

# Large time behavior of finite difference schemes for the transport equation

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We consider the transport equation ( $c > 0$ )

$$\begin{aligned}\partial_t u + c \partial_x u &= 0, & t \geq 0, & \quad x \in \mathbb{R}, \\ u_{t=0} &= u_0.\end{aligned}$$

We set  $\Delta t, \Delta x > 0, \lambda := \frac{\Delta t}{\Delta x}$ .

Example of explicit one-step finite difference approximation:

We consider  $n \in \mathbb{N}, j \in \mathbb{Z}$ .

- Upwind scheme:

$$u_j^{n+1} = c\lambda u_{j-1}^n + (1 - c\lambda)u_j^n.$$

- Modified Lax Friedrichs scheme: ( $0 < D < \frac{1}{2\lambda}$ )

$$u_j^{n+1} = \left(\lambda D + \frac{c\lambda}{2}\right) u_{j-1}^n + (1 - 2D\lambda)u_j^n + \left(\lambda D - \frac{c\lambda}{2}\right) u_{j+1}^n.$$

- O3 scheme

$$\begin{aligned}u_j^{n+1} &= -\frac{c\lambda(1 - (c\lambda)^2)}{6} u_{j-2}^n + \frac{c\lambda(1 + c\lambda)(2 - c\lambda)}{2} u_{j-1}^n \\ &\quad + \frac{(1 - (c\lambda)^2)(2 - c\lambda)}{2} u_j^n - \frac{c\lambda(1 - c\lambda)(2 - c\lambda)}{6} u_{j+1}^n.\end{aligned}$$

Modified Lax Friedrichs scheme:  $c\lambda = \frac{1}{2}$  and  $u^0 = \delta := (\delta_{j,0})_{j \in \mathbb{Z}}$

Decay rate :  $\frac{1}{\sqrt{n}}$

O3 scheme:  $c\lambda = \frac{1}{2}$  and  $u^0 = \delta := (\delta_{j,0})_{j \in \mathbb{Z}}$

Decay rate :  $\frac{1}{n^{\frac{1}{4}}}$

Lax Wendroff scheme:  $c\lambda = \frac{1}{2}$  and  $u^0 = \delta := (\delta_{j,0})_{j \in \mathbb{Z}}$

- 1 Introduction
- 2 Generalization of the Local Limit Theorem
- 3 Idea of the proof : Spatial dynamics

We consider the Modified Lax Friedrichs scheme

$$u_j^{n+1} = \left( \lambda D + \frac{c\lambda}{2} \right) u_{j-1}^n + (1 - 2D\lambda) u_j^n + \left( \lambda D - \frac{c\lambda}{2} \right) u_{j+1}^n$$

We consider the Modified Lax Friedrichs scheme

$$u_j^{n+1} = \underbrace{\left(\lambda D + \frac{c\lambda}{2}\right)}_{= \mathbf{a}_1} u_{j-1}^n + \underbrace{(1 - 2D\lambda)}_{= \mathbf{a}_0} u_j^n + \underbrace{\left(\lambda D - \frac{c\lambda}{2}\right)}_{= \mathbf{a}_{-1}} u_{j+1}^n$$



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Thus,

$$\forall n \in \mathbb{N}, \quad u^{n+1} = \mathbf{a} * u^n,$$

where  $\mathbf{a} = (\dots, 0, \mathbf{a}_{-1}, \mathbf{a}_0, \mathbf{a}_1, 0, \dots) \in \mathbb{C}^{\mathbb{Z}}$ .

An explicit one-step finite difference approximation can be written as

$$\begin{aligned}\forall n \in \mathbb{N}, \quad u^n &= \mathbf{a}^n * u^0, \\ u^0 &\in \mathbb{C}^{\mathbb{Z}},\end{aligned}$$

where  $\mathbf{a} \in \mathbb{C}^{\mathbb{Z}}$  is finitely supported and  $\mathbf{a}^n = \mathbf{a} * \cdots * \mathbf{a}$ .

**Goal:** We want to study  $\mathbf{a}^n$  for large values of  $n$  to better understand the large time behavior of the finite difference scheme.

## A detour through Probability - The Local Limit Theorem

We consider a random walk

$$S_n := X_1 + \cdots + X_n$$

where  $X_n$  are i.i.d. random variables (same law as some random variable  $X$  with values in  $\mathbb{Z}$ ) and we use the notation

$$\forall j \in \mathbb{Z}, \quad \mathbf{a}_j := P(X = j).$$

We then have

$$\forall n \in \mathbb{N}, \forall j \in \mathbb{Z}, \quad \mathbf{a}_j^n = P(S_n = j).$$

Central Limit Theorem :

$$\sqrt{n} \left( \frac{S_n}{n} - \mathbb{E}(X) \right) \xrightarrow{\mathcal{L}} \mathcal{N}(0, V(X)).$$

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Local Limit Theorem : Under suitable conditions on the sequence  $\mathbf{a}$

$$\mathbf{a}_j^n - \frac{1}{\sqrt{2\pi V(X)n}} \exp\left(-\frac{|j - n\mathbb{E}(X)|^2}{2nV(X)}\right) \underset{n \rightarrow +\infty}{=} o\left(\frac{1}{\sqrt{n}}\right) \quad \text{uniformly on } \mathbb{Z}.$$

## A detour through Probability - The Local Limit Theorem

$$\mathbf{a}_j^n \sim \frac{1}{\sqrt{2\pi V(X)n}} \exp\left(-\frac{|j - n\mathbb{E}(X)|^2}{2nV(X)}\right) - \frac{1}{n} q\left(\frac{j - n\mathbb{E}(X)}{\sqrt{V(X)n}}\right) \stackrel{n \rightarrow +\infty}{=} o\left(\frac{1}{n}\right),$$

where

$$\forall x \in \mathbb{R}, \quad q(x) := C(X)(x^3 - 3x)e^{-\frac{x^2}{2}}$$

with  $C(X)$  a constant depending on the random variable  $X$ .

**Problem:** For a finite difference scheme, the sequence  $\mathbf{a}$  does not always have non negative or even real coefficients.

Thus, we are looking for a generalization of the Local Limit Theorem.

We also want to find precise bounds for  $\mathbf{a}_j^n$ .

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## Fourier Series

For  $\mathbf{a} \in \ell^1(\mathbb{Z})$ , we introduce  $F_{\mathbf{a}}$  the Fourier series associated to  $\mathbf{a}$ :

$$\forall \kappa \in \mathbb{S}^1, \quad F_{\mathbf{a}}(\kappa) := \sum_{k \in \mathbb{Z}} \mathbf{a}_k \kappa^k.$$

If  $\mathbf{a}$  is finitely supported then  $F_{\mathbf{a}}$  can be holomorphically extended on  $\mathbb{C} \setminus \{0\}$ .

The consistency of the scheme implies

$$F_{\mathbf{a}}(1) = 1 \quad \text{and} \quad F'_{\mathbf{a}}(1) = c\lambda.$$

The  $\ell^2$ -stability is equivalent to having that

$$\max_{\mathbb{S}^1} |F_{\mathbf{a}}| \leq 1. \quad (\text{Von Neumann condition})$$

Stronger hypothesis (Dissipation)

We will suppose that

$$\forall \kappa \in \mathbb{S}^1 \setminus \{1\}, \quad |F_{\mathbf{a}}(\kappa)| < 1.$$

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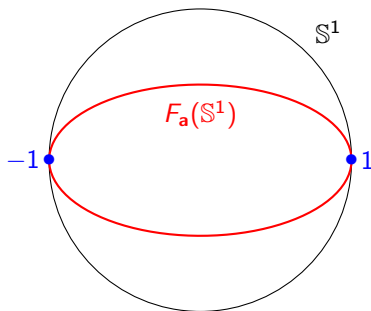
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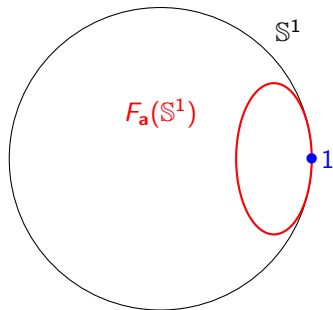
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Lax-Friedrichs scheme,  $c\lambda = \frac{1}{2}$

$$F_a(e^{it}) = \cos(t) + ic\lambda \sin(t).$$



Lax-Wendroff scheme,  $c\lambda = \frac{1}{2}$

$$F_a(e^{it}) = 1 - 2(c\lambda)^2 \sin^2\left(\frac{t}{2}\right) + ic\lambda \sin(t).$$

## Thomée's result

The goal of Thomée was to classify the  $\ell^\infty$  stable schemes. It is equivalent to searching for the sequences  $\mathbf{a}$  such that  $(\mathbf{a}^n)_n$  is bounded in  $\ell^1(\mathbb{Z})$ .

The logarithm of  $F_{\mathbf{a}}$  at 1 as an asymptotic expansion of the form

$$F_{\mathbf{a}}(e^{it}) \underset{t \rightarrow 0}{=} \exp \left( ic\lambda t + iq(t) - \beta t^{2\mu} + o(t^{2\mu}) \right)$$

where  $\beta \in \mathbb{C}$  with  $\Re(\beta) > 0$ ,  $\mu \in \mathbb{N}$  and  $q$  is a polynomial with real coefficients

$$q(t) = b_2 t^2 + \dots + b_{2\mu-1} t^{2\mu-1}.$$

- $ic\lambda$ : Consistency of the scheme
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We have

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Thomée '65 (generalizeable for multiple tangency points)

The family  $(\mathbf{a}^n)_{n \in \mathbb{N}}$  is bounded in  $\ell^1(\mathbb{Z})$  (and the scheme is  $\ell^\infty$  stable) if and only if  $q = 0$ . In this case, there exists algebraic bounds on  $\mathbf{a}_j^n$ .

We can observe that the condition  $q = 0$  implies that the modified equation associated to our scheme is

$$\partial_t u + c \partial_x u = (-1)^{\mu+1} \frac{\beta}{\Delta t} \Delta x^{2\mu} \partial_x^{2\mu} u.$$

We have

$$F_a(e^{it}) \underset{t \rightarrow 0}{=} \exp \left( ic\lambda t + iq(t) - \beta t^{2\mu} + o(t^{2\mu}) \right).$$

Scheme	$iq(t)$	$2\mu$	$\ell^\infty$ stability
Upwind scheme	0	2	Yes
Modified Lax Friedrichs scheme	0	2	Yes
O3 scheme	0	4	Yes
Lax Wendroff scheme	$-i \frac{c\lambda(1-(c\lambda)^2)}{6} t^3$	4	No

**Recent developments** We assume that the logarithm of  $F_a$  has an asymptotic expansion of the form

$$F_a(e^{i\xi}) \underset{\xi \rightarrow 0}{=} \exp\left(ic\lambda\xi - \beta\xi^{2\mu} + o(\xi^{2\mu})\right)$$

where  $\beta \in \mathbb{C}$  with  $\Re(\beta) > 0$  and  $\mu \in \mathbb{N}$ .

### Generalized Gaussian Bounds

Under the previous hypotheses (consistency, dissipation,  $\ell^2$  and  $\ell^\infty$  stability), there exist two constants  $C, c > 0$  such that

$$\forall n \in \mathbb{N} \setminus \{0\}, \forall j \in \mathbb{Z}, \quad |a_j^n| \leq \frac{C}{n^{\frac{1}{2\mu}}} \exp\left(-c \left(\frac{|j - nc\lambda|}{n^{\frac{1}{2\mu}}}\right)^{\frac{2\mu}{2\mu-1}}\right).$$

Due to Diaconis, Saloff-Coste '14, generalization in Coulombel, Faye '21



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We define

$$\forall x \in \mathbb{R}, \quad H_{2\mu}^\beta(x) := \frac{1}{2\pi} \int_{\mathbb{R}} e^{-ixu} e^{-\beta u^{2\mu}} du.$$

It is the fundamental solution of

$$\partial_t u + (-1)^\mu \beta \partial_x^{2\mu} u = 0.$$

For  $\mu = 1$ , we have

$$H_{2\mu}^\beta(x) = \frac{1}{\sqrt{4\pi\beta}} \exp\left(-\frac{|x|^2}{4\beta}\right).$$

It is the heat kernel / normal distribution.

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### Generalized Local Limit Theorem

Under the previous hypotheses, we have

$$\mathbf{a}_j^n - \frac{1}{n^{\frac{1}{2\mu}}} H_{2\mu}^\beta \left( \frac{j - nc\lambda}{n^{\frac{1}{2\mu}}} \right) \underset{n \rightarrow +\infty}{=} o \left( \frac{1}{n^{\frac{1}{2\mu}}} \right)$$

where the error term is uniform on  $\mathbb{Z}$

Due to Randles, Saloff-Coste '15 (also '17 for a multidimensional generalization)

## Main Theorem

Under the previous hypotheses, there exist two constants  $C, c > 0$  such that for all  $n \in \mathbb{N} \setminus \{0\}$  and  $j \in \mathbb{Z}$

$$\left| \mathbf{a}_j^n - \frac{1}{n^{\frac{1}{2\mu}}} H_{2\mu}^\beta \left( \frac{j - nc\lambda}{n^{\frac{1}{2\mu}}} \right) \right| \leq \frac{C}{n^{\frac{1}{\mu}}} \exp \left( -c \left( \frac{|j - nc\lambda|}{n^{\frac{1}{2\mu}}} \right)^{\frac{2\mu}{2\mu-1}} \right).$$

## Corollary

Under the previous hypotheses, there exists a constant  $C > 0$  such that for all  $u_0 \in \ell^2(\mathbb{Z})$

$$\forall n \in \mathbb{N} \setminus \{0\}, \quad \|\mathbf{a}^n * u_0 - H^n * u_0\|_{\ell^2(\mathbb{Z})} \leq \frac{C}{n^{\frac{1}{2\mu}}} \|u_0\|_{\ell^2(\mathbb{Z})},$$

where

$$H^n = \left( \frac{1}{n^{\frac{1}{2\mu}}} H_{2\mu}^\beta \left( \frac{j - nc\lambda}{n^{\frac{1}{2\mu}}} \right) \right)_{j \in \mathbb{Z}}.$$

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Upwind scheme:  $c\lambda = \frac{1}{2}$  and  $u^0 = \delta := (\delta_{j,0})_{j \in \mathbb{Z}}$

## Idea of the proof

- Fourier Analysis : We have for  $n \in \mathbb{N}$ ,  $j \in \mathbb{Z}$ ,

$$\mathbf{a}_j^n = \frac{1}{2\pi} \int_0^{2\pi} e^{ijt} F_{\mathbf{a}}(e^{it})^n dt.$$

We then use integrations by parts and contour deformations to obtain estimates on  $\mathbf{a}_j^n$ .

- Other approach : The goal is to use another representation of the coefficients  $\mathbf{a}_j^n$  using functional calculus.

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- Other approach : The goal is to use another representation of the coefficients  $\mathbf{a}_j^n$  using functional calculus.



We introduce the Laurent operator

$$L_a : u \in \ell^2(\mathbb{Z}) \mapsto a * u \in \ell^2(\mathbb{Z}).$$

Young's convolution inequality implies that this operator is well-defined and bounded. Fourier analysis implies that

$$\sigma(L_a) = F_a(\mathbb{S}^1).$$

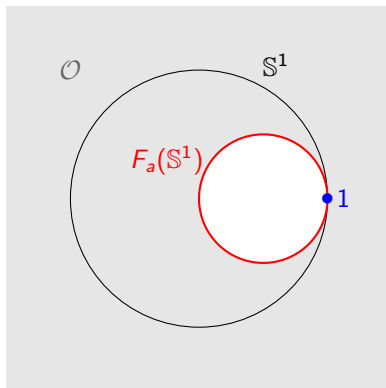
If you consider  $\Gamma$  a path that surrounds  $F_a(\mathbb{S}^1)$ , then for  $n \in \mathbb{N}^*$ ,  $j \in \mathbb{Z}$ ,

$$\mathbf{a}_j^n = \frac{1}{2i\pi} \int_{\Gamma} z^n ((zId - L_a)^{-1} \delta)_j dz. \quad (\text{Inverse Laplace transform})$$

We try our best to choose  $\Gamma$  that is inside the unit disk so that the factor  $z^n$  decreases exponentially. However, this is not possible near 1.

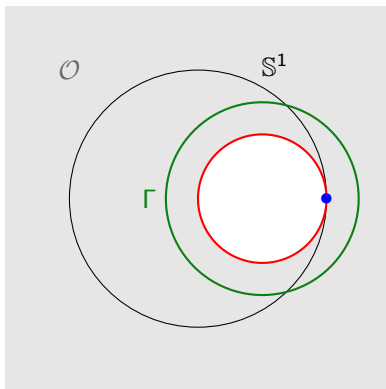
We have

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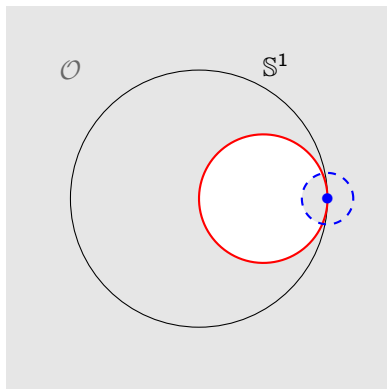
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This will not be a good enough choice of  $\Gamma$ .

We have

$$\forall n \in \mathbb{N}, \forall j \in \mathbb{Z}, \quad \mathbf{a}_j^n = \frac{1}{2i\pi} \int_{\Gamma} z^n ((zId - L_a)^{-1} \delta)_j dz.$$



We need a holomorphic extension of  $((zId - L_a)^{-1} \delta)_j$  in a neighborhood of 1 and a precise choice of  $\Gamma$  depending on  $n$  and  $j$ .

## Perspective

- There are still some results that have not yet been generalized for implicit schemes.
- Could we expect a better asymptotic expansion of  $\mathbf{a}^n$ ?
- Apply the same idea of proof to study the stability near discrete shock profiles for systems of conservation law. (Linearizing near particular non constant solutions, work in progress)

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