

# Hysteresis in Two-Phase Solids

## Numerical Analysis Based on Variational Model due to Mielke, Theil, Levitas

A mixture  $\chi$  is a  $N$ -tuple of characteristic functions

$$\chi = (\chi^{(1)}, \dots, \chi^{(N)}) \in \Xi \subset L^\infty(\Omega; \{0, 1\}^N), \text{ i.e.}$$

$\chi^{(j)}(x) = 1$  if the phase  $j$  is found at material point  $x$  and  $\chi^{(j)}(x) = 0$  otherwise. Below  $N=2$ .

Total free energy of arrangement  $\chi \in \Xi$  at time  $t$

$$\mathcal{E}(t, \chi) := \inf_{v \in V} \left\{ \int_{\Omega} \sum_{j=1}^N \chi^{(j)} W_j(\varepsilon(v)) dx + \mathcal{L}(t, v) \right\}$$

Re-arrangement of the phase configuration  $\chi$  to the actual arrangement  $\zeta$  dissipates energy

$$\mathcal{D}(\chi, \zeta) := \int_{\Omega} D(\chi, \zeta) dx := \int_{\Omega} \kappa |\chi - \zeta| dx.$$

Maximum dissipation principle and stability concept result in Incremental Problem.

## Cont. Hysteresis in Two-Phase Solids

### Incremental Problem

Given initial phase arrangement  $\chi_0$  and time-discretization  $0 < t_1 < \cdots < t_{j-1} < t_j < \cdots < t_J = T$ , seek minimizer  $\chi_j \in \Xi$ ,  $j = 1, 2, \dots, J$ , of

$$(IP) \quad \left\{ \mathcal{E}(t_j, \zeta) + \mathcal{D}(\chi_{j-1}, \zeta) \right\} \quad \text{over } \zeta \in \Xi.$$

Nonconvexity of  $\Xi$  results in nonattainment and generalized solutions  $\theta_j \in \Xi^{**}$ : Seek minimizer  $\chi_j \in \Xi$  of

$$(RIP) \quad \text{minimize } \mathcal{I}^{rlx}(t, \theta_{j-1}, \zeta) \text{ over } \zeta \in \Xi^{**},$$

where

$$\begin{aligned} \mathcal{I}^{rlx}(t, \eta, \theta) = \inf_{v \in V} & \left\{ \int_{\Omega} W_{\theta}^{qc}(Dv) dx \right. \\ & \left. + \mathcal{L}(t, v) \right\} + \int_{\Omega} \kappa |\eta - \theta| dx. \end{aligned}$$

## Cont. Hysteresis in Two-Phase Solids

### Fully Discrete Algorithm

(i) Compute a minimizer  $u_j \in V_h$  of

$$\begin{aligned} & \int_{\Omega} \left( W_2(\varepsilon(v_h)) + 2\gamma H(\vartheta_{j-1}, \ell(\varepsilon(v_h))) \right) dx \\ & + L(t_j, v_h) \quad \text{over } v_h \in V_h. \end{aligned}$$

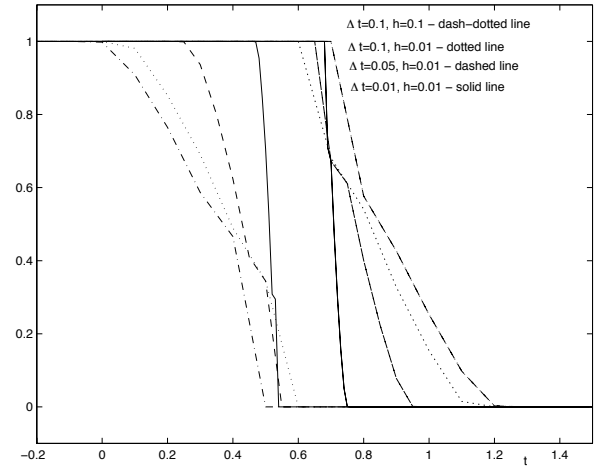
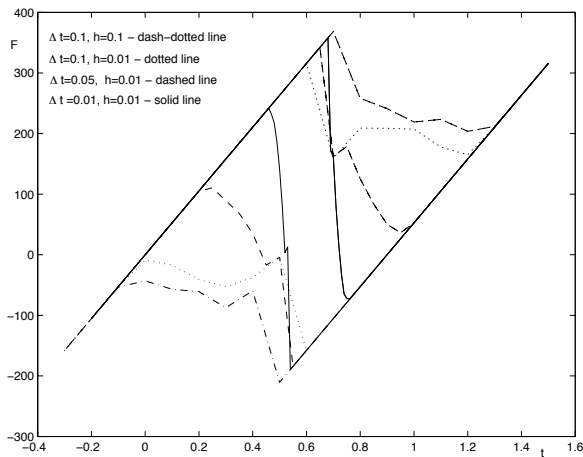
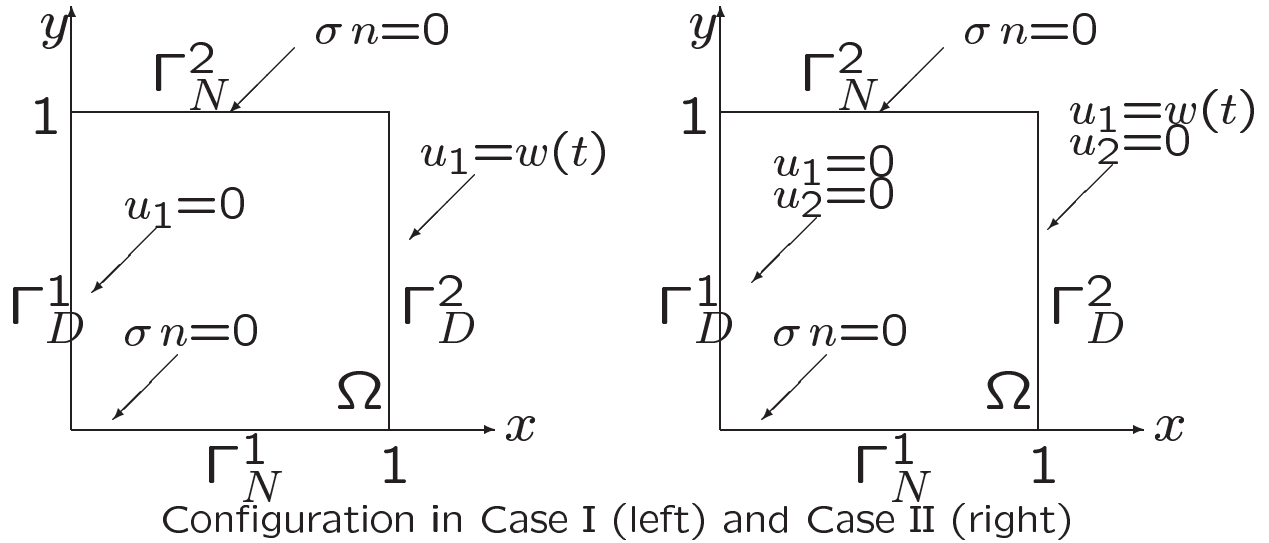
(ii) Compute new arrangements  $\theta_j =: (\vartheta_j, 1 - \vartheta_j)$

$$\vartheta_j = M(\vartheta_{j-1}, \ell(\varepsilon(u_j))).$$

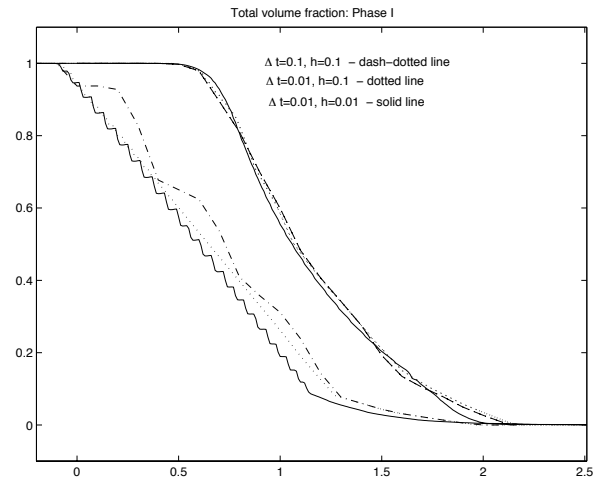
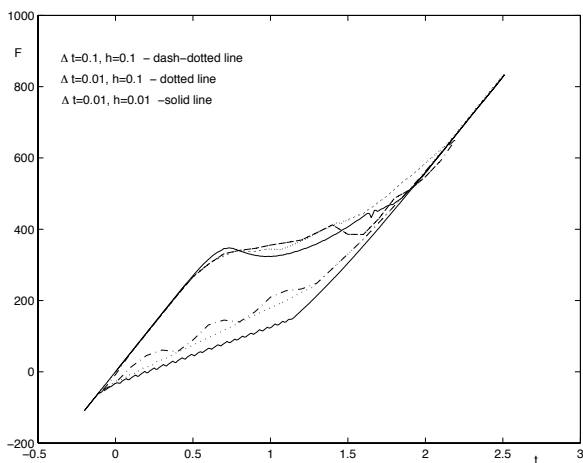
Explicit formulae include, for each  $E \in \mathbb{R}_{sym}^{d \times d}$  and real  $r$  and  $s$ ,

$$\begin{aligned} W_j(E) &= \langle E - E_j, \mathbb{C}(E - E_j) \rangle + W_j^0, \\ \ell(E) &= \frac{1}{2\gamma} (W_2(E) - W_1(E)) + \frac{1}{2}, \\ H(r, s) &= \begin{cases} \frac{\kappa}{2\gamma} r & \text{if } s \leq -\frac{\kappa}{2\gamma}, \\ \frac{\kappa}{2\gamma} r - \frac{1}{2} (s + \frac{\kappa}{2\gamma})^2 & \text{if } -\frac{\kappa}{2\gamma} \leq s \leq r - \frac{\kappa}{2\gamma}, \\ \frac{1}{2} r^2 - rs & \text{if } r - \frac{\kappa}{2\gamma} \leq s \leq r + \frac{\kappa}{2\gamma}, \\ -\frac{\kappa}{2\gamma} r - \frac{1}{2} (s - \frac{\kappa}{2\gamma})^2 & \text{if } r + \frac{\kappa}{2\gamma} \leq s \leq 1 + \frac{\kappa}{2\gamma}, \\ \frac{\kappa}{2\gamma} (1 - r) + \frac{1}{2} - s & \text{if } s \geq 1 + \frac{\kappa}{2\gamma}, \end{cases} \\ M(r, s) &= \begin{cases} 0 & \text{if } s \leq -\frac{\kappa}{2\gamma}, \\ s + \frac{\kappa}{2\gamma} & \text{if } -\frac{\kappa}{2\gamma} \leq s \leq r - \frac{\kappa}{2\gamma}, \\ r & \text{if } r - \frac{\kappa}{2\gamma} \leq s \leq r + \frac{\kappa}{2\gamma}, \\ s - \frac{\kappa}{2\gamma} & \text{if } r + \frac{\kappa}{2\gamma} \leq s \leq 1 + \frac{\kappa}{2\gamma}, \\ 1 & \text{if } s \geq 1 + \frac{\kappa}{2\gamma}. \end{cases} \end{aligned}$$

## Cont. Hysteresis in Two-Phase Solids Numerical Results



Load-displacement diagrams (left) and global volume fractions along the loop (right) for Material B in Case I.



Load-displacement diagrams (left) and global volume fractions along the loop (right) for Material B in Case II.