

# VANISHING CYCLES

## AND CARTAN EIGENVECTORS

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### Introduction

**0.1.** Let  $A(R)$  be the Cartan matrix of a finite root system  $R$ . The coordinates of its eigenvectors have an important meaning in the physics of integrable systems: namely, these numbers appear as the masses of particles (or, dually, as the energy of solitons) in affine Toda field theories, cf. [F], [D].

Historically, the first example of the system of type  $E_8$  appeared in the pioneering papers [Z] on the 2D critical Ising model in a magnetic field.

The aim of this note is a study of these numbers, and their  $q$ -deformations, using the motivation coming from the singularity theory.

Let us suppose that  $R$  is simply laced, i.e. of type  $A, D$ , or  $E$ . These root systems are in one-to-one correspondence with (classes of) simple singularities  $f : \mathbb{C}^N \rightarrow \mathbb{C}$ , cf. [AGV]. Under this correspondence, the root lattice  $Q(R)$  is identified with the lattice of vanishing cycles, and the Cartan matrix  $A(R)$  is the intersection matrix with respect to a *distinguished base*. The action of the Weyl group on  $Q(R)$  is realized by Gauss - Manin monodromies - this is the Picard - Lefschetz theory (for some details see §1 below).

Remarkably, this geometric picture provides a finer structure: namely, the symmetric matrix  $A = A(R)$  comes equipped with a decomposition

$$A = L + L^t \tag{0.1}$$

where  $L$  is a nondegenerate triangular "Seifert form", or "variation matrix". The matrix

$$C = -L^{-1}L^t \tag{0.2}$$

represents a Coxeter element of  $R$ ; geometrically it is the operator of "classical monodromy".

We call the relation (0.1) - (0.2) between the Cartan matrix and the Coxeter element the *Cartan/Coxeter correspondence*.

Incidentally, in a particular case (corresponding to a bipartition of the Dynkin graph) this relation is equivalent to an observation by R. Steinberg, cf. [Stein], cf. 2.4 below. It enables one to relate the eigenvectors of  $A$  and  $C$ , cf. Theorem 2.5.

A decomposition (0.1) will be called a *polarization* of the Cartan matrix  $A$ . In 2.3 below we introduce an operation of *Sebastiani - Thom*, or *joint product*  $A * B$  of Cartan matrices (or of polarized lattices)  $A$  and  $B$ . With respect to this operation the Coxeter eigenvectors factorize very simply.

In this note we will mainly concentrate on the example of  $E_8$ ; this lattice decomposes into three "quarks":

$$E_8 = A_4 * A_2 * A_1 \tag{0.3}$$

This decomposition is the main message from the singularity theory, and we discuss it in detail in this note.

We use (0.3) and the Cartan/Coxeter correspondence to obtain some expressions for all Cartan eigenvectors of  $E_8$ ; this is the first main result of this note, see 3.9 below.

(An elegant expression for all the Cartan eigenvectors of all finite root systems was obtained by P. Dorey, cf. [D] (a), Table 2 on p. 659.)

**0.2.** In the paper [Giv] A. Givental has proposed a  $q$ -twisted version of the Picard - Lefschetz theory, which gave rise to a  $q$ -deformation of  $A$ ,

$$A(q) = L + qL^t.$$

In the last section, §4, we calculate the eigenvalues and eigenvectors of  $A(q)$  in terms of the eigenvalues and eigenvectors of  $A$ . This is the second main result of this note.

It turns out that if  $\lambda$  is an eigenvalue of  $A$  then

$$\lambda(q) = 1 + (\lambda - 2)\sqrt{q} + q \tag{0.4}$$

will be an eigenvalue of  $A(q)$ . The coordinates of the corresponding eigenvector  $v(q)$  are obtained from the coordinates of  $v = v(1)$  by multiplication by appropriate powers of  $q$ ; this is related to the fact that the Dynkin graph of  $A$  is a tree, cf. 4.9.

For an example of  $E_8$ , see (4.5.1).

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## §1. Recollections from singularity theory

Here we recall some classical constructions and statements, cf. [AGV].

This section plays a motivational role: formally the results of the following sections are purely algebraic, and do not depend on it, however, one could not arrive at them without this geometric motivation.

**1.1.** Let  $f : (\mathbb{C}^N, 0) \rightarrow (\mathbb{C}, 0)$  be the germ of a holomorphic function with an isolated critical point at 0, with  $f(0) = 0$ . We will be interested only in polynomial functions (from the list below, cf. 1.4), so  $f \in \mathbb{C}[x_1, \dots, x_N]$ . The *Milnor ring* of  $f$  is defined by

$$\text{Miln}(f, 0) = \mathbb{C}[[x_1, \dots, x_N]] / (\partial_1 f, \dots, \partial_N f)$$

where  $\partial_i := \partial/\partial x_i$ ; it is a finite-dimensional commutative  $\mathbb{C}$ -algebra. (In fact, it is a Frobenius, or, equivalently, a Gorenstein algebra.) The number

$$\mu := \dim_{\mathbb{C}} \text{Miln}(f, 0)$$

is called the multiplicity or Milnor number of  $(f, 0)$ .

A *Milnor fiber* is

$$V_z = f^{-1}(z) \cap \bar{B}_\rho$$

where

$$\bar{B}_\rho = \{(x_1, \dots, x_N) \mid \sum |x_i|^2 \leq \rho\}$$

for  $1 \gg \rho \gg |z| > 0$ .

For  $z$  belonging to a small disc  $D_\epsilon = \{z \in \mathbb{C} \mid |z| < \epsilon\}$ , the space  $V_z$  is a complex manifold with boundary, homotopically equivalent to a bouquet  $\vee S^{N-1}$  of  $\mu$  spheres, [M].

The family of free abelian groups

$$Q(f; z) := \tilde{H}_{N-1}(V_z; \mathbb{Z}) \cong \mathbb{Z}^\mu, \quad z \in \dot{D}_\epsilon := D_\epsilon \setminus \{0\}, \quad (1.1.1)$$

( $\tilde{H}$  means that we take the reduced homology for  $N = 1$ ), carries a flat Gauss - Manin connection.

Take  $t \in \mathbb{R}_{>0} \cap \dot{D}_\epsilon$ ; the lattice  $Q(f; t)$  does not depend, up to a canonical isomorphism, on the choice of  $t$ . Let us call this lattice  $Q(f)$ . The linear operator

$$T(f) : Q(f) \xrightarrow{\sim} Q(f) \quad (1.1.2)$$

induced by the path  $p(\theta) = e^{i\theta}t$ ,  $0 \leq \theta \leq 2\pi$ , is called the classical monodromy of the germ  $(f, 0)$ .

In all the examples below  $T(f)$  has finite order  $h$ . The eigenvalues of  $T(f)$  have the form  $e^{2\pi ik/h}$ ,  $k \in \mathbb{Z}$ . The set of suitable chosen  $k$ 's for each eigenvalue are called the *spectrum* of our singularity.

**1.2. Morse deformations.** The  $\mathbb{C}$ -vector space  $\text{Miln}(f, 0)$  may be identified with the tangent space to the base  $B$  of the miniversal deformation of  $f$ . For

$$\lambda \in B^0 = B \setminus \Delta$$

where  $\Delta \subset B$  is an analytic subset of codimension 1, the corresponding function  $f_\lambda : \mathbb{C}^N \rightarrow \mathbb{C}$  has  $\mu$  nondegenerate Morse critical points with distinct critical values, and the algebra  $\text{Miln}(f_\lambda)$  is semisimple, isomorphic to  $\mathbb{C}^\mu$ .

Let  $0 \in B$  denote the point corresponding to  $f$  itself, so that  $f = f_0$ , and pick  $t \in \mathbb{R}_{>0} \cap \overset{\bullet}{D}_\epsilon$  as in 1.1.

Afterwards pick  $\lambda \in B^0$  close to 0 in such a way that the critical values  $z_1, \dots, z_\mu$  of  $f_\lambda$  have absolute values  $\ll t$ .

As in 1.1, for each

$$z \in \tilde{D}_\epsilon := D_\epsilon \setminus \{z_1, \dots, z_\mu\}$$

the Milnor fiber  $V_z$  has the homotopy type of a bouquet  $\vee S^{N-1}$  of  $\mu$  spheres, and we will be interested in the middle homology

$$Q(f_\lambda; z) = \tilde{H}_{N-1}(V_z; \mathbb{Z}) \cong \mathbb{Z}^\mu$$

( $\tilde{H}$  means that we take the reduced homology for  $N = 1$ ).

The lattices  $Q(f_\lambda; z)$  carry a natural bilinear product induced by the cup product in the homology which is symmetric (resp. skew-symmetric) when  $N$  is odd (resp. even).

The collection of these lattices, when  $z \in \tilde{D}_\epsilon$  varies, carries a flat Gauss - Manin connection.

Consider an "octopus"

$$\text{Oct}(t) \subset \mathbb{C}$$

with the head at  $t$ : a collection of non-intersecting paths  $p_i$  ("tentacles") connecting  $t$  with  $z_i$  and not meeting the critical values  $z_j$  otherwise. It gives rise to a base

$$\{b_1, \dots, b_\mu\} \subset Q(f_\lambda) := Q(f_\lambda; t)$$

(called "distinguished") where  $b_i$  is the cycle vanishing when being transferred from  $t$  to  $z_i$  along the tentacle  $p_i$ , cf. [Gab], [AGV].

The Picard - Lefschetz formula describe the action of the fundamental group  $\pi_1(\tilde{D}_\epsilon; t)$  on  $Q(f_\lambda)$  with respect to this basis. Namely, consider a loop  $\gamma_i$  which turns around  $z_i$  along the tentacle  $p_i$ , then the corresponding transformation of

$Q(f_\lambda)$  is the reflection (or transvection)  $s_i := s_{b_i}$ , cf. [Lef], Théorème fondamental, Ch. II, p. 23.

The loops  $\gamma_i$  generate the fundamental group  $\pi_1(\tilde{D}_\epsilon)$ . Let

$$\rho : \pi_1(\tilde{D}_\epsilon; t) \longrightarrow GL(Q(f_\lambda))$$

denote the monodromy representation. The image of  $\rho$ , denoted by  $G(f_\lambda)$  and called the *monodromy group of  $f_\lambda$* , lies inside the subgroup  $O(Q(f_\lambda)) \subset GL(Q(f_\lambda))$  of linear transformations respecting the above mentioned bilinear form on  $Q(f_\lambda)$ .

The subgroup  $G(f_\lambda)$  is generated by  $s_i, 1 \leq i \leq \mu$ .

As in 1.1, we have the monodromy operator

$$T(f_\lambda) \in G(f_\lambda),$$

the image by  $\rho$  of the path  $p \subset \tilde{D}_\epsilon$  starting at  $t$  and going around all points  $z_1, \dots, z_\mu$ .

This operator  $T(f_\lambda)$  is now a product of  $\mu$  simple reflections

$$T(f_\lambda) = s_1 s_2 \dots s_\mu,$$

- this is because the only critical value 0 of  $f$  became  $\mu$  critical values  $z_1, \dots, z_\mu$  of  $f_\lambda$ .

One can identify the relative (reduced) homology  $\tilde{H}_{N-1}(V_t, \partial V_t; \mathbb{Z})$  with the dual group  $\tilde{H}_{N-1}(V_t; \mathbb{Z})^*$ , and one defines a map

$$\text{var} : \tilde{H}_{N-1}(V_t, \partial V_t; \mathbb{Z}) \longrightarrow \tilde{H}_{N-1}(V_t; \mathbb{Z}),$$

called a *variation operator*, which translates to a map

$$L : Q(f_\lambda)^* \xrightarrow{\sim} Q(f_\lambda)$$

("Seifert form") such that the matrix  $A(f_\lambda)$  of the bilinear form in the distinguished basis is

$$A(f_\lambda) = L + (-1)^{N-1} L^t,$$

and

$$T(f_\lambda) = (-1)^{N-1} L L^{-t}.$$

SIGNS !!!

A choice of a path  $q$  in  $B$  connecting 0 with  $\lambda$ , enables one to identify  $Q(f)$  with  $Q(f_\lambda)$ , and  $T(f)$  will be identified with  $T(f_\lambda)$ .

The image  $G(f)$  of the monodromy group  $G(f_\lambda)$  in  $GL(Q(f)) \cong GL(Q(f_\lambda))$  is called the monodromy group of  $f$ ; it does not depend on a choice of a path  $q$ .

**1.3. Sebastiani - Thom factorization.** If  $g \in \mathbb{C}[y_1, \dots, y_M]$  is another function, the sum, or **join** of two singularities  $f \oplus g : \mathbb{C}^{N+M} \rightarrow \mathbb{C}$  is defined by

$$(f \oplus g)(x, y) = f(x) + g(y)$$

Obviously we can identify

$$\text{Miln}(f \oplus g) \cong \text{Miln}(f) \otimes \text{Miln}(g)$$

Note that the function  $g(y) = y^2$  is a unit for this operation.

It follows that the singularities  $f(x_1, \dots, x_N)$  and

$$f(x_1, \dots, x_N) + x_{M+1}^2 + \dots + x_{N+M}^2$$

are "almost the same". In order to have good signs (and for other purposes) it is convenient to add some squares to a given  $f$  to get  $N \equiv 3 \pmod{4}$ .

The fundamental Sebastiani - Thom theorem, [ST], says that there exists a natural isomorphism of lattices

$$Q(f \oplus g) \cong Q(f) \otimes_{\mathbb{Z}} Q(g),$$

and under this identification the full monodromy decomposes as

$$T_{f \oplus g} = T_f \otimes T_g$$

Thus, if

$$\text{Spec}(T_f) = \{e^{\mu_p \cdot 2\pi i / h_1}\}, \quad \text{Spec}(T_g) = \{e^{\nu_q \cdot 2\pi i / h_2}\}$$

then

$$\text{Spec}(T_{f \oplus g}) = \{e^{(\mu_p h_2 + \nu_q h_1) \cdot 2\pi i / h_1 h_2}\}$$

**1.4. Simple singularities.** Cf. [AGV] (a), 15.1. They are:

$$x^{n+1}, \quad n \geq 1, \quad (A_n)$$

$$x^2 y + y^{n-1}, \quad n \geq 4 \quad (D_n)$$

$$x^4 + y^3 \quad (E_6)$$

$$x y^3 + x^3 \quad (E_7)$$

$$x^5 + y^3 \quad (E_8)$$

Their names come from the following facts:

— their lattices of vanishing cycles may be identified with the corresponding root lattices;

— the monodromy group is identified with the corresponding Weyl group;

— the classical monodromy  $T_f$  is a Coxeter element, therefore its order  $h$  is equal to the Coxeter number, and

$$\text{Spec}(T_f) = \{e^{2\pi i k_1 / h}, \dots, e^{2\pi i k_r / h}\}$$

where the integers

$$1 = k_1 < k_2 < \dots < k_r = h - 1,$$

are the exponents of our root system.

We will discuss the case of  $E_8$  in some details below.

## §2. Cartan - Coxeter correspondence and join product

### 2.1. Lattices, polarization, Coxeter elements.

Let us call a *lattice* a pair  $(Q, A)$  where  $Q$  is a free abelian group, and

$$A : Q \times Q \longrightarrow \mathbb{Z}$$

a symmetric bilinear map ("Cartan matrix"). We shall identify  $A$  with a map

$$A : Q \longrightarrow Q^\vee := \text{Hom}(Q, \mathbb{Z}).$$

A *polarized lattice* is a triple  $(Q, A, L)$  where  $(Q, A)$  is a lattice, and

$$L : Q \xrightarrow{\sim} Q^\vee$$

("variation", or "Seifert matrix") is an isomorphism such that

$$A = A(L) := L + L^\vee \tag{2.1.1}$$

where

$$L^\vee : Q = Q^{\vee\vee} \xrightarrow{\sim} Q^\vee$$

is the conjugate to  $L$ .

The *Coxeter automorphism* of a polarized lattice is defined by

$$C = C(L) = -L^{-1}L^\vee \in GL(Q). \tag{2.1.2}$$

We shall say that the operators  $A$  and  $C$  are in a *Cartan - Coxeter correspondence*.

**2.1.1. Example.** Let  $(Q, A)$  be a lattice, and  $\{e_1, \dots, e_n\}$  an ordered  $\mathbb{Z}$ -base of  $Q$ . With respect to this base  $A$  is expressed as a symmetric matrix  $A = (a_{ij}) = A(e_i, e_j) \in \mathfrak{gl}_n(\mathbb{Z})$ . Let us suppose that all  $a_{ii}$  are even. We define the matrix of  $L$  to be the unique upper triangular matrix  $(\ell_{ij})$  such that  $A = L + L^t$  (in particular  $\ell_{ii} = a_{ii}/2$ ; in our examples we will have  $a_{ii} = 2$ .) We will call  $L$  the *standard polarization* associated to an ordered base.  $\square$

Polarized lattices form a groupoid:

an isomorphism of polarized lattices  $f : (Q_1, A_1, L_1) \xrightarrow{\sim} (Q_2, A_2, L_2)$  is by definition an isomorphism of abelian groups  $f : Q_1 \xrightarrow{\sim} Q_2$  such that

$$L_1(x, y) = L_2(f(x), f(y))$$

(and whence  $A_1(x, y) = A_2(f(x), f(y))$ ).

**2.2. Lemma.** (i) (orthogonality)

$$A(x, y) = A(Cx, Cy).$$

(ii) (gauge transformations) For any  $P \in GL(Q)$

$$A(P^\vee LP) = P^\vee A(L)P, \quad C(P^\vee LP) = P^{-1}C(L)P.$$

□

**2.3. Join product.** Suppose we are given two polarized lattices  $(Q_i, A_i, L_i)$ ,  $i = 1, 2$ .

Set  $Q = Q_1 \otimes Q_2$ , whence

$$L := L_1 \otimes L_2 : Q \xrightarrow{\sim} Q^\vee,$$

and define

$$A := A_1 * A_2 := L + L^\vee : Q \xrightarrow{\sim} Q^\vee$$

The triple  $(Q, A, L)$  will be called the **join**, or **Sebastiani - Thom**, product of the polarized lattices  $Q_1$  and  $Q_2$ , and denoted by  $Q_1 * Q_2$ .

Obviously

$$C(L) = -C(L_1) \otimes C(L_2) \in GL(Q_1 \otimes Q_2).$$

It follows that if  $\text{Spec}(C(L_i)) = \{e^{2\pi i k_i/h_i}, k_i \in K_i\}$  then

$$\text{Spec}(C(L)) = \{-e^{2\pi i(k_1/h_1 + k_2/h_2)}, (k_1, k_2) \in K_1 \times K_2\} \quad (2.3.1)$$

**2.4. Black/white decomposition and a Steinberg's theorem.** Cf. [Stein], [C], [BS]. Let  $\alpha_1, \dots, \alpha_r$  be a base of simple roots of a finite reduced irreducible root system  $R$ . Suppose that  $R$  is simply laced.

Let

$$A = (a_{ij}) = (\langle \alpha_i, \alpha_j^\vee \rangle)$$

be the Cartan matrix; suppose that  $R$  is simply laced, so that  $A$  is symmetric.

Choose a black/white coloring of the set of vertices of the corresponding Dynkin graph in such a way that any two neighbouring vertices have different colours.

Let us choose an ordering of simple roots in such a way that the first  $r$  roots are black, and the last  $r - p$  roots are white. In this base  $A$  has a block form

$$A = \begin{pmatrix} 2I_p & X \\ X^t & 2I_{r-p} \end{pmatrix}$$

Consider a Coxeter element

$$C = s_1 s_2 \dots s_r = C_B C_W, \quad (2.4.1)$$



where

$$C_B = \prod_{i=1}^p s_i, \quad C_W = \prod_{i=p+1}^r s_i.$$

Here  $s_i$  denotes the simple reflection corresponding to the root  $\alpha_i$ .

The matrices of  $C_B, C_W$  with respect to the base  $\{\alpha_i\}$  are

$$C_B = \begin{pmatrix} -I & -X \\ 0 & I \end{pmatrix}, \quad C_W = \begin{pmatrix} I & 0 \\ -X^t & -I \end{pmatrix},$$

so that

$$C_B + C_W = 2I - A. \quad (2.4.2)$$

This is an observation due to R.Steinberg, cf. [Stein], p. 591.

**2.4.1. Remark: a matrix quadratic polynomial.** We can express the basic relations (2.4.1), (2.4.2) by saying that  $C_B, C_W$  are the "roots" of a matrix quadratic polynomial

$$t^2 + (A - 2)t + C = (t - C_B)(t - C_W),$$

□.

We can also rewrite this as follows. Set

$$L = \begin{pmatrix} I & X \\ 0 & I \end{pmatrix}.$$

Then  $A = L + L^t$ , and one checks easily that

$$C = -L^{-1}L^t, \quad (2.4.3)$$

so we are in the situation 2.1. This explains the name "Cartan - Coxeter coresspondence".

**2.5. Eigenvectors' correspondence.** The following corollary was obtained by the physicists, cf. [F], or [BS], Corollary 3.7.

**Theorem.** *In the notations of 2.3, a vector*

$$x = \sum x_i \alpha_i$$

*is an eigenvector of  $A$  with the eigenvalue  $2(1 - \cos \theta)$  iff the vector*

$$x_c := \sum e^{\pm i\theta/2} x_i$$

*where the sign in  $e^{\pm i\theta/2}$  is plus if  $i$  is a white vertex, and minus otherwise, is an eigenvector of  $c$  with eigenvalue  $e^{2i\theta}$ .*

□

### §3. Sebastiani - Thom factorization of $E_8$

**3.1. Root systems  $A_n$ .** We consider the Dynkin graph of  $A_n$  with the obvious numbering of the vertices.

The Coxeter number  $h = n + 1$ , the set of exponents:

$$\text{Exp}(A_n) = \{1, 2, \dots, n\}$$

The eigenvalues of any Coxeter element are  $e^{i\theta_k}$ , and the eigenvalues of the Cartan matrix  $A(A_n)$  are  $2 - 2 \cos \theta_k$ ,  $\theta_k = 2\pi k/h$ ,  $k \in \text{Exp}(A_n)$ .

An eigenvector of  $A(A_n)$  with the eigenvalue  $2 - 2 \cos \theta$  has the form

$$x(\theta) = \left( \sum_{k=0}^{n-1} e^{i(n-1-2k)\theta}, \sum_{k=0}^{n-2} e^{i(n-2-2k)\theta}, \dots, 1 \right) \quad (3.1.1)$$

Denote by  $C(A_n)$  the Coxeter element

$$C(A_n) = s_1 s_2 \dots s_n$$

Its eigenvector with the eigenvalue  $e^{2i\theta}$  is:

$$X_{C(A_n)} = \left( \sum_{k=0}^{n-j} e^{2ik\theta} \right)_{1 \leq j \leq n}$$

For example, for  $n = 4$ :

$$C_{A_4} = \begin{pmatrix} 0 & 0 & 0 & -1 \\ 1 & 0 & 0 & -1 \\ 0 & 1 & 0 & -1 \\ 0 & 0 & 1 & -1 \end{pmatrix}$$

and

$$X_{C(A_4)} = \begin{pmatrix} 1 + e^{2i\theta} + e^{4i\theta} + e^{6i\theta} \\ 1 + e^{2i\theta} + e^{4i\theta} \\ 1 + e^{2i\theta} \\ 1 \end{pmatrix}$$

is an eigenvector with eigenvalue  $e^{2i\theta}$ .

Similarly, for  $n = 2$ :

$$C_{A_2} = \begin{pmatrix} 0 & -1 \\ 1 & -1 \end{pmatrix}, \quad X_{C(A_2)} = \begin{pmatrix} 1 + e^{2i\theta} \\ 1 \end{pmatrix}$$

### 3.2. $E_8$ versus $A_4 * A_2 * A_1$ : elementary analysis.

The ranks:

$$r(E_8) = 8 = r(A_4)r(A_2)r(A_1);$$

the Coxeter numbers:

$$h(E_8) = h(A_4)h(A_2)h(A_1) = 5 \cdot 3 \cdot 2 = 30.$$

It follows that

$$|R(E_8)| = 240 = |R(A_4)||R(A_2)||R(A_1)|.$$

The exponents of  $E_8$  are:

$$1, 7, 13, 19, 11, 17, 23, 29.$$

All these numbers, except 1, are primes, and these are all primes  $\leq 30$ , not dividing 30.

They may be determined from the formula

$$\frac{i}{5} + \frac{j}{3} + \frac{1}{2} = \frac{30 + k(i, j)}{30}, \quad 1 \leq i \leq 4, \quad 1 \leq j \leq 2,$$

so

$$k(i, 1) = 1 + 6(i - 1) = 1, 7, 13, 19;$$

$$k(i, 2) = 1 + 10 + 6(i - 1) = 11, 17, 23, 29.$$

This shows that the exponents of  $E_8$  are the same as the exponents of  $A_4 * A_2 * A_1$ .

The following theorem is more delicate.

**3.3. Theorem** (Gabrielov). *There exists a polarization of the root lattice  $Q(E_8)$  and an isomorphism of polarized lattices*

$$\Gamma : Q(A_4) * Q(A_2) * Q(A_1) \xrightarrow{\sim} Q(E_8). \quad (3.3.1)$$

In the left hand side  $Q(A_n)$  means the root lattice of  $A_n$  with the standard Cartan matrix and the standard polarization

$$A(A_n) = L(A_n) + L(A_n)^t$$

where the Seifert matrix  $L(A_n)$  is upper triangular.

In the process of the proof, given in 3.4 - 3.6 below, the isomorphism  $\Gamma$  will be written down explicitly.

### 3.4. Beginning of the proof.

For  $n = 4, 2, 1$ , we consider the bases of simple roots  $e_1, \dots, e_n$  in  $Q(A_n)$ , with scalar products given by the Cartan matrices  $A(A_n)$ .

The tensor product of three lattices

$$Q_* = Q(A_4) \otimes Q(A_2) \otimes Q(A_1)$$

will be equipped with the "factorizable" basis in the lexicographic order:

$$(f_1, \dots, f_8) := (e_1 \otimes e_1 \otimes e_1, e_1 \otimes e_2 \otimes e_1, e_2 \otimes e_1 \otimes e_1, e_2 \otimes e_2 \otimes e_1, \\ e_3 \otimes e_1 \otimes e_1, e_3 \otimes e_2 \otimes e_1, e_4 \otimes e_1 \otimes e_1, e_4 \otimes e_2 \otimes e_1).$$

Introduce a scalar product  $(x, y)$  on  $Q_*$  given, in the basis  $\{f_i\}$ , by the matrix

$$A_* = A_4 * A_2 * A_1.$$

### 3.5. Gabrielov - Picard - Lefschetz transformations $\alpha_m, \beta_m$ .

Let  $(Q, (,))$  be a lattice of rank  $r$ . We introduce the following two sets of transformations  $\{\alpha_m\}, \{\beta_m\}$  on the set *Bases - cycl(Q)* of cyclically ordered bases of  $Q$ .

If  $x = (x_i)_{i \in \mathbb{Z}/r\mathbb{Z}}$  is a base, and  $m \in \mathbb{Z}/r\mathbb{Z}$ , we set

$$(\alpha_m(x))_i = \begin{cases} x_{m+1} + (x_{m+1}, x_m)x_m & \text{if } i = m \\ x_m & \text{if } i = m + 1 \\ x_i & \text{otherwise} \end{cases}$$

and

$$(\beta_m(x))_i = \begin{cases} x_m & \text{if } i = m - 1 \\ x_{m-1} + (x_{m-1}, x_m)x_m & \text{if } i = m \\ x_i & \text{otherwise} \end{cases}$$

We define also a transformation  $\gamma_m$  by

$$(\gamma_m(x))_i = \begin{cases} -x_m & \text{if } i = m \\ x_i & \text{otherwise} \end{cases}$$

**3.6. Passage from  $A_4 * A_2 * A_1$  to  $E_8$ .** Consider the base  $f = \{f_1, \dots, f_8\}$  of the lattice  $Q_* := Q(A_4) \otimes Q(A_2) \otimes Q(A_1)$  described in 3.4, and apply to it the following transformation

$$G' = \gamma_2 \gamma_1 \beta_4 \beta_3 \alpha_3 \alpha_4 \beta_4 \alpha_5 \alpha_6 \alpha_7 \alpha_1 \alpha_2 \alpha_3 \alpha_4 \beta_6 \beta_3 \alpha_1, \quad (3.6.1)$$

cf. [Gab], Example 3.

Then the base  $G'(f)$  has the intersection matrix given by the Dynkin graph of  $E_8$ , with the ordering indicated on Figure 1 below.



**3.9. Lemma.** Let  $A(E_8)$  be the standard Cartan matrix of  $E_8$  from [B]:

$$A(E_8) = \begin{pmatrix} 2 & 0 & -1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 2 & 0 & -1 & 0 & 0 & 0 & 0 \\ -1 & 0 & 2 & -1 & 0 & 0 & 0 & 0 \\ 0 & -1 & -1 & 2 & -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 & 2 & -1 & 0 & 0 \\ 0 & 0 & 0 & 0 & -1 & 2 & -1 & 0 \\ 0 & 0 & 0 & 0 & 0 & -1 & 2 & -1 \\ 0 & 0 & 0 & 0 & 0 & 0 & -1 & 2 \end{pmatrix}.$$

Then

$$A(E_8) = G^t A_* G$$

and

$$C_G(E_8) = G^{-1} C_* G$$

where

$$C_* = C(Q(A_4) * Q(A_2) * Q(A_1)) = C(A_4) \otimes C(A_2) \otimes C(A_1),$$

is the factorized Coxeter element, and

$$G = \begin{pmatrix} 0 & 0 & 0 & 1 & -1 & 0 & 0 & 0 \\ -1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -1 & 1 & 0 & 0 & 0 & 0 \\ -1 & 1 & -1 & 0 & 0 & 1 & 0 & 0 \\ 0 & 1 & -1 & 0 & 0 & 0 & 1 & 0 \\ -1 & 1 & -1 & 0 & 0 & 0 & 1 & 0 \\ 0 & 1 & -1 & 0 & 0 & 0 & 0 & 1 \\ 0 & 1 & -1 & 0 & 0 & 0 & 0 & 0 \end{pmatrix} \quad (3.8.1)$$

Here

$$G = G' P$$

where  $P$  is the permutation matrix of passage from the Gabrielov's ordering on Fig. 1 to the Bourbaki ordering on Fig. 2.

**3.9. Cartan eigenvectors of  $E_8$ .** To obtain the Cartan eigenvectors of  $E_8$ , one should pass from  $C_G(E_8)$  to a "black/white" coxeter element (as in 2.4)

$$C_{BW}(E_8) = s_1 \circ s_4 \circ s_6 \circ s_8 \circ s_2 \circ s_3 \circ s_5 \circ s_7$$

Any two Coxeter elements are conjugate in the Weyl group  $W(E_8)$ .

The elements  $C_G(E_8)$  and  $C_{BW}(E_8)$  are conjugate by the following element of  $W(E_8)$ :

$$C_G(E_8) = w^{-1} C_{BW}(E_8) w$$

where

$$w = s_7 \circ s_5 \circ s_3 \circ s_2 \circ s_6 \circ s_4 \circ s_5 \circ s_1 \circ s_3 \circ s_2 \circ s_4 \circ s_1 \circ s_3 \circ s_2 \circ s_1 \circ s_2$$

This expression for  $w$  can be obtained using an algorithm described in [C].

Thus, if  $x_*$  is an eigenvector of  $C_*(E_8)$  then

$$x_{BW} = wG^{-1}x_*$$

is an eigenvector of  $C_{BW}(E_8)$ . But we know the eigenvectors of  $C_*(E_8)$ , they are all factorizable.

This provides the eigenvectors of  $C_{BW}(E_8)$ , which in turn have very simple relation to the eigenvectors of  $A(E_8)$ , due to Theorem 2.5.

**Conclusion: an expression for the eigenvectors of  $A(E_8)$ .**

Let  $\theta = \frac{a\pi}{5}$ ,  $1 \leq a \leq 4$ ,  $\gamma = \frac{b\pi}{3}$ ,  $1 \leq b \leq 2$ ,  $\delta = \frac{\pi}{2}$ ,

$$\alpha = \theta + \gamma + \delta = 1 + \frac{k\pi}{30},$$

$$k \in \{1, 7, 11, 13, 17, 19, 23, 29\}.$$

The 8 eigenvalues of  $A(E_8)$  have the form

$$\lambda(\alpha) = \lambda(\theta, \gamma) = 2 - 2 \cos \alpha$$

An eigenvector of  $A(E_8)$  with the eigenvalue  $\lambda(\theta, \gamma)$  is

$$X_{E_8}(\theta, \gamma) = \begin{pmatrix} \cos(\gamma + \theta - \delta) + \cos(\gamma - 3\theta - \delta) + \cos(\gamma - \theta - \delta) \\ \cos(2\gamma + 2\theta) \\ \cos(2\gamma) + \cos(2\gamma + 2\theta) + \cos(2\gamma - 2\theta) + \cos(4\theta) + \cos(2\theta) \\ \cos(\gamma + 3\theta - \delta) + \cos(\gamma + \theta - \delta) + \cos(-\gamma + 3\theta - \delta) \\ 2 \cos(2\gamma) + 2 \cos(2\gamma + 2\theta) + \cos(2\gamma - 2\theta) + \cos(2\gamma + 4\theta) + \cos(4\theta) + 2 \cos(2\theta) + 1 \\ \cos(\gamma + 3\theta - \delta) + \cos(\gamma + \theta - \delta) \\ \cos(2\gamma) + \cos(2\theta - 2\delta) \\ \cos(\gamma - \theta - \delta) \end{pmatrix}$$

One can simplify it as follows:

$$X_{E_8}(\theta, \gamma) = \begin{pmatrix} 2 \cos(4\theta) \cos(\gamma - \theta - \delta) \\ -\cos(2\gamma + 2\theta) \\ 2 \cos^2(\theta) \\ -2 \cos(\gamma) \cos(3\theta - \delta) - \cos(\gamma + \theta - \delta) \\ -2 \cos(2\gamma + 3\theta) \cos(\theta) + \cos(2\gamma) \\ -2 \cos \theta \cos(\gamma + 2\theta - \delta) \\ -2 \cos(\gamma + \theta - \delta) \cos(\gamma - \theta + \delta) \\ -\cos(\gamma - \theta - \delta) \end{pmatrix} \quad (3.9.1)$$

**3.10. Perron - Frobenius and all that.** The Perron - Frobenius eigenvector corresponds to the eigenvalue

$$2 - 2 \cos \frac{\pi}{30},$$

and may be chosen as

$$v_{PF} = \begin{pmatrix} 2 \cos \frac{\pi}{5} \cos \frac{11\pi}{30} \\ \cos \frac{\pi}{15} \\ 2 \cos^2 \frac{\pi}{5} \\ 2 \cos \frac{2\pi}{30} \cos \frac{\pi}{30} \\ 2 \cos \frac{4\pi}{15} \cos \frac{\pi}{5} + \frac{1}{2} \\ 2 \cos \frac{\pi}{5} \cos \frac{7\pi}{30} \\ 2 \cos \frac{\pi}{30} \cos \frac{11\pi}{30} \\ \cos \frac{11\pi}{30} \end{pmatrix}$$

If we order its coordinates in the increasing order, we get

$$v_{PF<} = \begin{pmatrix} \cos \frac{11\pi}{30} \\ 2 \cos \frac{\pi}{5} \cos \frac{11\pi}{30} \\ 2 \cos \frac{\pi}{30} \cos \frac{11\pi}{30} \\ \cos \frac{\pi}{15} \\ 2 \cos \frac{\pi}{5} \cos \frac{7\pi}{30} \\ 2 \cos^2 \frac{\pi}{5} \\ 2 \cos \frac{4\pi}{15} \cos \frac{\pi}{5} + \frac{1}{2} \\ 2 \cos \frac{2\pi}{30} \cos \frac{\pi}{30} \end{pmatrix}$$

Zamolodchikov gives in [Z] the following expression for a PF vector:

$$v_{Zam}(m) = \begin{pmatrix} m \\ 2m \cos \frac{\pi}{5} \\ 2m \cos \frac{\pi}{30} \\ 4m \cos \frac{\pi}{5} \cos \frac{7\pi}{30} \\ 4m \cos \frac{\pi}{5} \cos \frac{2\pi}{15} \\ 4m \cos \frac{\pi}{5} \cos \frac{\pi}{30} \\ 8m \cos^2 \frac{\pi}{5} \cos \frac{7\pi}{30} \\ 8m \cos^2 \frac{\pi}{5} \cos \frac{2\pi}{15} \end{pmatrix}$$

Setting  $m = \cos \frac{11\pi}{30}$ , we have indeed :

$$v_{PF<} = v_{Zam}\left(\cos \frac{11\pi}{30}\right)$$



**3.11. Another form of the eigenvectors' matrix.** As was noticed in [BS], the coordinates of all eigenvectors of  $A(E_8)$  may be obtained from the coordinates of the PF vector by some permutations and sign changes.

Namely, if  $(z_1, \dots, z_8)$  is a PF vector then the other eigenvectors are the columns of the matrix

$$Z = \begin{pmatrix} z_1 & z_7 & z_4 & z_2 & z_2 & z_4 & z_7 & z_1 \\ z_2 & z_1 & -z_7 & -z_4 & z_4 & z_7 & -z_1 & -z_2 \\ z_3 & z_6 & z_5 & z_8 & -z_8 & -z_5 & -z_6 & -z_3 \\ z_4 & z_2 & -z_1 & -z_7 & -z_7 & -z_1 & z_2 & z_4 \\ z_5 & -z_8 & -z_3 & z_6 & -z_6 & z_3 & z_8 & -z_5 \\ z_6 & -z_5 & -z_8 & z_3 & z_3 & -z_8 & -z_5 & z_6 \\ z_7 & -z_4 & z_2 & -z_1 & z_1 & -z_2 & z_4 & -z_7 \\ z_8 & -z_3 & z_6 & -z_5 & -z_5 & z_6 & -z_3 & z_8 \end{pmatrix}$$

However, these eigenvectors differ from the ones given by the formula (3.9.1): the latter ones are proportional to the former ones.

## §4. Givental's $q$ -deformations

**4.1. Definition.** Let  $(Q, A, L)$  be a polarized lattice. We define a  $q$ -deformed Cartan matrix by

$$A(q) = L + qL^t.$$

This definition is inspired by the  $q$ -deformed Picard - Lefschetz theory developed by Givental, [Giv].

**4.2. "Black/white"  $q$ -deformation.** Let

$$A = \begin{pmatrix} 2 & B \\ C & 2 \end{pmatrix}$$

be a block matrix, and  $v = \begin{pmatrix} x \\ y \end{pmatrix}$  its eigenvector with the eigenvalue  $\lambda$ . This means that

$$By = (\lambda - 2)x, \quad Cx = (\lambda - 2)y. \quad (4.2.1)$$

Consider the matrix

$$A(q) = \begin{pmatrix} 1+q & B \\ qC & 1+q \end{pmatrix},$$

and let us look for its eigenvector in the form  $v = \begin{pmatrix} x \\ by \end{pmatrix}$  with eigenvalue  $\mu$ . This would mean that

$$By = \frac{\mu - q - 1}{b}x, \quad Cx = \frac{b(\mu - q - 1)}{q}x. \quad (4.2.2)$$

Comparing (4.2.1) and (4.2.2) we conclude that

$$\frac{\mu'}{b} = \frac{b\mu'}{q}, \quad \mu' = \mu - q - 1,$$

whence  $b = \sqrt{q}$ , and that

$$\lambda - 2 = \frac{\mu'}{\sqrt{q}}.$$

Conclusion: *the vector*

$$v(q) = \begin{pmatrix} x \\ q^{1/2}y \end{pmatrix}$$

*is an eigenvector of  $A(q)$  with eigenvalue*

$$\lambda(q) = 1 + (\lambda - 2)q^{1/2} + q. \quad (4.2.3)$$

For another approach to the same values see 4.10 below.

**4.2.1. A generalization.** More generally, let

$$M = \begin{pmatrix} 2I & A & 0 \\ C & 2I & B \\ 0 & D & 2I \end{pmatrix}$$

be a block Jacobi matrix which admits an eigenvector  $v = \begin{pmatrix} x \\ y \\ z \end{pmatrix}$  for the eigenvalue  $\lambda = 2 - 2 \cos \theta$ .

Then its  $q$ -deformation

$$M(q) = \begin{pmatrix} (1+q)I & A & 0 \\ qC & (1+q)I & D \\ 0 & qD & (1+q)I \end{pmatrix}$$

admits an eigenvector

$$v(q) = \begin{pmatrix} x \\ q^{1/2}y \\ qz \end{pmatrix}$$

with the eigenvalue

$$\lambda(q) = 1 + q - 2q^{\frac{1}{2}} \cos \theta.$$

This can be generalized to the block Jacobi matrices of any size.

**4.3. Remark** (M.Finkelberg). The expression (4.2.3) resembles the number of points of an elliptic curve  $X$  over a finite field  $\mathbb{F}_q$ . To better appreciate this resemblance, note that in all our examples  $\lambda$  will have the form

$$\lambda = 2 - 2 \cos \theta,$$

so if we set

$$\alpha = \sqrt{q}e^{i\theta}$$

(“a Frobenius root”) then  $|\alpha| = \sqrt{q}$ , and

$$\lambda(q) = 1 - \alpha - \bar{\alpha} + q,$$

cf. [IR], Chapter 11, §1, [Kn], Chapter 10, Theorem 10.5.

So, the Coxeter eigenvalues  $e^{2i\theta}$  may be seen as analogs of “Frobenius roots of an elliptic curve over  $\mathbb{F}_1$ ”.

**4.4. Standard deformation for  $A_n$ .** Let us consider the following  $q$ -deformation of  $A = A(A_n)$ :

$$A(q) = \begin{pmatrix} 1+q & -1 & 0 & \dots & 0 \\ -q & 1+q & -1 & \dots & 0 \\ \dots & \dots & \dots & \dots & \dots \\ 0 & \dots & 0 & -q & 1+q \end{pmatrix}$$

Then

$$\text{Spec}(A(q)) = \{\lambda(q) := 1 + (\lambda - 2)\sqrt{q} + q \mid \lambda \in \text{Spec}(A(1))\}.$$

If  $x = (x_1, \dots, x_n)$  is an eigenvector of  $A = A(1)$  with eigenvalue  $\lambda$  then

$$x(q) = (x_1, q^{1/2}x_2, \dots, q^{(n-1)/2}x_n)$$

is an eigenvector of  $A(q)$  with eigenvalue  $\lambda(q)$ .

**4.5. Standard deformation for  $E_8$ .** A  $q$ -deformation:

$$A_{E_8}(q) = \begin{pmatrix} 1+q & 0 & -1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1+q & 0 & -1 & 0 & 0 & 0 & 0 \\ -q & 0 & 1+q & -1 & 0 & 0 & 0 & 0 \\ 0 & -q & -q & 1+q & -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & -q & 1+q & -1 & 0 & 0 \\ 0 & 0 & 0 & 0 & -q & 1+q & -1 & 0 \\ 0 & 0 & 0 & 0 & 0 & -q & 1+q & -1 \\ 0 & 0 & 0 & 0 & 0 & 0 & -q & 1+q \end{pmatrix}$$

Its eigenvalues are

$$\lambda(q) = 1 + q + (\lambda - 2)\sqrt{q} = 1 + q - 2\sqrt{q} \cos \theta$$

where  $\lambda = 2 - 2 \cos \theta$  is an eigenvalue of  $A(E_8)$ .

If  $X = (x_1, x_2, x_3, x_4, x_5, x_6, x_7, x_8)$  is an eigenvector of  $A(E_8)$  for the eigenvalue  $\lambda$ , then

$$X = (x_1, \sqrt{q}x_2, \sqrt{q}x_3, qx_4, q\sqrt{q}x_5, q^2x_6, q^2\sqrt{q}x_7, q^3x_8) \quad (4.5.1)$$

is an eigenvector of  $A_{E_8}(q)$  for the eigenvalue  $\lambda(q)$ .

We see the patterns here:

(i) The eigenvalues are always  $\lambda(q)$ , as in the black/white case, so the  $q$ -deformed Cartan matrices for different orderings are conjugate.

(ii) The coordinates of the  $q$ -deformed eigenvectors are equal to  $q^i \times$  the coordinates of the original eigenvector.

For an explanation, see 4.7, 4.8 below.

#### 4.6. Examples of conjugation.

(a) Let

$$A(A_4; q) = \begin{pmatrix} 1+q & -1 & 0 & 0 \\ -q & 1+q & -1 & 0 \\ 0 & -q & 1+q & -1 \\ 0 & 0 & -q & 1+q \end{pmatrix}$$

be the  $q$ -deformation of the standard  $A(A_4)$ , and

$$A_{BW}(A_4; q) = \begin{pmatrix} 1+q & 0 & -1 & 0 \\ 0 & 1+q & -1 & -1 \\ -q & -q & 1+q & 0 \\ 0 & -q & 0 & 1+q \end{pmatrix}$$

that of the "black/white" one. Then

$$A(A_4; q) = P(q) \cdot A_{BW}(A_4; q) \cdot P^{-1}(q),$$

with

$$P(q) = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & q & 0 & 0 \\ 0 & 0 & 0 & q \end{pmatrix} \text{ and } P^{-1}(q) = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & q^{-1} & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & q^{-1} \end{pmatrix}$$

The matrix  $P(q)$  is a deformation of a permutation matrix.

(b) A similar example for  $E_8$ :

$$A(E_8; q) = P(q) \cdot A_{BW}(E_8; q) \cdot P^{-1}(q),$$

with

$$P(q) = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & q & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & q & 0 \\ 0 & 0 & q^2 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & q^2 \\ 0 & 0 & 0 & q^3 & 0 & 0 & 0 & 0 \end{pmatrix}$$

In fact, always

$$P(q) = D(q)P(1) = P(1)D'(q)$$

where  $P(1)$  is a permutation matrix, and  $D(q), D'(q)$  are diagonal matrices, with some natural powers  $q^n$  on the diagonal.

Below follows an explanation of what is going on.

#### 4.7. Conjugacy of different $q$ -deformations.

Let  $A = (a_{ij}) \in \mathfrak{gl}_r(\mathbb{C})$  be a symmetric matrix, and

$$A = L + L^t,$$

the standard polarization, with  $L$  upper triangular. Thus,  $L = (\ell_{ij})$ , with  $\ell_{ii} = a_{ii}/2$ , and

$$\ell_{ij} = \begin{cases} a_{ij} & \text{if } i < j \\ 0 & \text{if } i > j \end{cases}$$

For a bijection

$$\sigma : \{1, \dots, r\} \xrightarrow{\sim} \{1, \dots, r\},$$

define a matrix  $L_\sigma = (\ell_{ij}^\sigma)$  with  $\ell_{ii}^\sigma = a_{ii}/2$ , and

$$\ell_{ij}^\sigma = \begin{cases} a_{ij} & \text{if } \sigma(i) < \sigma(j) \\ 0 & \text{if } \sigma(i) > \sigma(j) \end{cases}$$

Obviously

$$A = L_\sigma + L_\sigma^t,$$

i.e. we have got a different, "σ - twisted", polarization of  $A$ .

Consider the  $q$ -deformations of  $A$  corresponding to these two polarizations:

$$A(q) = L + qL^t$$

and

$$A(\sigma; q) = L_\sigma + qL_\sigma^t.$$

Let us ask a

**4.7.1. Question.** *Find a diagonal matrix of the form*

$$D = D(\sigma; q) = \text{Diag}(q^{n_1}, \dots, q^{n_r}) \quad (4.7.1)$$

such that

$$A(\sigma; q) = D^{-1}A(q)D. \quad (4.7.2)$$

This equation means that

$$A(\sigma; q)_{ij} = q^{-n_i+n_j} A(q)_{ij}$$

for all  $i, j$ . This is an overdetermined system of  $r(r-1)/2$  equations on  $r$  variables  $n_i$ , not solvable in general.

However, we have something positive to say.

Let  $A = (a_{ij})$  be any matrix having the property

$$a_{ij} \neq 0 \text{ implies } a_{ji} \neq 0 \quad (4.7.3)$$

Let us assign to  $A$  its "Dynkin graph"  $\Gamma(A)$  having  $[r] := \{1, \dots, r\}$  as the set of vertices, vertices  $i$  and  $j$  being connected iff  $a_{ij} \neq 0$ .

Now let us return to our symmetric matrix  $A$ .

**4.7.1. Tree lemma.** *If  $\Gamma(A)$  is a tree then for each  $\sigma \in S_r = \text{Aut}([r])$  there exists  $D = D(\sigma; q)$  such that (4.7.2) holds true.*

A proof will be given in 4.9 below (the reader may also wish to take it as an exercise).

For the moment we assume this assertion.

**4.8. Now let us turn a permutation matrix  $P$  on.**

Define a permutation matrix  $P = P(\sigma) = (p_{ij}) \in GL_r(\mathbb{Z})$ , by

$$p_{ij} = \delta_{i, \sigma(j)}.$$

Then its inverse

$$P^{-1} = (p'_{ij}), \quad p'_{ij} = \delta_{\sigma(i),j} = \delta_{i,\sigma^{-1}(j)} = P^t \quad (4.8.1)$$

Define

$$A' = P^{-1}AP = (a'_{ij}).$$

Then

$$a'_{ij} = a_{\sigma(i),\sigma(j)}. \quad (4.8.2)$$

It follows that  $A'$  is symmetric as well.

Decompose

$$A' = L' + L'^t, \quad (4.8.3)$$

with  $L'$  upper triangular.

Define a  $q$ -deformation

$$A'(q) = L' + qL'^t,$$

Let us look for a matrix  $P(q)$  of the form

$$P(q) = PD(q),$$

$D(q) = D(\sigma; q)$  being as in (4.7.1), such that

$$A'(q) = P(q)^{-1}A(q)P(q). \quad (4.8.4)$$

(Note that

$$D(q)P = PD'(q)$$

where

$$D'(q) = \text{diag}(q^{n'_1}, \dots, q^{n'_r}), \quad n'_i = n_{\sigma^{-1}(i)},$$

cf. (4.8.2).)

The matrix  $A'$  has two polarizations: the first, the standard one, (4.8.3), and the second, the  $\sigma$ -twisted one:

$$A' = L_\sigma + L_\sigma^t, \quad L_\sigma := P^{-1}LP.$$

So, we are in the situation 4.7, and we are looking for a diagonal matrix  $D$ .

If we are lucky, and  $D$  exists, for example, according to the Tree lemma 4.7.1, if  $\Gamma(A)$  is a tree, then the problem (4.8.4) is solved.

This is the case for the finite Cartan matrices of types  $A, D$  or  $E$ .

**4.9. Homology of the Dynkin graph, the diagonal conjugacy, and a proof of Tree Lemma.** Under the assumptions of Lemma 4.7.1, pick any orientation on the Dynkin graph  $\Gamma = \Gamma(A)$ ; let  $\vec{\Gamma}$  denote the oriented graph thus obtained.

Consider a cochain complex

$$0 \longrightarrow C^0(\vec{\Gamma}; \mathbb{C}^*) \xrightarrow{d} C^1(\vec{\Gamma}; \mathbb{C}^*) \longrightarrow 0$$

We have  $\text{Coker}(d) = H^1(\vec{\Gamma}; \mathbb{C}^*) = 0$  since  $\Gamma$  is a tree. This means the following:

(i) *For any collection of numbers  $\{b_{ij} \in \mathbb{C}^*, (i, j) \in [r]^2, i \neq j\}$  such that*

$$b_{ji} = b_{ij}^{-1}, \quad (4.9.1)$$

*there exists a collection  $\{c_i \in \mathbb{C}^*, i \in [r]\}$  such that*

$$b_{ij} = c_j/c_i. \quad (4.9.2)$$

*Moreover, we can choose the numbers  $c_i$  in such a way that they are products of some  $b_{ij}$ .*

To prove the last assertion, let us choose a trunk of our tree  $\Gamma$ , and partially order its vertices by taking the minimal vertex  $i_0$  to be the beginning of the trunk, and then going "upstairs". This defines an orientation on  $\Gamma$ . Now, given a 1-cochain  $(b_{ij})$ , we set  $c_{i_0} = 1$  and then define the other  $c_i$  one by one, by going upstairs, and using as a definition

$$c_j = b_{ij}c_i, \quad i < j.$$

Obviously, the numbers  $c_i$  defined in such a way, are products of  $b_{ab}$ .

□

If we don't suppose  $\Gamma$  to be a tree then to get a solution of (4.9.2), in addition to (4.9.1) one should impose on  $\{b_{ij}\}$  one more condition:

for any non-contractible loop  $i_1 \longrightarrow i_2 \longrightarrow \dots \longrightarrow i_k \longrightarrow i_1$  in  $\Gamma$ ,

$$b_{i_1 i_2} b_{i_2 i_3} \dots b_{i_k i_1} = 1. \quad (4.9.1a)$$

(One could restrict to the loops representing some set of generators of  $H_1(\vec{\Gamma}; \mathbb{Z})$ .)

Let us return to the conditions of Lemma 4.7.1.

(ii) *Let  $A' = (a'_{ij})$  be another matrix with  $\Gamma(A') = \Gamma(A)$  such that  $a'_{ii} = a_{ii}$  for all  $i$ , and*

$$a'_{ij}/a_{ij} = a_{ji}/a'_{ji}$$

*for all  $i \neq j$ . Then there exists a diagonal matrix*

$$D = \text{Diag}(c_1, \dots, c_r)$$

*such that  $A' = D^{-1}AD$ .*

This is a corollary of (i), namely, set  $b_{ij} = a'_{ij}/a_{ij}$ .

Now we can prove the Tree Lemma. In fact, two matrices  $A(q)$  and  $A_\sigma(q)$  satisfy the conditions of (ii), with  $b_{ij} = 1, q$ , or  $q^{-1}$ .

We can choose  $c_i$  to be the integer powers of  $q$  due to the last assertion of (i).



**4.10. Another way to compute the eigenvalues and eigenvectors of  $A(q)$ .** Let  $A$  be a symmetric generalized Cartan matrix, and  $A(q)$  its standard  $q$ -deformation.

We can apply 4.9 (ii) to the matrices  $A(q)$  and a symmetric

$$A'(q) = \sqrt{q}A + (1 - \sqrt{q})^2I$$

So, there exists a diagonal matrix  $D$  as above such that

$$A(q) = D^{-1}A'(q)D.$$

But the eigenvalues of  $A'(q)$  are obviously

$$\lambda(q) = \sqrt{q}\lambda + (1 - \sqrt{q})^2 = 1 + (\lambda - 2)\sqrt{q} + q.$$

If  $v$  is an eigenvector of  $A$  for  $\lambda$  then  $v$  is an eigenvector of  $A'(q)$  for  $\lambda(q)$ , and  $Dv$  will be an eigenvector of  $A(q)$  for  $\lambda(q)$ .

## References

[AGV] V.Arnold, S.Gusein-Zade, A.Varchenko, Singularities of differentiable maps, vol. II, Monodromy and asymptotic of integrals. В.И.Арнольд, А.Н.Варченко, С.М.Гусейн-Заде, Особенности дифференцируемых отображений, т. 2, М., 1984.

[B] N.Bourbaki, Groupes et algèbres de Lie, Chapitres IV - VI.

[BS] L.Brillon, V.Schechtman, Coxeter element and particle masses, <http://www.math.univ-toulouse.fr/schechtman/masse.pdf>.

[C] B.Casselman, Essays on Coxeter groups. Coxeter elements in finite Coxeter groups, <https://www.math.ubc.ca/cass/research/pdf/Element.pdf>

[D] P.Dorey, (a) Root systems and purely elastic  $S$ -matrices, *Nucl. Phys.* **B358** (1991), 654 - 676. (b) The exact  $S$ -matrices of affine Toda field theories, Durham PhD thesis (1990).

[F] M.D.Freeman, (a) On the mass spectrum of an affine Toda field theory, *Phys. Let.* **B 261** (1991), 57 - 61. (b) Conserved charges and soliton solutions in affine Toda theory, arXiv:hep/9408092.

[Gab] A.Gabriellov, Intersection matrices for some singularities. А.М.Габриэлов, Матрицы пересечений для некоторых особенностей, *Функц. анализ и его прил.*, **7** (1973), вып. 3, 18 - 32.

[Giv] A.Givental, Twisted Picard - Lefschetz formulas. А.Б.Гивенталь, Скрученные формулы Пикара - Лефшеца, *Функц. анализ и его прил.*, **22** (1988), вып. 1, 12 - 22.

- [IR] K.Ireland, M.Rosen, A classical introduction to modern Number theory.
- [Kn] A.Knapp, Elliptic curves.
- [L] K.Lamotke, Die Homologie isolierter Singularitäten, *Math. Z.* **143** (1975), 27 - 44.
- [Lef] S.Lefschetz, L'analysis situs et la géométrie algébrique, Gauthier-Villars, 1950.
- [M] J.Milnor, Singular points of complex hypersurfaces.
- [ST] M.Sebastiani, R.Thom, Un résultat sur la monodromie, *Inv. Math.* **13** (1971), 90 - 96.
- [Serre] J.-P.Serre, Cours d'arithmétique.
- [Stein] R.Steinberg, Finite subgroups of  $SU_2$ , Dynkin diagrams and affine Coxeter elements, *Pacific J. Math.* **118** (1985), 587 - 598.
- [Z] A.Zamolodchikov, (a) Integrable field theory from Conformal field theory, in: Integrable Systems in Quantum Field Theory and Statistical Mechanics, *Adv. Studies in Pure Math.* **19** (1989), 641 - 674. (b) Integrals of motion and  $S$ -matrix of the (scaled)  $T = T_c$  Ising model with magnetic field.