

Diffusion d'une particule marquée dans un gaz dilué de sphères dures

Thierry Bodineau

Isabelle Gallagher, Laure Saint-Raymond

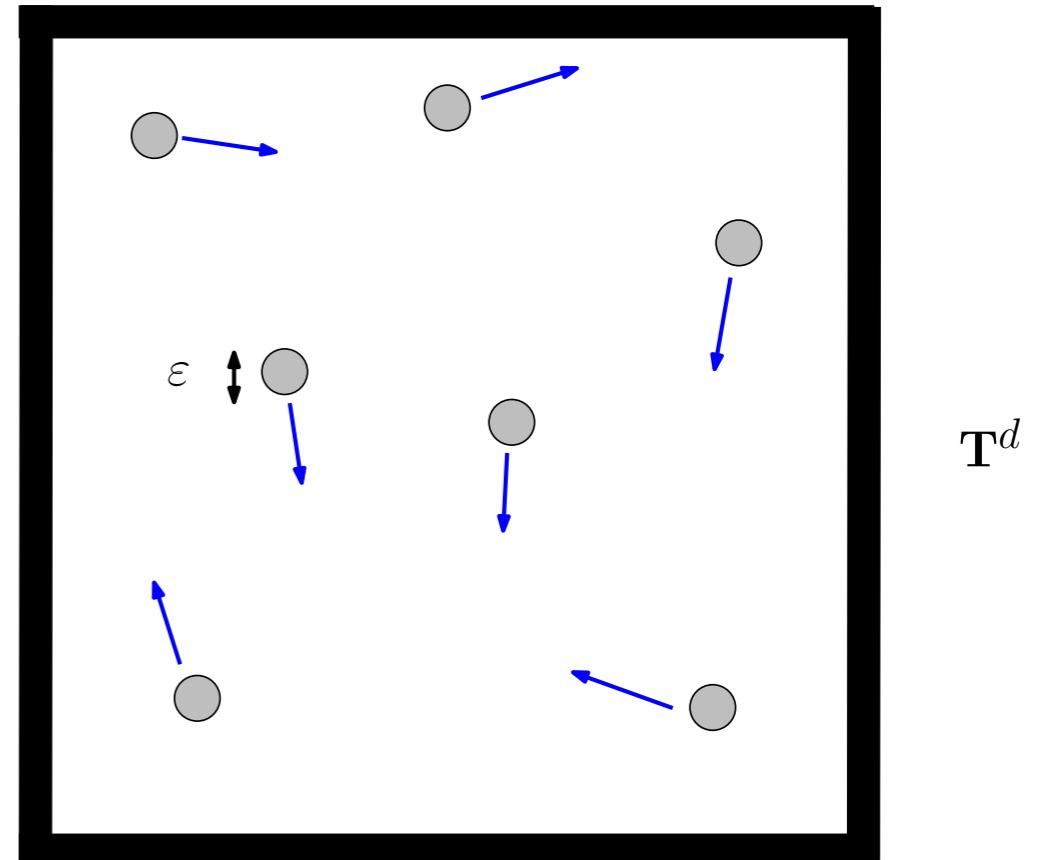
Journées MAS 2014

Outline.

- Particle Diffusion
- Introduction

Diluted Gas of hard spheres

Gas of N hard spheres with deterministic Newtonian dynamics (elastic collisions).



Dimension : $d \geq 2$

Periodic domain: $T^d = [0, 1]^d$

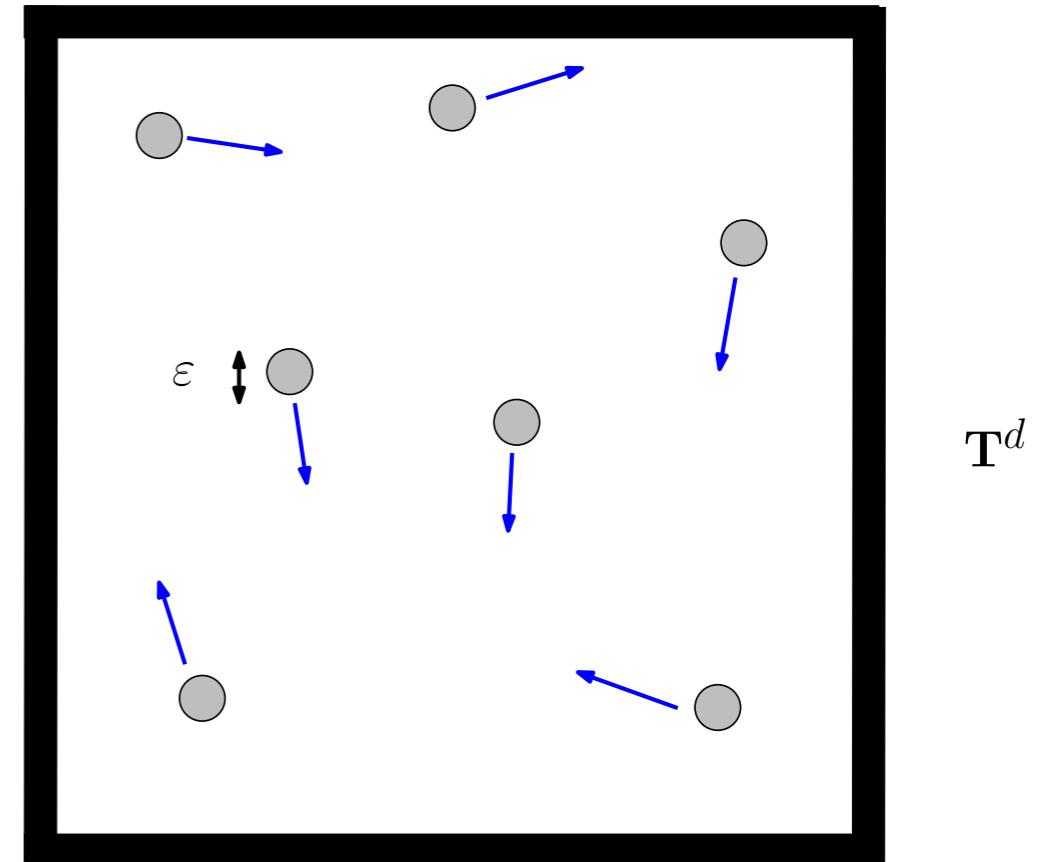
Sphere radius = ε

Boltzmann-Grad scaling

$$N\varepsilon^{d-1} = \alpha$$

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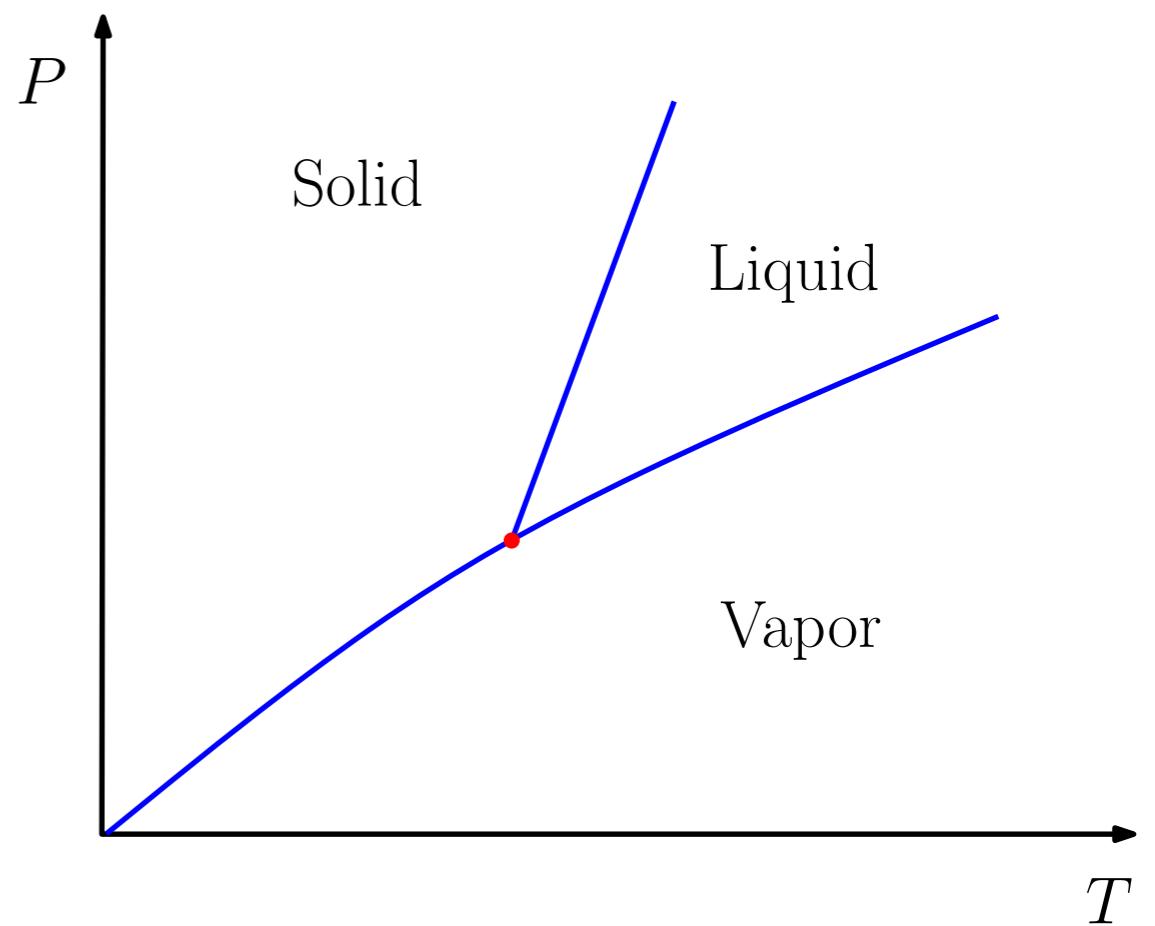
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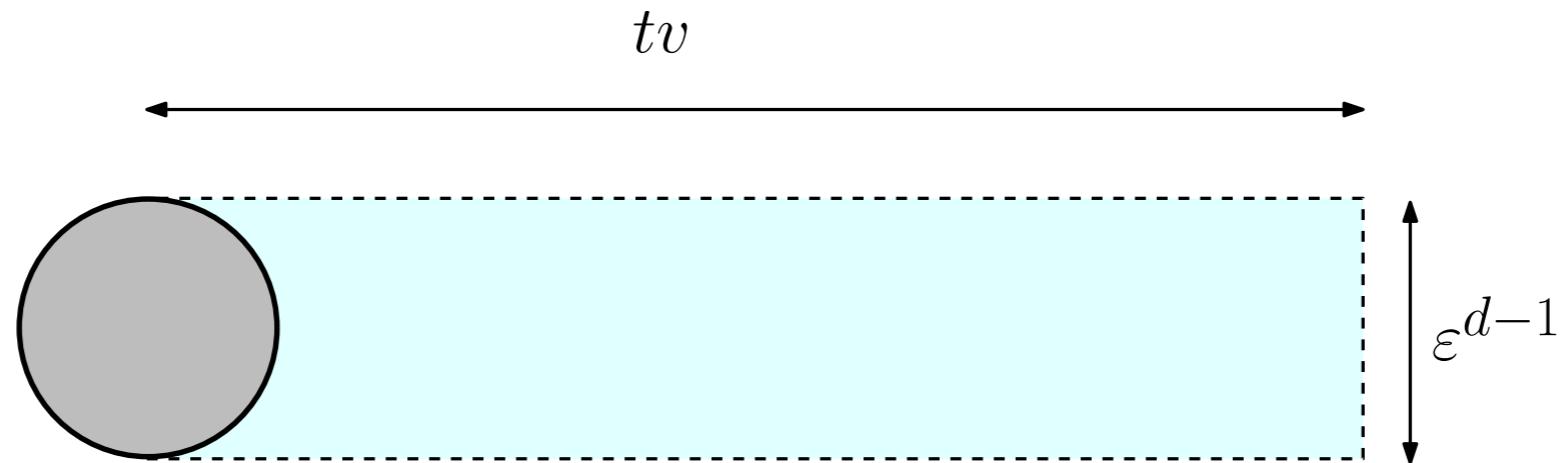
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Boltzmann-Grad scaling



- Volume covered by a particle $= tv\varepsilon^{d-1}$
- On average N particles per unit volume

On average, a particle has α collisions per unit of time

$$N \times \varepsilon^{d-1} \equiv \alpha$$

Diluted Gas of hard spheres

Gas of N hard spheres with deterministic Newtonian dynamics (elastic collisions).

Initial data at equilibrium and a tagged particle (x_1, v_1)

Questions.

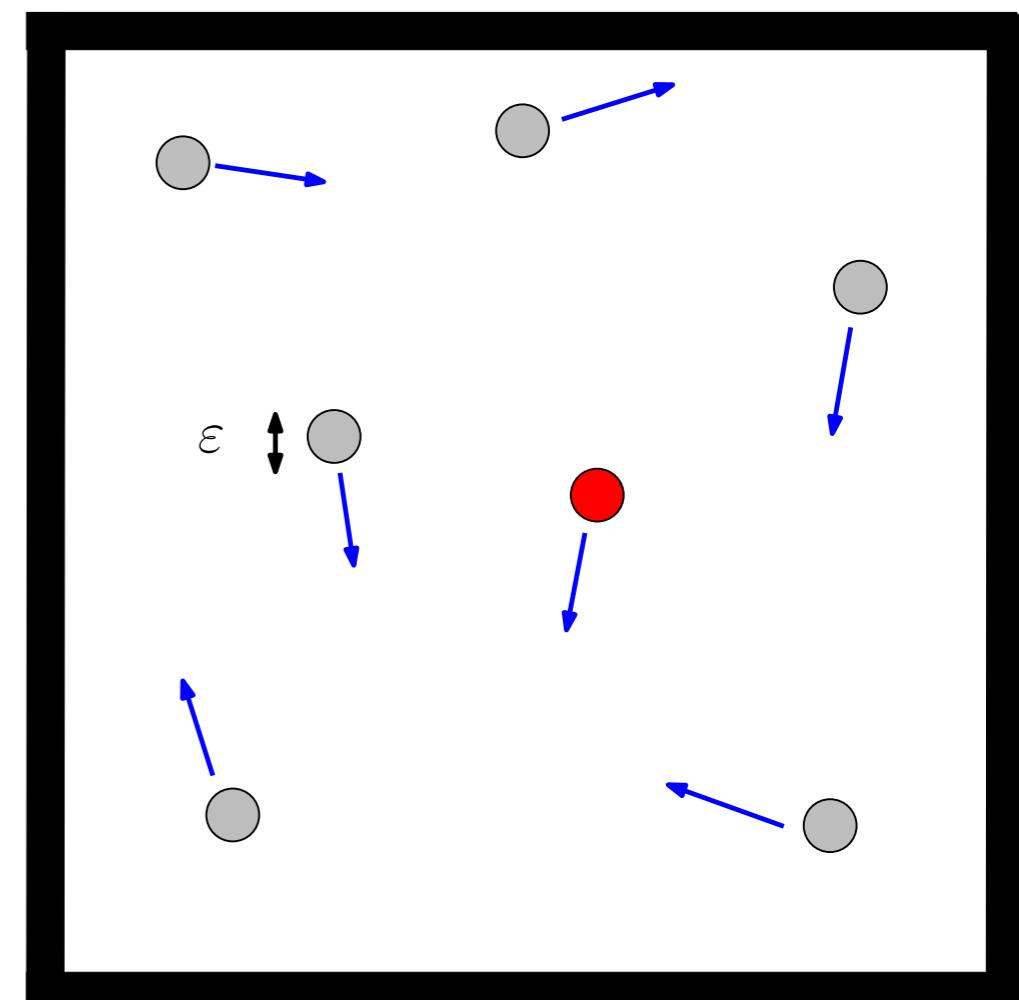
In the Boltzmann-Grad scaling

$$N \times \varepsilon^{d-1} \equiv \alpha \text{ and } N \rightarrow \infty$$

1. Distribution of $(x_1(t), v_1(t))$

2. Position of the tagged particle

$$x_1(\alpha t) \text{ when } \alpha \rightarrow \infty$$



Hard Sphere dynamics

Gas of N hard spheres : $Z_N = \{(x_i(t), v_i(t)\}_{i \leq N}$

$$\frac{dx_i}{dt} = v_i, \quad \frac{dv_i}{dt} = 0 \quad \text{as long as } |x_i(t) - x_j(t)| > \varepsilon,$$

and elastic collisions if $|x_i(t) - x_j(t)| = \varepsilon$

$$\begin{cases} v_i(t^+) = v_i(t^-) - \frac{1}{\varepsilon^2} (v_i - v_j) \cdot (x_i - x_j) (x_i - x_j)(t^-) \\ v_j(t^+) = v_j(t^-) + \frac{1}{\varepsilon^2} (v_i - v_j) \cdot (x_i - x_j) (x_i - x_j)(t^-) \end{cases}$$

Liouville equation for the particle density $f_N(t, Z_N)$

$$\partial_t f_N + \sum_{i=1}^N v_i \cdot \nabla_{x_i} f_N = 0$$

in the phase space

$$\mathcal{D}_\varepsilon^N := \{Z_N \in \mathbf{T}_\lambda^{dN} \times \mathbb{R}^{dN} / \forall i \neq j, \quad |x_i - x_j| > \varepsilon\}$$

with specular reflection on the boundary $\partial \mathcal{D}_\varepsilon^N$.

The tagged particle

Equilibrium distribution

$$M_{N,\beta}(Z_N) = \frac{1}{\mathcal{Z}_{N,\beta}} \exp \left(-\frac{\beta}{2} \sum_{i=1}^N |\nu_i|^2 \right) \prod_{i \neq j} 1_{|x_i - x_j| > \varepsilon}$$

Particle $Z_1 = (x_1, \nu_1)$ is tagged. Initial distribution :

$$f_N^0(Z_N) = M_{N,\beta}(Z_N) \rho^0(x_1)$$

Uniform bound: $\rho^0(x_1) \leq \mu$

Notation: Marginals

$$t \geq 0, \forall s \geq 1, \quad f_N^{(s)}(t, Z_s) = \iint \dots \int f_N(t, Z_N) dz_{s+1} \dots dz_N$$

Tagged particle distribution $f_N^{(1)}(t, (x_1, \nu_1))$

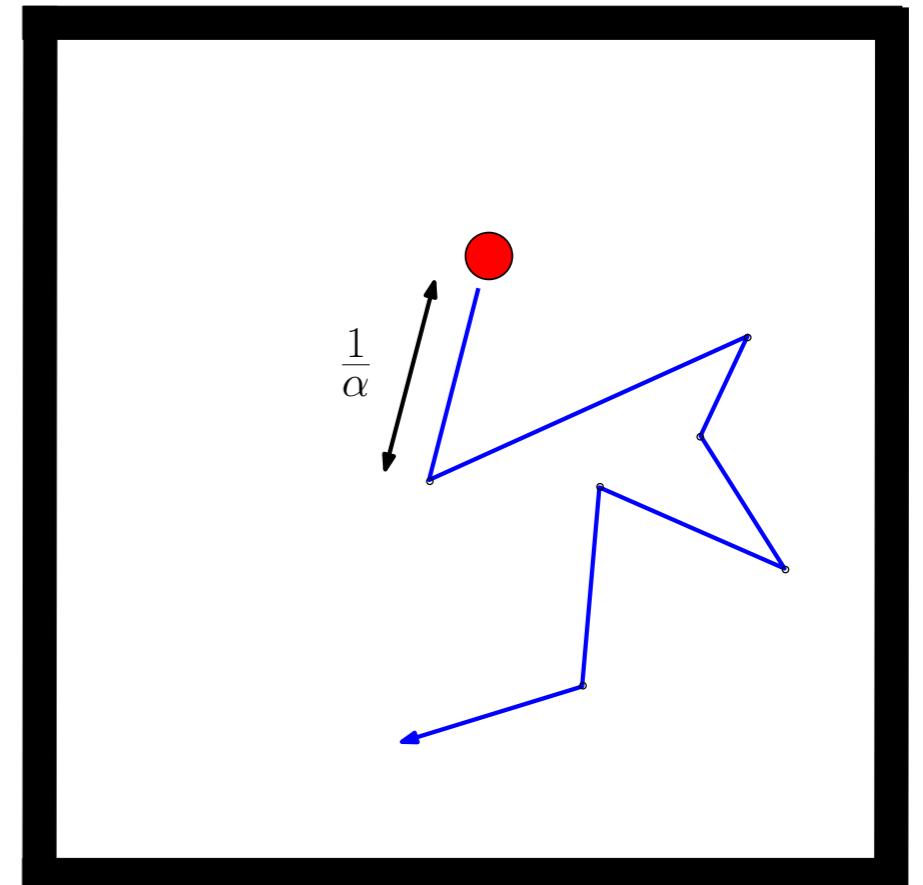
Limiting stochastic process

single particle dynamics

Position : $x(t) = \int_0^t v(u)du$

Markov process on the velocities

$\{v(t)\}_{t \geq 0}$ with generator αL



$$Lg(v) := \iint [g(v) - g(v')] \left((v - v_1) \cdot \nu \right)_+ M_\beta(v_1) dv_1 d\nu$$

$$v' = v + (\nu \cdot (v_1 - v)) \nu, \quad v'_1 = v_1 - (\nu \cdot (v_1 - v)) \nu$$

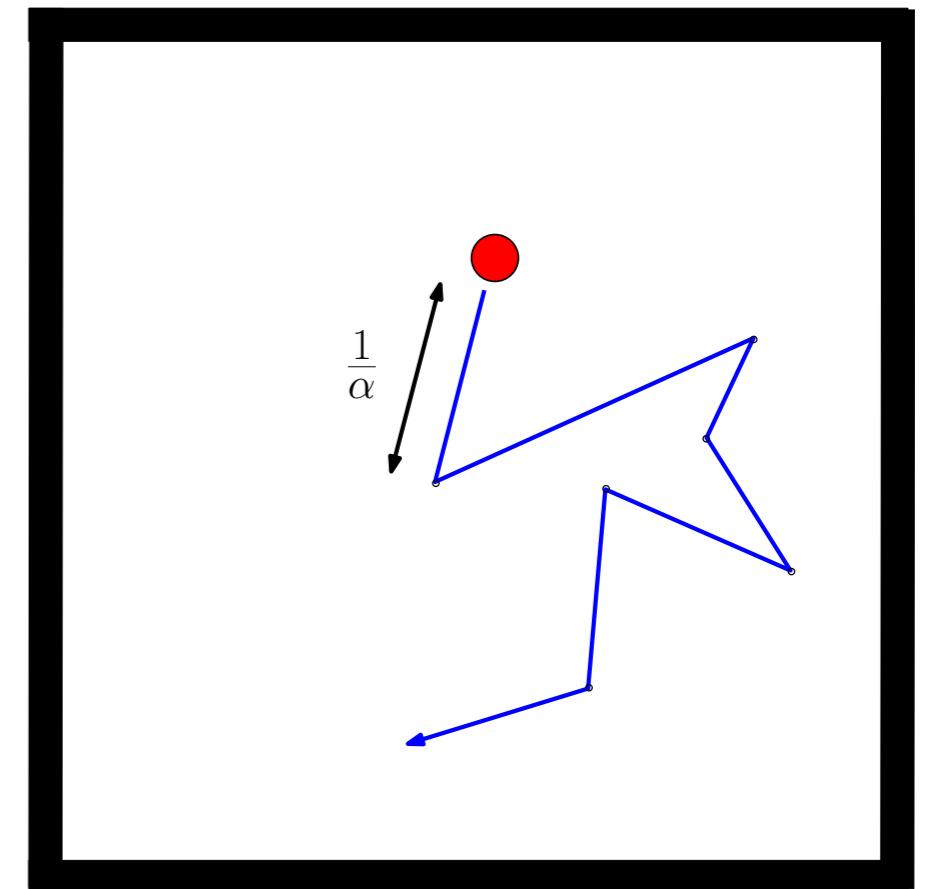
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Particle distribution $M_\beta(v)\varphi_{\color{red}\alpha}(x, v, t)$ follows the

Linear Boltzmann equation

$$\partial_t \varphi + v \cdot \nabla_x \varphi = -\alpha L \varphi$$

Probabilist approaches :

Tanaka, Sznitman, Méléard, Graham, Fournier ...

[van Beijeren, Lanford, Lebowitz, Spohn]

N particle
system

$$f_N^{(1)}(\textcolor{red}{x}_1, \textcolor{red}{v}_1, t)$$

$$\alpha = N \varepsilon^{d-1}$$

$$N \rightarrow \infty$$



$$t > 0$$

Linear Boltzmann
equation

$$\varphi_\alpha(\textcolor{red}{x}_1, \textcolor{red}{v}_1, t) M_\beta(\textcolor{red}{v}_1)$$

[van Beijeren, Lanford, Lebowitz, Spohn]

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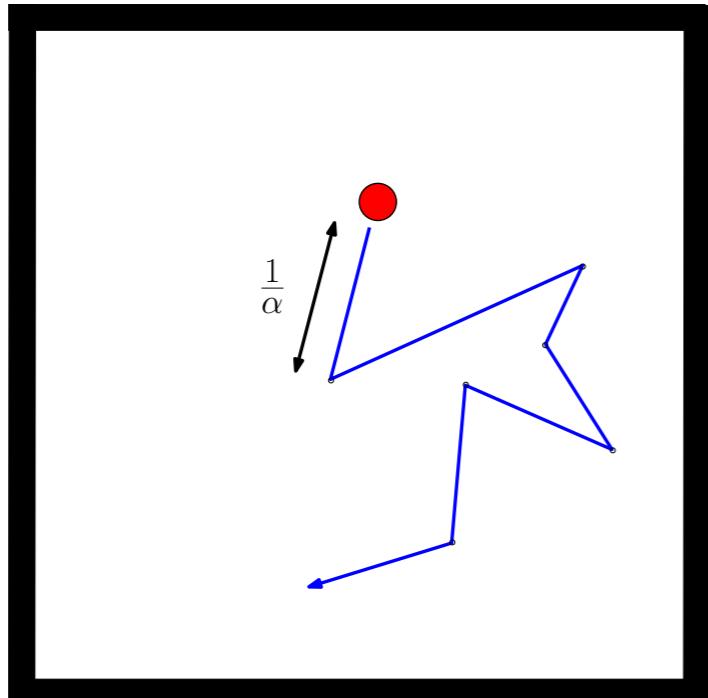
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Large time asymptotic

$$\begin{aligned} t &= \alpha \tau \\ \alpha &\rightarrow \infty \end{aligned}$$

Heat equation

Brownian motion

[van Beijeren, Lanford, Lebowitz, Spohn]

N particle system

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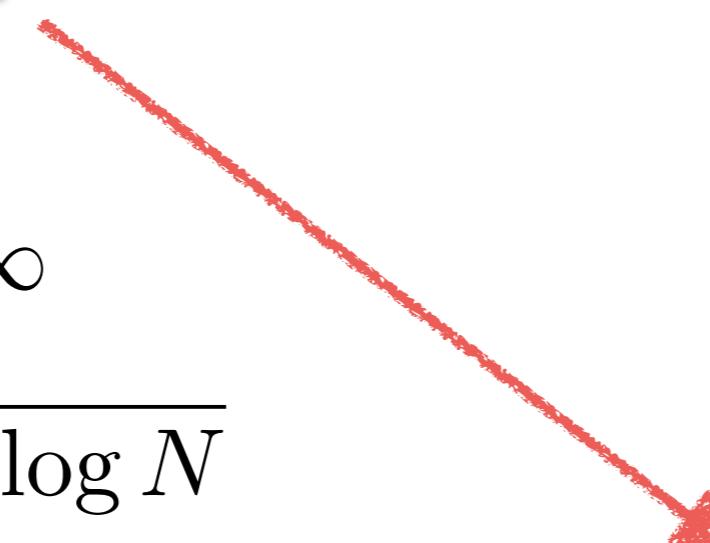
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Linear Boltzmann equation

$$\varphi_\alpha(\mathbf{x}_1, \mathbf{v}_1, t) M_\beta(\mathbf{v}_1)$$

$$N \rightarrow \infty$$

$$\alpha = \sqrt{\log \log N}$$



$$\begin{aligned} t &= \alpha \tau \\ \alpha &\rightarrow \infty \end{aligned}$$

Heat equation

Brownian motion

Large time asymptotic

Convergence to the Brownian motion

Rescaled position of the tagged particle

$$\chi(\tau) = x_1(\alpha\tau) \quad \text{with} \quad \alpha = \sqrt{\log \log N}$$

Initial data $f_N^0(Z_N) = M_{N,\beta}(Z_N) \rho^0(x_1)$

Theorem [B., Gallagher, Saint-Raymond]

χ converges weakly to a brownian motion with variance κ_β

The distribution of the tagged particle $f_N^{(1)}(x_1, v_1, \alpha\tau)$

converges as $N \rightarrow \infty$ to $M_\beta(v_1) \rho(x_1, \tau)$

$$\partial_\tau \rho = \kappa_\beta \Delta_x \rho \quad \text{on } \mathbb{R}^+ \times [0, 1]^d, \quad \rho|_{\tau=0} = \rho^0$$

Quantum brownian motion: [Erdös, Salmhofer, Yau]
Lorentz gas

Boltzmann equation

Theorem.

For chaotic initial data $f_N^0(Z_N) \simeq \prod_{i=1}^N f^0(z_i)$ the density of the particle system converges up to a time $t > 0$ to the solution of the Boltzmann equation when $N \rightarrow \infty$, $N\varepsilon^{d-1} = \alpha$

$$\begin{aligned} & \partial_t f + v \cdot \nabla_x f \\ &= \iint_{\mathbf{S}^{d-1} \times \mathbb{R}^d} [f(v')f(v'_1) - f(v)f(v_1)] ((v - v_1) \cdot \nu)_+ dv_1 d\nu \end{aligned}$$

$$\text{with } v' = v + \nu \cdot (v_1 - v) \nu, \quad v'_1 = v_1 - \nu \cdot (v_1 - v) \nu$$

[**Lanford**], [King], [Alexander], [Uchiyama], [Cercignani, Illner, Pulvirenti], [Simonella], [*Gallagher, Saint-Raymond, Texier*], [Pulvirenti, Saffirio, Simonella]

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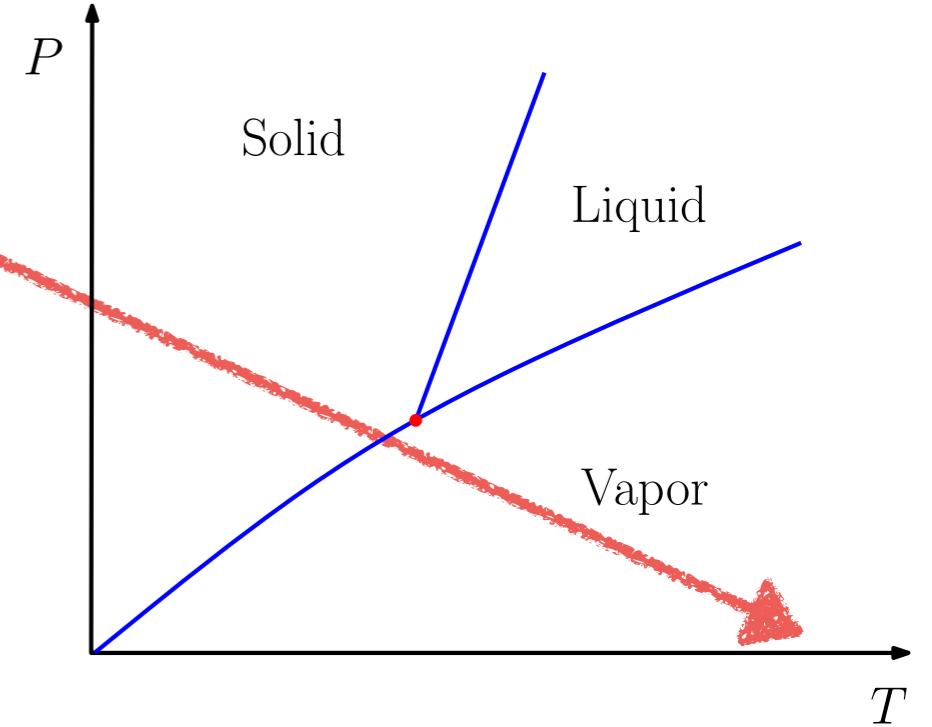
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Lanford's strategy leads to a short time convergence which depends on f^0 . The convergence time remains short even if initially the system starts from equilibrium !!!

Equilibrium Properties

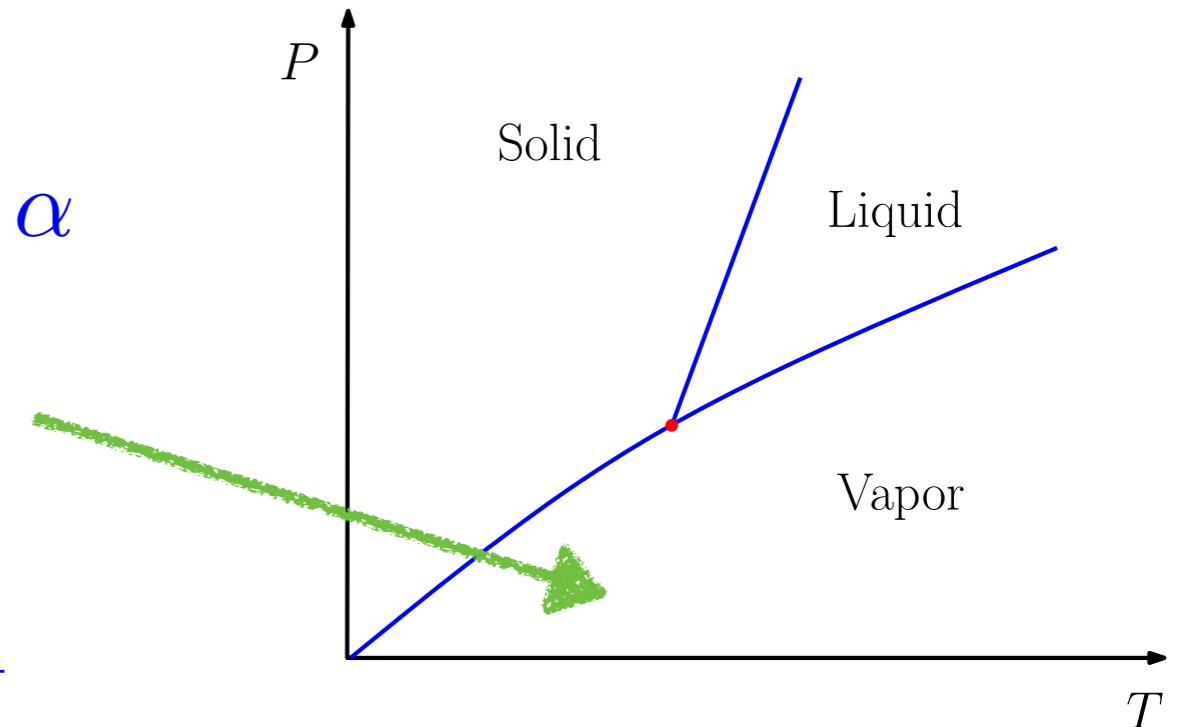
- Boltzmann-Grad : $N \times \varepsilon^{d-1} \equiv \alpha$
- Low density : $N \times \varepsilon^d = \rho \ll 1$
- High density : $N \times \varepsilon^d = \rho \gg 1$



Open Problem : Vapor/Solid phase transition.

Equilibrium Properties

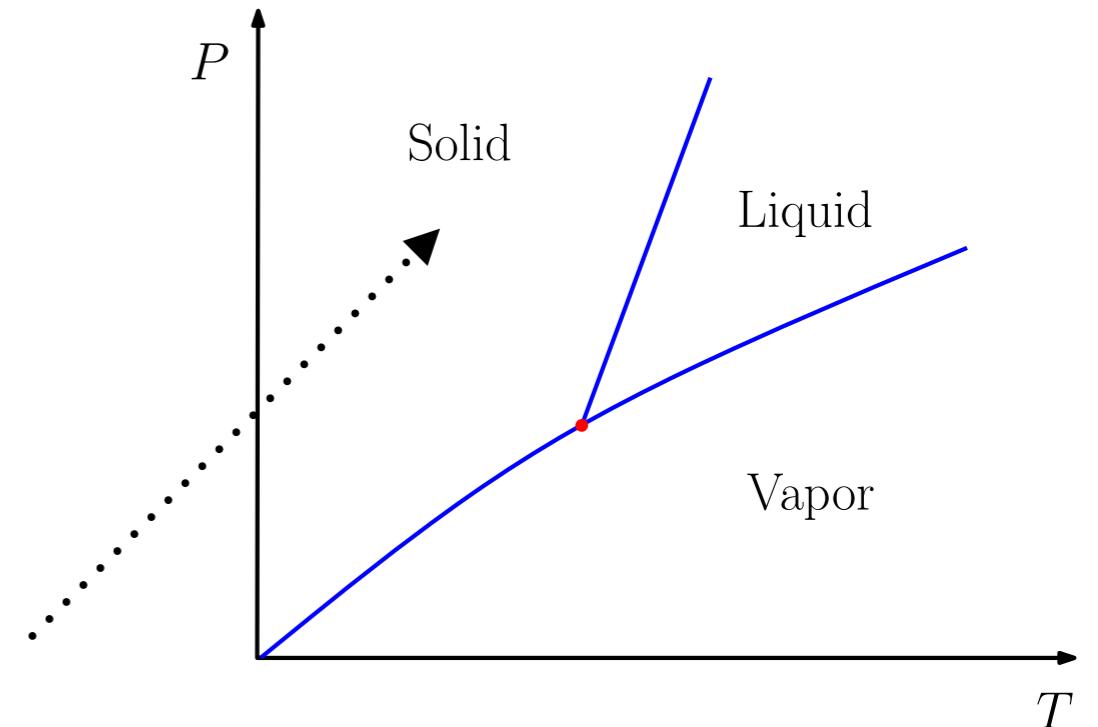
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Liquid/Vapor phase transition

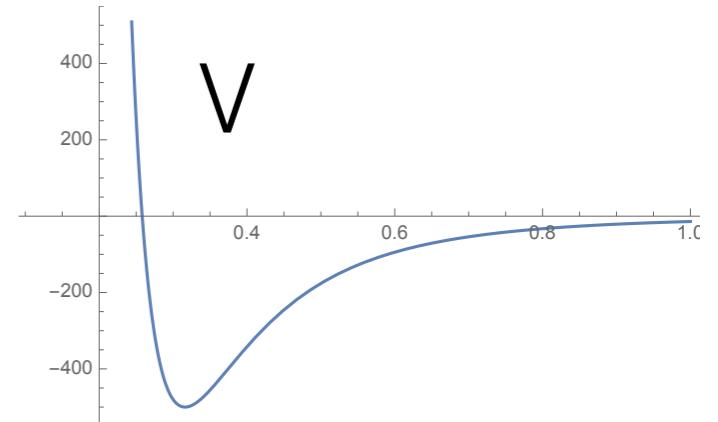
Interaction potential

Lennard-Jones potential

$$H(X) = - \sum_{i,j} V\left(\frac{x_i - x_j}{\varepsilon}\right)$$

Gibbs measure $X = \{x_i \in [0, 1]^d; i \leq N\}$

$$\mu_{N,T}(X) = \frac{1}{Z_{N,T}} \exp\left(-\frac{1}{T} H(X)\right)$$



Liquid/Vapor phase transition

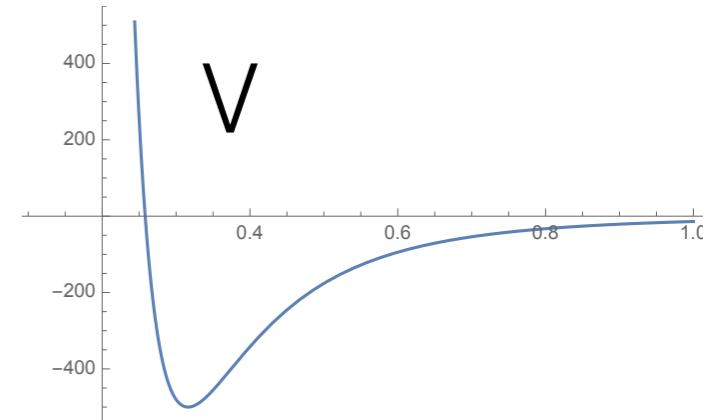
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There is a density $\rho = N\varepsilon^d$ and a temperature T

$$\mu_{N,T} \left(\text{Phase transition} \right) \xrightarrow{N \rightarrow \infty} 1$$

[Lebowitz, Mazel, Presutti]

Continuum percolation

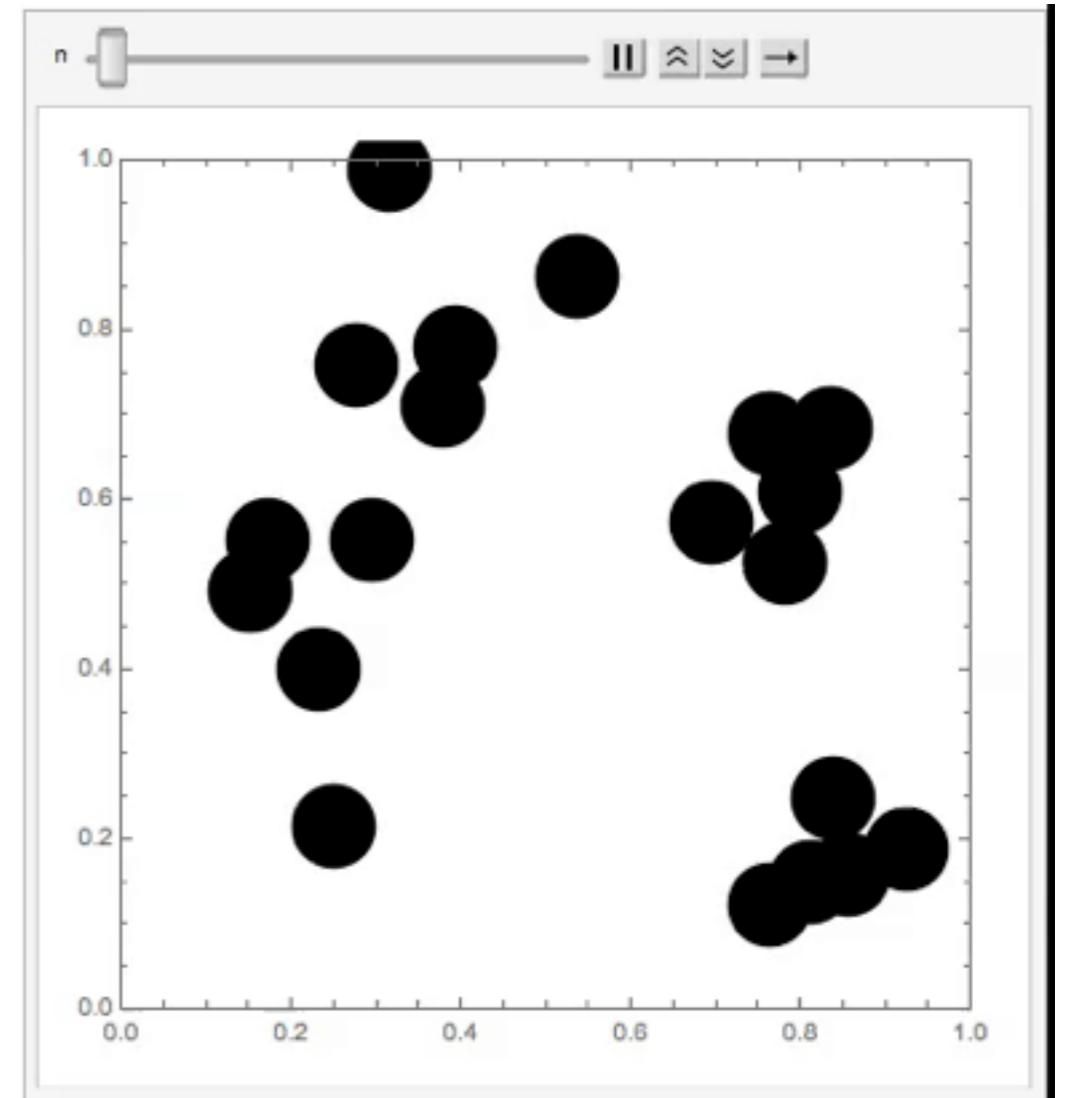
Disks of radius ε randomly distributed

No interaction, No exclusion.

density $\rho = N\varepsilon^d$

Phase transition : $N \rightarrow \infty$

- $\rho < \rho_c$: small percolation clusters
- $\rho > \rho_c$: percolation of a macroscopic cluster



Many more models : R. Marchand

Continuum percolation

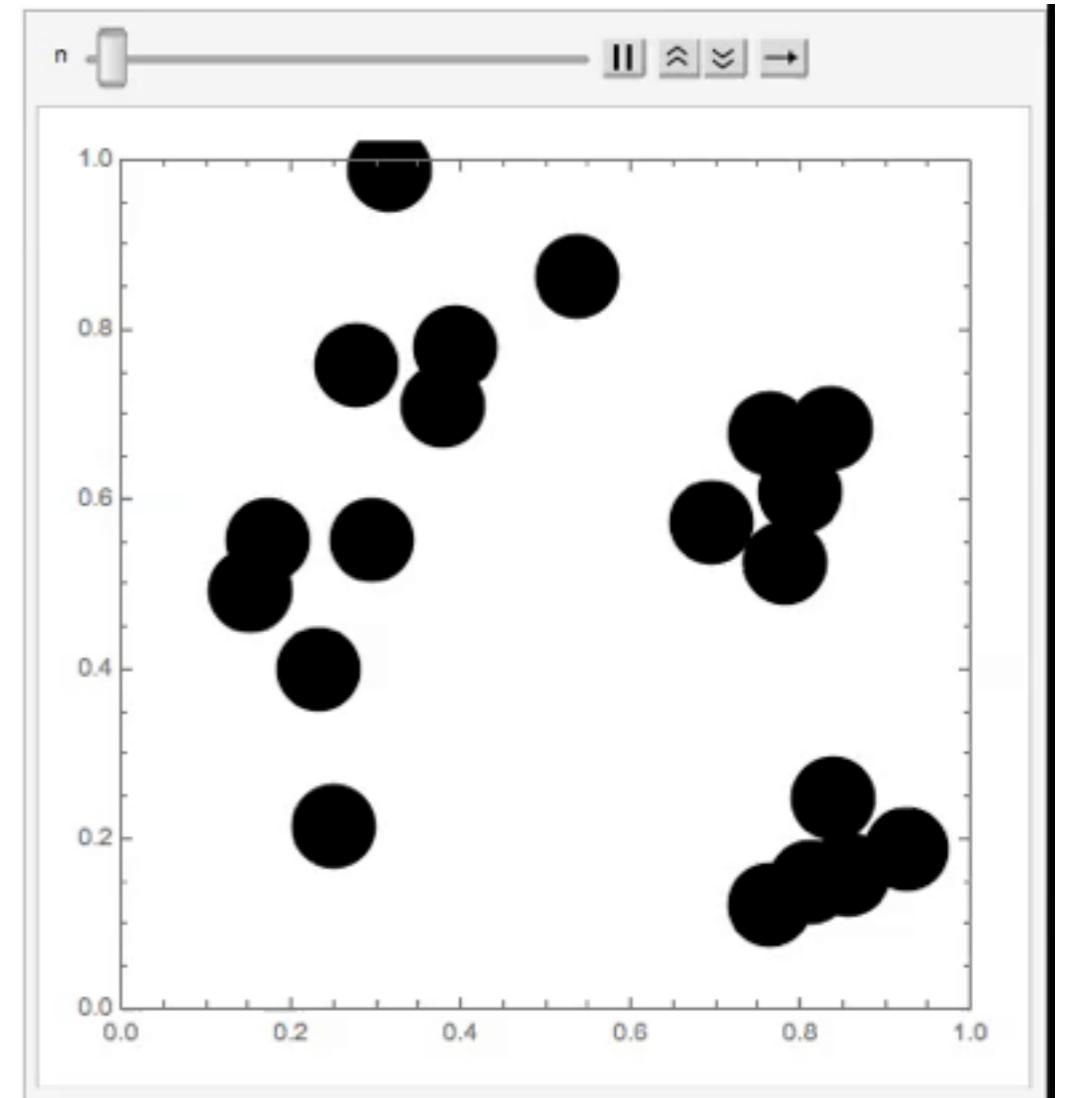
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Lattice gas models

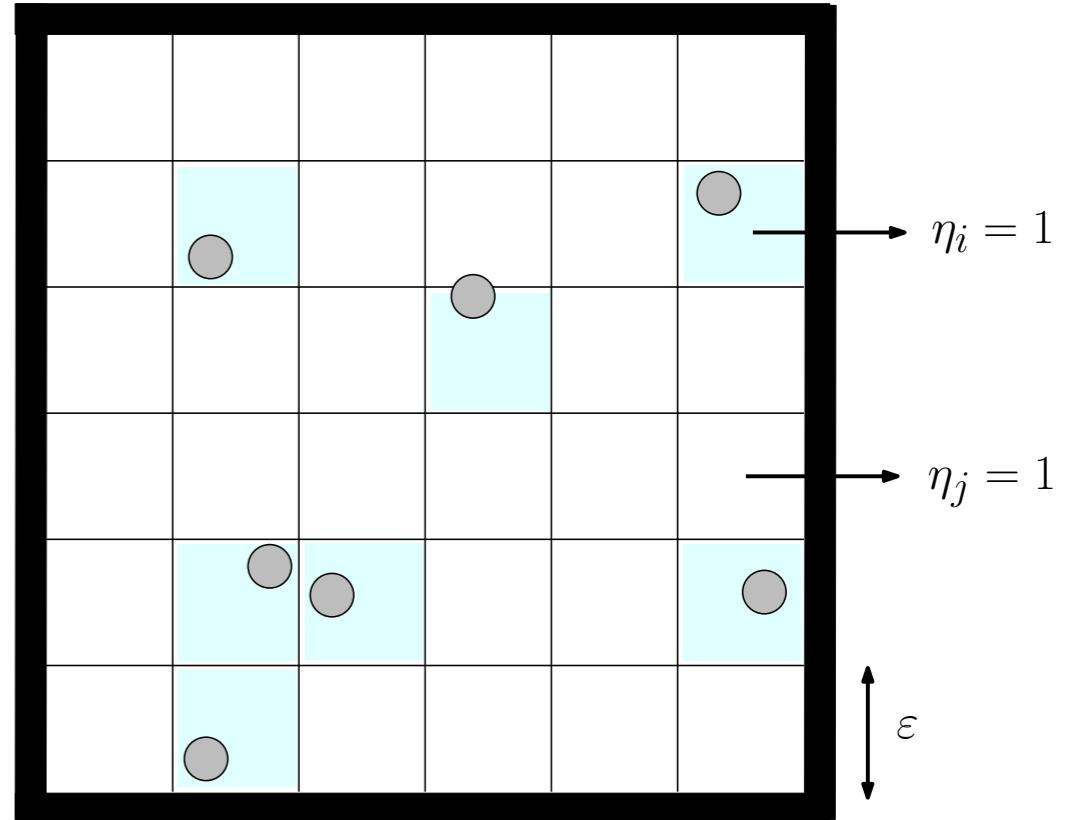
Simplifying the model:

- discrete variables

$$\eta_i \in \{0, 1\} \text{ with } i \in \left\{1, \dots, \frac{1}{\varepsilon}\right\}^d$$

- nearest neighbor interaction

$$H(\eta) = - \sum_{i \sim j} \eta_i \eta_j - h \sum_i \eta_i$$



Gibbs measure

$$\mu_{N,T}(\eta) = \frac{1}{Z_{N,T}} \exp\left(-\frac{1}{T} H(\eta)\right)$$

Lattice gas models

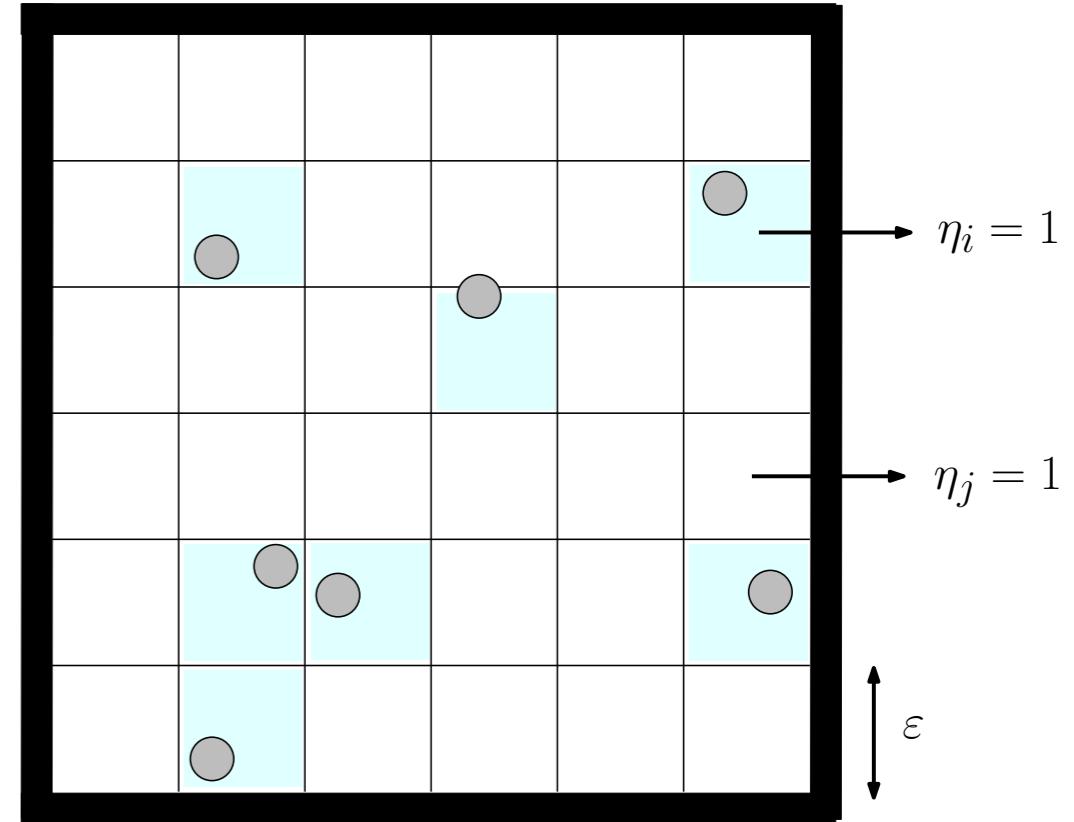
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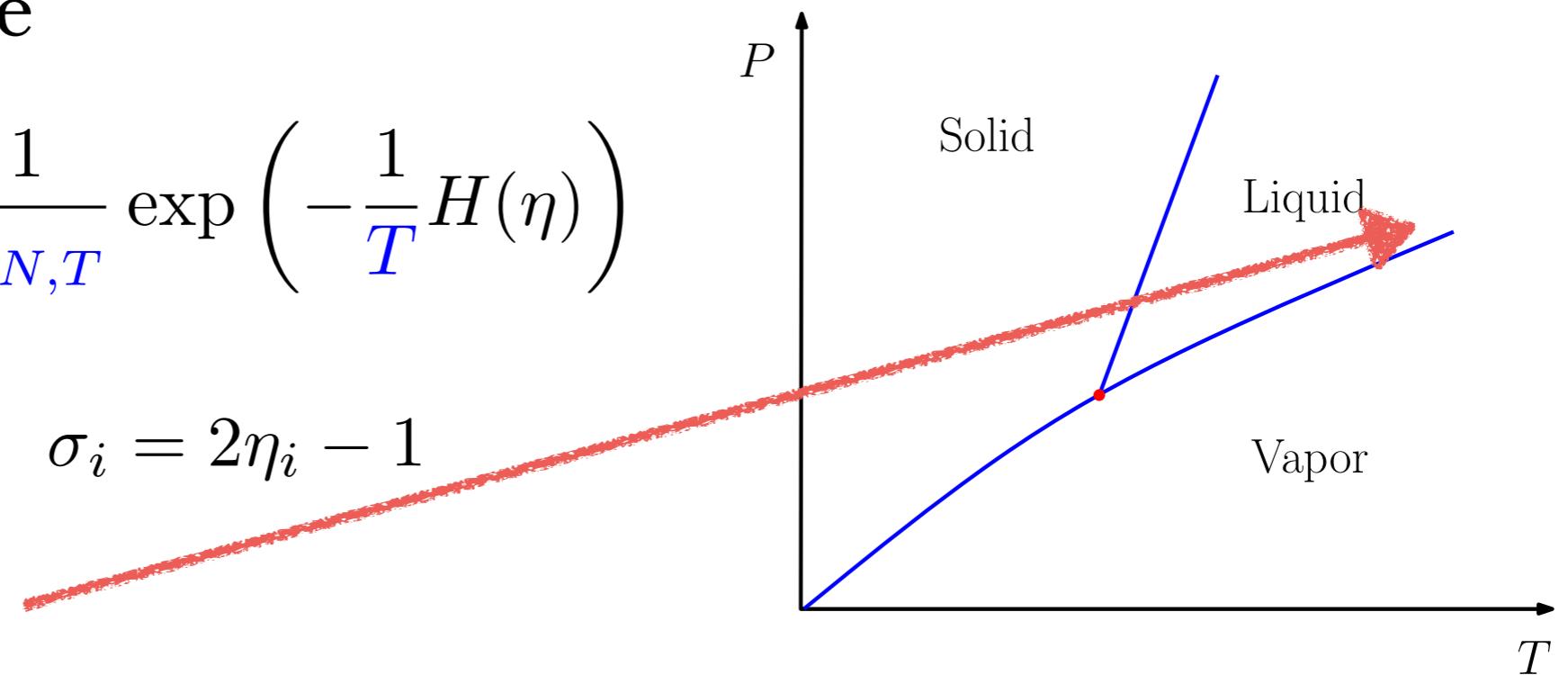


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Ising Model : $\sigma_i = 2\eta_i - 1$

B. de Tilière



- **Equilibrium systems** : Gibbs measure

Parameters : (T, h) ; (T, P)

Critical curve : $T = F(h)$

- **Non-equilibrium systems** : Local dynamical rules
 - Sandpile models; Driven particle systems ...
 - Long range correlations; Self-organized criticality (in general no obvious Hamiltonian)

M. Gorny : A mean field approach