

# Thresholds for entanglement criteria in quantum information theory

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- ▶ Entanglement via positive maps approach
- ▶ Reduction Criterion (RED)
- ▶ Absolute Reduction Criterion (ARED)
- ▶ Approximations of (A)SEP
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# Introduction: Entanglement versus separability

## Entanglement=inseparability

### Separable state:

$$\rho = \sum_i p_i e_i e_i^* \otimes f_i f_i^*, p_i \geq 0, \sum_i p_i = 1, e_i \in \mathbb{C}^n, f_i \in \mathbb{C}^k$$

**Goal:** efficient methods to characterize entangled states

**PPT criterion:**<sup>1</sup> if a state is separable, then the partial transpose respect to one of the subsystems is positive-semidef.

$$SEP := \{\rho_{AB} : \rho_{AB} - sep\} \subset PPT = \{\rho_{AB} / \rho^\Gamma \geq 0\}$$

Tool to detect entanglement: if the partial transpose is not positive-semidefinite, then the state is entangled

**Question:** exists other positive maps  $\varphi$  such that  $(id \otimes \varphi)(\rho) \not\geq 0$ , for some  $\rho$  entangled state

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<sup>1</sup>A. Peres, Separability criterion for density matrices, PRL 77,1996 

# Separability criteria based on positive maps approach

## Mathematical formulation<sup>2</sup>:

$\rho \in \mathcal{M}_n \otimes \mathcal{M}_k$  separable iff  
 $\rho^\varphi := [\text{id} \otimes \varphi](\rho) \geq 0, \forall \varphi \geq 0, \varphi : \mathcal{M}_k \rightarrow \mathcal{M}_m$ , all positive  
integers  $m \in \mathbb{N}$

$$SEP \subset \{ \rho : \rho^\varphi \geq 0 \}$$

- ▶ transposition map  $\varphi \Rightarrow$  Positive Partial Transposition (PPT)
- ▶ reduction map:  $\varphi(X) := I \cdot \text{Tr}X - X \Rightarrow$

## Reduction Criterion (RC)<sup>3</sup>:

$$\rho_{AB} - \text{sep} \Rightarrow \rho_A \otimes I_B - \rho_{AB} \geq 0 \quad , \quad I_A \otimes \rho_B - \rho_{AB} \geq 0 \quad (1)$$

$\rho_A = [\text{id} \otimes \text{Tr}](\rho_{AB})$  partial trace over the second subsystem

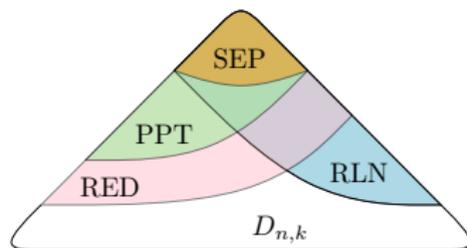
<sup>2</sup>Horodecki M, Separability of mixed states: necessary and sufficient conditions, Phys. Lett A, 1996

<sup>3</sup>Horodecki and all.'99, Cerf and all. '99

# Separability via Reduction Criterion

- ▶  $SEP \subset RED := \{\rho : \rho^{red} := \rho_A \otimes I_B - \rho_{AB} \geq 0\}$
- ▶  $\rho_{AB}$  rank one entangled state, then  $\rho^{red} \not\geq 0$
- ▶  $\rho^\Gamma \geq 0 \Rightarrow \rho^{red} \geq 0, (PPT \subset RED)$
- ▶ If  $dim B = 2$ , then  $PPT = RC$
- ▶ RC connected to entanglement distillation: all states that violate RC are distillable.

$$SEP \subset PPT \subset RED$$



# Absolutely Separable States

Knill's question<sup>4</sup>: *given an self-adjoint positive semi-definite operator  $\rho$ , which are the conditions on the spectrum of  $\rho$  such that  $\rho$  is separable respect to any decomposition?*

**ASEP**: states that remain separable under any unitary transformation

$$ASEP_{n,k} = \bigcap_{U \in \mathcal{U}_{nk}} USEP_{n,k} U^*$$

**Goal**= conditions on the spectrum such that to be separable!

$$APPT_{n,k} := \{\rho \in \mathcal{D}_{n,k} / \forall U \in \mathcal{U}_{nk} : (U\rho U^*)^\Gamma \geq 0\} = \bigcap_{U \in \mathcal{U}_{nk}} UPPT_{n,k} U^*$$

$$ARED_{n,k} := \{\rho \in \mathcal{D}_{n,k} / \forall U \in \mathcal{U}_{nk} : (U\rho U^*)^{red} \geq 0\} = \bigcap_{U \in \mathcal{U}_{nk}} URED_{n,k} U^*$$

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<sup>4</sup>Open Problems in Quantum Information Theory, <http://www.imaph.tu-bs.de/qi/problems>

## Absolutely PPT states

- ▶ necessary and sufficient conditions<sup>5</sup> on the spectrum under which the absolute PPT property holds
- ▶ the condition is to check the positivity of an exponential number of Hermitian matrices (the number of LMI is bounded above by  $e^{2p \ln p(1+o(1))}$ ,  $p = \min(n, k)$ )
- ▶ if  $\rho \in \mathcal{H}_{2n} = \mathcal{H}_2 \otimes \mathcal{H}_n$ , then  $\rho \in APPT$  iff

$$\lambda_1 \leq \lambda_{2n-1} + 2\sqrt{\lambda_{2n}\lambda_{2n-2}}$$

- ▶ if  $\rho \in \mathcal{H}_{3n} = \mathcal{H}_3 \otimes \mathcal{H}_n$ , then  $\rho \in APPT$  iff

$$\begin{pmatrix} 2\lambda_{3n} & \lambda_{3n-1} - \lambda_1 & \lambda_{3n-2} - \lambda_2 \\ \lambda_{3n-1} - \lambda_1 & 2\lambda_{3n-3} & \lambda_{3n-4} - \lambda_3 \\ \lambda_{3n-2} - \lambda_2 & \lambda_{3n-4} - \lambda_3 & 2\lambda_{3n-5} \end{pmatrix} \geq 0, \\ \begin{pmatrix} 2\lambda_{3n} & \lambda_{3n-1} - \lambda_1 & \lambda_{3n-3} - \lambda_2 \\ \lambda_{3n-1} - \lambda_1 & 2\lambda_{3n-2} & \lambda_{3n-4} - \lambda_3 \\ \lambda_{3n-3} - \lambda_2 & \lambda_{3n-4} - \lambda_3 & 2\lambda_{3n-5} \end{pmatrix} \geq 0. \quad (2)$$

<sup>5</sup>Hildebrand, Positive partial transpose from spectra, PRA, 2007.

# Absolutely Reduced pure states

$$\text{ARED}_{n,k} := \{\rho \in D_{n,k} \mid \forall U \in \mathcal{U}_{nk} : (U\rho U^*)^{\text{red}} \geq 0\} \quad (3)$$

## Reduction of pure state<sup>6</sup>:

Given a vector  $\psi \in \mathbb{C}^n \otimes \mathbb{C}^k$  with Schmidt coefficients  $\{x_i\}_{i=1}^r$ , the eigenvalues of the reduced matrix  $(\psi\psi^*)^{\text{red}}$  are

$$\text{spec} \left( (\psi\psi^*)^{\text{red}} \right) = \underbrace{(x_1, \dots, x_1)}_{k-1 \text{ times}}, \eta_1, x_2, \dots, \eta_{r-1}, \underbrace{(x_r, \dots, x_r)}_{k-1 \text{ times}}, \underbrace{(0, \dots, 0)}_{(n-r)k \text{ times}}, \eta_r$$

where  $x_i \geq \eta_i \geq x_{i+1}$  for  $i \in [r-1]$  and  $\eta_r = -\sum_{i=1}^{r-1} \eta_i \leq 0$ .

The set  $\{\eta_i\}_{i=1}^r \setminus \{x_i\}_{i=1}^r$  is the set of solutions  $\eta \in \mathbb{R} \setminus \{x_i\}_{i=1}^r$  of the equation

$$\sum_{i=1}^r \frac{x_i}{x_i - \eta} = 1$$

<sup>6</sup>M.A. Jivulescu, N. Lupa, I. Nechita, D. Reeb, Positive reduction from spectra, Linear Algebra and its Applications, 2014

# Characterizing Absolutely Reduced states

$$\text{ARED}_{n,k} = \{\rho \in D_{n,k} : \forall X \in \Delta_{\min(n,k)}, \langle \lambda_\rho^\downarrow, \hat{x}^\uparrow \rangle \geq 0\}, \quad (4)$$

where  $\lambda_\rho^\downarrow$  is the vect. of the eigenvalues of  $\rho$  and  $\hat{x}^\uparrow$  is the vector of Schmidt coefficients.

$$\hat{x} := (\underbrace{x_1, \dots, x_1}_{k-1 \text{ times}}, \eta_1, \underbrace{x_2, \dots, x_2}_{k-1 \text{ times}}, \dots, \eta_{r-1}, \underbrace{x_r, \dots, x_r}_{k-1 \text{ times}}, \underbrace{0, \dots, 0}_{(n-r)k \text{ times}}, \eta_r),$$

$\eta_i$  are the solutions of the equation  $F_x(\lambda) := \sum_{i=1}^q \frac{m_i x_i}{x_i - \lambda} - 1 = 0$ .

- ▶ necessary and sufficient condition on the spectrum as family of linear inequalities in terms of the spectrum of reduced of a pure state
- ▶ given  $\rho \in M_2(\mathbb{C}) \otimes M_k(\mathbb{C})$ , then  $\rho \in \text{ARED}_{2,k}$  if and only if

$$\lambda_1 \leq \lambda_{k+1} + 2\sqrt{(\lambda_2 + \dots + \lambda_k)(\lambda_{k+2} + \dots + \lambda_{2k})}.$$

## Approximations of SEP

$$\text{SEPBALL}_{n,k} = \left\{ \rho \in D_{n,k} \mid \text{Tr}(\rho^2) \leq \frac{1}{nk-1} \right\}$$

- ▶ largest Euclidian ball<sup>7</sup> inside  $D_{n,k}$ , centered at  $\frac{I}{nk}$
- ▶ contains states on the boundary of  $D_{n,k}$
- ▶ all states within SEPBALL are separable
- ▶ depends only on the spectrum, i.e.  $\text{SEPBALL}_{n,k} \subset \text{ASEP}$
- ▶ it is smaller than other sets

$$\text{GER}_{n,k} = \left\{ \lambda \in \Delta_{nk} : \sum_{i=1}^{r-1} \lambda_i^\downarrow \leq 2\lambda_{nk}^\downarrow + \sum_{i=1}^{r-1} \lambda_{nk-i}^\downarrow \right\}$$

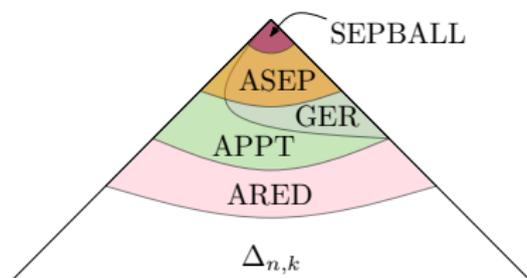
- ▶ the defining equation<sup>8</sup> represents the sufficient condition provided by Gershgorin's theorem for all Hildebrand APPT matrix inequalities to be satisfied
- ▶ lower approximation of APPT, since  $\text{GER}_{n,k} \subset \text{APPT}_{n,k}$
- ▶ provides easily-checkable sufficient condition to be APPT, much simpler than Hildebrand's conditions

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<sup>7</sup>Gurvitz L., PRA 2002

<sup>8</sup>M.A. Jivulescu, N. Lupa, I. Nechita, D. Reeb, Positive reduction from spectra, Linear Algebra and its Applications, 2014

# Approximations of (A)SEP



$$LS_p := \{\lambda \in \Delta_{nk} : \lambda_1^\downarrow \leq \lambda_{nk-p+1}^\downarrow + \lambda_{nk-p+2}^\downarrow + \cdots + \lambda_{nk}^\downarrow\}$$

- ▶  $LS_p$  -the sets of eigenvalue vectors for which the largest eigenvalue is less or equal than the sum of the  $p$  smallest (arbitrary  $p \in [nk]$ )
- ▶ For  $n, k \geq 3$ ,  $APPT \subseteq LS_3 \subseteq LS_k \subseteq ARED_{n,k} \subseteq LS_{2k-1}$ .

# Threshold concept

**Threshold**=the value  $c$  of the parameter, giving an scaling of the environment, **value at which a sharp phase transition of the system occurs!**

## Mathematical characterisation<sup>9</sup>:

Consider a random bipartite quantum state

$\rho_{AB} \in M_n(\mathbb{C}) \otimes M_k(\mathbb{C})$ , obtained by partial tracing over  $\mathbb{C}^s$  a uniformly distributed, pure state  $x \in \mathbb{C}^n \otimes \mathbb{C}^k \otimes \mathbb{C}^s$ .

When one (or both) of the system dimensions  $n$  and  $k$  are large, a threshold phenomenon occurs:

if  $s \sim cnk$ , then there is a *threshold value*  $c_0$  such that

1. for all  $c < c_0$ , as dimension  $nk$  grows,  $P(\rho_{AB} \text{ satisfies the entangled criterion}) = 0$ ;
2. for all  $c > c_0$ , as dimension  $nk$  grows,  $P(\rho_{AB} \text{ satisfies the entangled criterion}) = 1$ ;

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<sup>9</sup>Aubrun, G. *Partial transposition of random states and non-centered semicircular distributions*. Random Matrices: Theory Appl. (2012)

## Reduction Criterion: RMT approach

$W = XX^* \in M_d(\mathbb{C})$  -Wishart matrix of parameters  $d$  and  $s$ .  
**Wishart matrices – physically reasonable models for random density matrices** on a tensor product space.

The spectral properties of  $\rho^{red} \rightarrow$  *reduced matrix*

$$R = W^{red} := W_A \otimes I_k - W_{AB},$$

where  $W_{AB}$  is a Wishart matrix of parameters  $nk$  and  $s$ ,  $W_A$  is its partial trace with respect to the second subsystem  $B$ .

Issues:

- ▶ study the distribution of the eigenvalues of the random matrix  $R = W^{red}$
- ▶ evaluating the probability that  $R$  is positive semidefinite

# Moment formula for $R$

## Theorem

The moments of the random matrix

$R = W^{red} = W_A \otimes I_k - W_{AB} \in M_{nk}(\mathbb{C})$  are given by<sup>10</sup>

$$\forall p \geq 1, \quad \mathbb{E}\text{Tr}(R^p) = \sum_{\alpha \in \mathcal{S}_p, f \in \mathcal{F}_p} (-1)^{|f^{-1}(2)|} s^{\#\alpha} n^{\#(\gamma^{-1}\alpha)} k^{\mathbf{1}_{f \equiv 1} + \#(P_f^{-1}\alpha)}, \quad (5)$$

(# -number of cycles of  $\alpha$ ,  $f : \{1, \dots, p\} \rightarrow \{1, 2\}$ ).

Examples:

$$\mathbb{E}\text{Tr}(R) = nk(k-1)s$$

$$\mathbb{E}\text{Tr}(R^2) = (k-2)[(ks)^2n + ksn^2] + nks^2 + (nk)^2s.$$

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<sup>10</sup>M.A. Jivulescu, N. Lupa, I. Nechita, On the reduction criterion for random quantum states, Journal of Mathematical Physics, Volume: 55, Issue: 11, 2014

# Moment formula

The proof is based on

- ▶ the development of  $R^p$  using the non-commutative binomial formula  $R^p = \sum_{f \in \mathcal{F}_p} (-1)^{|f^{-1}(2)|} R_f$
- ▶  $f : \{1, \dots, p\} \rightarrow \{1, 2\}$  encodes the choice of the term (choose the  $f(i)$ -th term) in each factor in the product

$$R^p = (W_A \otimes I_k - W_{AB})(W_A \otimes I_k - W_{AB}) \cdots (W_A \otimes I_k - W_{AB}).$$

- ▶  $R_f$  denotes the ordered product

$$R_f = R_{f(1)} R_{f(2)} \cdots R_{f(p)} = \overrightarrow{\prod_{1 \leq i \leq p} R_{f(i)}},$$

for the two possible values of the factors

$$R_1 = W_A \otimes I_k, R_2 = W_{AB}.$$

- ▶ Wick graphical calculus to compute  $\mathbb{E} \operatorname{Tr} R_f$ ;

# Three asymptotics regimes

We aim to study the behavior of the combinatorial powers of  $n$ ,  $k$  and  $s$  (dominant term) from the moment formula in three asymptotics regimes:

**Balanced asymptotics:**  $(\exists) \quad c, t > 0$  such that

$$n \rightarrow \infty; k \rightarrow \infty, \quad k/n \rightarrow t; s \rightarrow \infty, \quad s/(nk) \rightarrow c.$$

**Unbalanced asymptotics, first case:**  $(\exists) \quad c > 0$  such that

$$n - \text{fixed}; k \rightarrow \infty; s \rightarrow \infty, \quad s/(nk) \rightarrow c.$$

**Unbalanced asymptotics, second case:**  $(\exists) \quad c > 0$  such that

$$n \rightarrow \infty; k - \text{fixed}; s \rightarrow \infty, \quad s/(nk) \rightarrow c.$$

# Asymptotics regimes: balanced and unbalanced, first case

**Balanced asymptotics:** : the spectrum of the reduced Wishart matrix  $R$  becomes trivial when  $n \rightarrow \infty$ , in the sense that  $R/(ks) \approx I$ .

**Unbalanced asymptotics: first case:** the spectrum of the reduced Wishart matrix  $R$  becomes trivial when  $k \rightarrow \infty$ , in the sense that  $R/(ks) \approx I$ .

Asymptotically, all random quantum states satisfy the reduction criterion ( $C_{red} = 0$ ).

## Three asymptotics regimes: unbalanced, second case

### Theorem

*The moments of the rescaled random matrix  $R$  converge to the following combinatorial quantity:*

$$\forall p \geq 1, \quad \lim_{n \rightarrow \infty} \mathbb{E} \frac{1}{nk} \text{Tr} \left( \frac{R}{n} \right)^p = \sum_{\alpha \in \text{NC}(p)} \prod_{b \in \alpha} c \left[ (1-k)^{|b|} + k^2 - 1 \right]. \quad (6)$$

*Therefore, the empirical eigenvalue distribution  $\mu_n$  of  $\frac{R}{n}$  converges, in moments, to a compound free Poisson distribution  $\mu_{k,c} = \pi_{\nu_{k,c}}$ , where*

$$\nu_{k,c} = c\delta_{1-k} + c(k^2 - 1)\delta_1.$$

*Moreover, the above convergence holds in a strong sense: the extremal eigenvalues converge, almost surely, to the edges of the support of the limiting measure  $\mu_{k,c}$ .*

Its support is positive if and only if  $c > c_{\text{red}} := \frac{(1+\sqrt{k+1})^2}{k(k-1)}$ .

# Thresholds for RED in different asymptotics regimes

1. **Unbalanced asymptotics, second case**<sup>11</sup> : ( $n \rightarrow \infty, k$  fixat,  $s \sim cnk$ ),

$$C_{red} = \frac{(\sqrt{k+1} + 1)^2}{k(k-1)}$$

2. **Balanced asymptotics**<sup>12</sup>: ( $n, k_n \rightarrow \infty, s \sim cn$ )

$$C_{red} = 1;$$

3. **Unbalanced asymptotics, first case**<sup>12</sup>: ( $k \rightarrow \infty, n, s$  fixed)

$$C_{red} = n.$$

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<sup>11</sup>M.A. Jivulescu, N. Lupa, I. Nechita, On the reduction criterion for random quantum states, Journal of Mathematical Physics, 2014

<sup>12</sup>M.A. Jivulescu, N. Lupa, I. Nechita, Thresholds for reduction-related entanglement criteria in quantum information theory, Quantum Information and Computation, 2015:

# Spectrum of Wishart matrices

## Theorem

Let  $\{\lambda_i\}$ ,  $\lambda_i \geq 0$  the eigenvalues of a Wishart matrix  $(d, s)$ .

Then, when  $d \rightarrow \infty$  and  $s = s_d \sim cd$  for some constant  $c > 0$ ,

1. The empirical eigenvalue distribution  $\mu_d = \frac{1}{d} \sum_{i=1}^d \delta_{d^{-1}\lambda_i}$  converges to Marčenko-Pastur distribution

$$\pi_c = \max(1-c, 0)\delta_0 + \sqrt{4c - (x-1-c)^2} \mathbf{1}_{[(\sqrt{c}-1)^2, (\sqrt{c}+1)^2]}(x) dx;$$

2. For any function  $j_d = o(d)$ , almost surely, as  $d \rightarrow \infty$ , the rescaled eigenvalues  $\tilde{\lambda}_i = d^{-1}\lambda_i$  have the following limits<sup>13</sup>

$$\tilde{\lambda}_d, \tilde{\lambda}_{d-1}, \dots, \tilde{\lambda}_{d-j_d+1} \rightarrow a_c = \begin{cases} 0, & \text{if } c \leq 1, \\ (\sqrt{c}-1)^2, & \text{if } c > 1, \end{cases}$$

and

$$\tilde{\lambda}_1, \tilde{\lambda}_2, \dots, \tilde{\lambda}_{j_d} \rightarrow b_c = (\sqrt{c}+1)^2;$$

# Comparing entanglement criteria via thresholds

	$n, k \rightarrow \infty$	$m = \min(n, k)$ fixed, $\max(n, k) \rightarrow \infty$	
SEP <sup>14</sup>	$n^3 \lesssim s \lesssim n^3 \log^2 n$	$mnk \lesssim s \lesssim mnk \log^2(nk)$	
PPT <sup>15</sup>	$s \sim cnk$	$s \sim cnk$	
	$c = 4$	$c = 2 + 2\sqrt{1 - \frac{1}{m^2}}$	
RLN <sup>16</sup>	$s \sim cnk$	$s$ fixed	
	$c = (8/3\pi)^2$	$s = m^2$	
RED	$s \sim cn$	$m = n, s$ is fixed	$m = k, s = cnk$
	$c = 1$	$s = n$	$c = \frac{(1+\sqrt{k+1})^2}{k(k-1)}$

Conclusion: RLN is weaker than PPT (asymptotically);

<sup>14</sup>Aubrun, G., Szarek, S.J., and Ye, D. *Entanglement thresholds for random induced states*. Comm. Pure Appl. Math. (2014).

<sup>15</sup>Aubrun, G. *Partial transposition of random states and non-centered semicircular distributions*. Random Matrices: Theory Appl. (2012).

<sup>16</sup>Aubrun, G. and Nechita, I. *Realigning random states*, J.M.P. (2012).

# Thresholds for ARED in different asymptotics regimes<sup>17</sup>

## 1. Unbalanced asymptotics, first case :

( $k \rightarrow \infty, n - \text{fixed}, s \sim ck$ )

$$c_{ared} = n - 2$$

## 2. Balanced asymptotics: ( $n, k_n \rightarrow \infty, s \sim cnk$ )

$$c_{ared} = 1;$$

## 3. Unbalanced asymptotics, second case : ( $n \rightarrow \infty, k$ fixed, $s \sim cnk$ ),

$$c_{ared} = \left(1 + \frac{2}{k} + \frac{2}{k}\sqrt{k+1}\right)^2$$

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<sup>17</sup>M.A. Jivulescu, N. Lupa, I. Nechita, Thresholds for reduction-related entanglement criteria in quantum information theory, QIC, 2015: 

# Comparing entanglement criteria via thresholds

	$n, k \rightarrow \infty$	$m = \min(n, k)$ fixed $\max(n, k) \rightarrow \infty$	
APPT <sup>18</sup>	$s \sim c \min(n, k)^2 nk$	$s \sim cnk$	
	$c = 4$	$c = \left(m + \sqrt{m^2 - 1}\right)^2$	
GER	$s \sim c \min(n, k)^2 nk$	$s \sim cnk$	
	$c = 4$	$c = \left(m + \sqrt{m^2 - 1}\right)^2$	
ARED	$s \sim cnk$	$m = n, s \sim ck$	$m = k, s = cnk$
	$c = 1$	$c = n - 2$	$c = \left(\frac{(2+k+2)\sqrt{k+1}}{k}\right)^2$

- ▶ the thresholds for GER and APPT are the same, strengthening the claim that GER is a good approximation of APPT

<sup>18</sup>Collins, B., Nechita, I., and Ye, D. *The absolute positive partial transpose property for random induced states*. Random Matrices: Theory Appl. 01, 1250002 (2012).

## Thresholds for related sets

	$d = nk \rightarrow \infty$		
SEPBALL	$s \sim cd^2$		
	$c = 1$		
$LS_p$	$s \sim cd$		
	$p \geq 2$ fixed	$1 \ll p = o(d)$	$p = \lfloor td \rfloor$
	$c = \left(1 + \frac{2}{\sqrt{p-1}}\right)^2$	$c = 1$	$c = 1 - t$

- ▶ both sets depend on the product  $d = nk$ , so it is sufficient to consider one asymptotic regime  $d \rightarrow \infty$
- ▶ SEPBALL case: the size of the environment  $s_d$  scales like the square of the total system  $d = nk$

# Perspectives

- ▶ description of ARLN and to find thresholds for ARLN
- ▶ integrate into the theory the recent results about thresholds for k-extendibility criterion<sup>19</sup>
- ▶ validate/invalidate the conjecture that  $ASEP=APPT$ <sup>20</sup>

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<sup>19</sup>Lancien C. k-extendibility of high-dimensional bipartite quantum states, arxiv: 1504. 06459

<sup>20</sup>Arunachalam, S., Johnston, N., and Russo, V. *Is absolute separability determined by the partial transpose?* Quantum Inform. Comput. 2015

**Thank you for your attention!**

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