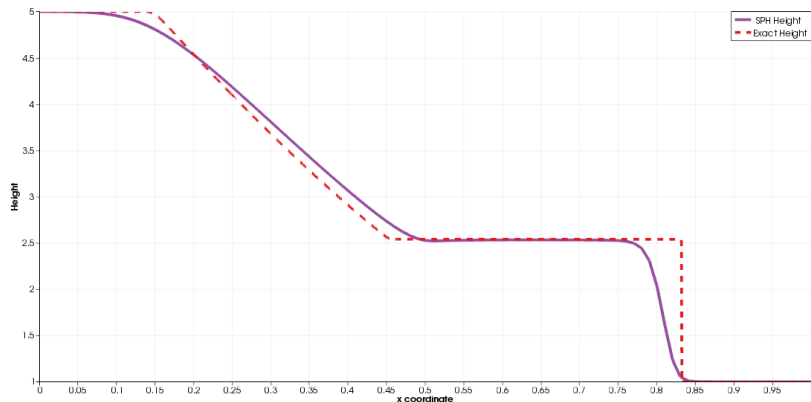
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Inconsistency of the SPH method

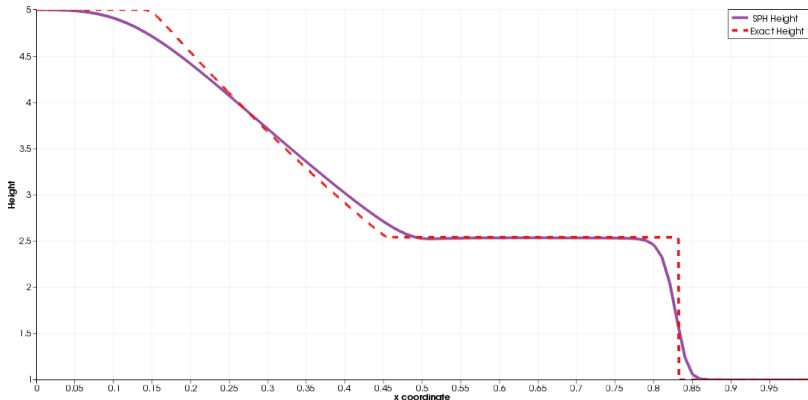
consider the dam-break defined by $\Phi = \Phi_L$ if $x < 0.5$, Φ_R otherwise

$$\Phi_L = \begin{pmatrix} h_L \\ h_L u_L \end{pmatrix} = \begin{pmatrix} 5 \\ 0 \end{pmatrix} \text{ and } \Phi_R = \begin{pmatrix} h_R \\ h_R u_R \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$



Inconsistency of the SPH method

renormalization: $D_{h,r}(f)_i = B_i^{-1} \sum_{j \in \mathcal{P}} \omega_j (f_j - f_i) W'_{ij}$, with $B_i = \sum_{j \in \mathcal{P}} \omega_j (x_j - x_i) W'_{ij}$



Non-well-balancedness in 2D

1 General introduction to the SPH method

- Brief history
- Core of the SPH method
- Discretization of the SPH equations
- Application to shallow-water equations

2 A well-balanced scheme

- Goals
- A change of variables
- Conservative fix
- Main result

3 Numerical results

- Inconsistency and non-well-balancedness
- Validation of the scheme

4 Conclusion and perspectives

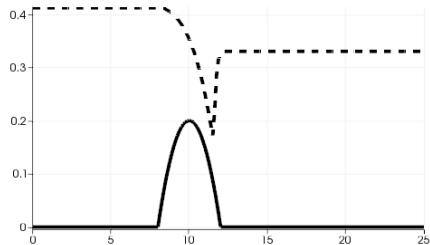
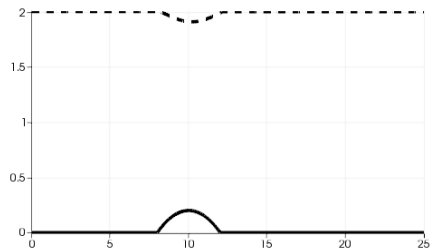
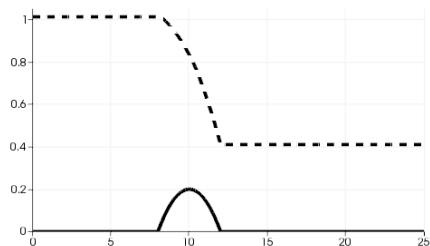
Steady test cases

Three steady test cases (Goutal - Maurel, 1997)

- computational domain: $x \in [0, 25]$
- topography: $Z(x) = \begin{cases} 0.2 - 0.05(x - 10)^2 & \text{if } 8 < x < 12, \\ 0 & \text{otherwise} \end{cases}$
- boundary conditions: $h(0, t)u(0, t) = q_L$ and $h(25, t) = h_R$
- initial conditions: $h(x, 0)u(x, 0) = 0$ and $h(x, 0) + Z(x) = h_R$
- individual values:
 - transcritical flow without shock (TF): $q_L = 1.53$; $h_R = 0.66$
 - transcritical flow with shock: (TFS) $q_L = 0.18$; $h_R = 0.33$
 - subcritical flow (SF): $q_L = 4.42$; $h_R = 2$

steady test cases: a *steady state* (i.e. constant discharge) is obtained after a transition period

Steady test cases



full line: topography
 dashed line: free surface
 results shown after 600s
 top left: TF; bottom left: TFS;
 right: SF

Steady test cases: discharge error

| Test case | Hydrostatic upwind | | SPH scheme | |
|-----------|--------------------|------------------|-------------|------------------|
| | L^2 error | L^∞ error | L^2 error | L^∞ error |
| TF | 5.98E-2 | 1.87E-2 | 5.67E-2 | 1.85E-2 |
| TFS | 4.68E-2 | 2.85E-2 | 5.50E-2 | 4.02E-2 |
| SF | 9.78E-2 | 2.70E-2 | 9.83E-2 | 2.74E-2 |

comparison between the hydrostatic upwind scheme (Berthon - Foucher, 2012) and the SPH scheme (both well-balanced)

Lake at rest without well-balanced SPH scheme

Lake at rest with well-balanced SPH scheme

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Conclusion

- SPH hybrid scheme for a reformulation of the shallow-water equations
- suitable discretization of the topography source term
- scheme still conservative
- well-balancedness confirmed

Perspectives

- multiple dimensions: easy (replace W'_{ij} with ∇W_{ij})
- well-balanced SPH scheme for Euler equations with gravity

Thanks for your attention!