On asymptotic properties for some parameter dependent variational inequalities

Mihai Bostan*, Eric Canon*, Patrick Hild *
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Abstract

This paper is concerned with asymptotic and monotonicity properties of some parameter dependent variational inequalities. The main part of the study deals with inequalities modelling friction problems with normal compliance or Tresca's conditions in which the parameter stands for the friction coefficient. The corresponding inequalities are (generalizations) of variational inequalities of the second kind. We then study an inequality of the first kind representing the elastoplastic torsion problem where the parameter represents the plasticity yield.

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Abbreviated title: Analysis of some parameter dependent variational inequalities.

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1 Introduction

Problems written as weak formulations involving variational inequalities represent various nonlinear phenomena which occur in mechanics and physics. The nonlinearity corresponding to the inequality in the problem is located either in the entire domain of study (e.g., obstacle problem, elastoplastic torsion problem, Bingham fluid...) or only on the boundary of the domain (e.g., Signorini or contact problems, friction models...). Generally we deal with first and second kind classes of variational inequalities. We recall that a variational inequality of the first (resp. second) kind is of the form:

$$u \in C$$
: $a(u, v - u) \ge L(v - u), \quad \forall v \in C,$

^{*}Laboratoire de Mathématiques de Besançon, UMR CNRS 6623, Université de Franche-Comté, 16 route de Gray, 25030 Besançon Cedex, France. Email : mihai.bostan@univ-fcomte.fr, eric.canon@univ-fcomte.fr, patrick.hild@univ-fcomte.fr. This work is supported by "l'Agence Nationale de la Recherche", project ANR-05-JCJC-0182-01.

respectively,

$$u \in X$$
: $a(u, v - u) + j(v) - j(u) \ge L(v - u), \quad \forall v \in X,$

where X is an Hilbert space, $C \subset X$ is a nonempty closed convex set, $a(\cdot, \cdot)$ is bilinear, X-elliptic and continuous on $X \times X$, $L(\cdot)$ is linear and continuous on X, $j(\cdot)$ is proper convex and lower semicontinuous on X (with values in $\mathbb{R} \cup \{+\infty\}$). More details concerning variational inequalities of the first and second kind can be found in e.g., [1, 7]. Our aim in this paper is to establish some properties for parameter dependent variational inequalities of the first, second kind and also more general inequalities when the parameter varies. To our knowledge the kind of estimates we prove have not been considered previously. An outline of the paper is as follows.

Section 2 is concerned with variational inequalities modelling friction phenomena which correspond to nonlinear boundary value problems. We first consider the well-known compliance model in elasticity introduced in [19, 17] (see also [13]). The corresponding inequality is neither of the first nor of the second kind since it corresponds to a more complex type. Denoting by \mathbf{u}_{ε} the solution of the problem with a friction parameter ε , we show that $\varepsilon^{-1}(\mathbf{u}_{\varepsilon} - \mathbf{u}_{0})$ converges as ε vanishes and we characterize the limit as a solution of a specific variational inequality. Then we consider an elementary friction model introduced in [7] where the unknown of the variational inequality (of the second kind) is a scalar valued function. We obtain similar results as in the previous case. Moreover we show that a seizure property occurs when the friction is large enough. In addition we prove in the scalar case a monotonicity result in which the solution decreases when the parameter increases. Finally we mention that some of the previous studies can be applied to the so-called Tresca's friction model involving the elasticity operator and vector valued functions.

In Section 3 we consider a variational inequality of the first kind, the so-called elastoplastic torsion problem. The nonnegative parameter denoted ε , stands for the plasticity yield and u_{ε} is the corresponding solution. We prove that $u_{\varepsilon}/\varepsilon$ converges as ε vanishes towards a function which we characterize and which can be explicitly given in physical meaningful cases. We end this section with a monotonicity result for $u_{\varepsilon}/\varepsilon$.

2 Inequalities modelling friction problems

In this section we consider some inequality problems modelling friction phenomena. For the derivation and the discussion of various friction models in a more general framework we refer the reader to, e.g., [21].

2.1 A friction problem with normal compliance

We consider an elastic body that occupies a region Ω in \mathbb{R}^n (n=2,3). The body is held fixed on a part Γ_D of the boundary $\partial\Omega = \overline{\Gamma_D} \cup \overline{\Gamma_C}$, can come into contact with a reactive foundation over the part Γ_C and is subject to volume forces $\mathbf{f} \in (L^2(\Omega))^n$. On the contact surface we use a normal compliance condition (see [15, 16] and the references therein). Let us denote by $\mathbf{u} = (u_i)_{1 \leq i \leq n}$ the displacement vector and by $\boldsymbol{\sigma} = (\sigma_{ij})_{1 \leq i,j \leq n}$ the stress tensor such that

$$\sigma_{ij}(\mathbf{u}) = a_{ijhk} \frac{\partial u_h}{\partial x_k}, \quad i, j, h, k \in \{1, ..., n\},$$

where the summation convention of repeated indices is adopted. The functions $a_{ijhk} \in L^{\infty}(\Omega)$ are the coefficients of a fourth order tensor, representing the elastic properties of the material. As usual we assume that $a_{ijhk} = a_{jihk} = a_{hkij}$ and the ellipticity condition $a_{ijhk}\xi_{ij}\xi_{hk} \geq \alpha|\xi|^2$, $\forall \xi_{ij} = \xi_{ji}$, for some $\alpha > 0$.

The frictional contact problem with normal compliance in elastostatics is to find the displacement field \mathbf{u} such that equations (1)–(3) hold:

$$\mathbf{div}\ \boldsymbol{\sigma}(\mathbf{u}) + \mathbf{f} = \mathbf{0} \quad \text{in } \Omega, \tag{1}$$

$$\mathbf{u} = \mathbf{0} \quad \text{on } \Gamma_D, \tag{2}$$

where **div** denotes the divergence operator of tensor valued functions. For each displacement field \mathbf{v} and for each density of surface forces $\sigma(\mathbf{v})\mathbf{n}$ defined on $\partial\Omega$ we adopt the following notation (\mathbf{n} stands for the outward unit normal on $\partial\Omega$):

$$\mathbf{v} = v_N \mathbf{n} + \mathbf{v}_T$$
 and $\boldsymbol{\sigma}(\mathbf{v}) \mathbf{n} = \sigma_N(\mathbf{v}) \mathbf{n} + \boldsymbol{\sigma}_T(\mathbf{v})$.

Then the conditions of normal compliance with friction on Γ_C are:

$$\begin{cases}
\sigma_{N}(\mathbf{u}) = -c_{N}(u_{N})_{+}^{m_{N}}, \\
|\boldsymbol{\sigma}_{T}(\mathbf{u})| \leq c_{T}(u_{N})_{+}^{m_{T}} & \text{if} \quad \mathbf{u}_{T} = \mathbf{0}, \\
\begin{cases}
|\boldsymbol{\sigma}_{T}(\mathbf{u})| = c_{T}(u_{N})_{+}^{m_{T}} \\
\mathbf{u}_{T} = -\lambda \boldsymbol{\sigma}_{T}(\mathbf{u}), \quad \lambda \geq 0,
\end{cases} & \text{if} \quad \mathbf{u}_{T} \neq \mathbf{0},
\end{cases}$$
(3)

where (.)₊ stands for the positive part so that $(u_N)_+$ represents the penetration of the body into the foundation. The constants $m_N \geq 1$, $m_T \geq 1$ as well as the nonnegative functions c_N and c_T in $L^{\infty}(\Gamma_C)$ stand for interface parameters characterizing the contact behavior between the body and the rigid foundation. The set of admissible displacements

$$\mathbf{V} = \left\{ \mathbf{v} \in (H^1(\Omega))^n : \mathbf{v} = \mathbf{0} \text{ on } \Gamma_D \right\}$$

is endowed with the norm of $(H^1(\Omega))^n$. We denote by $a(\cdot,\cdot)$ the standard bilinear form of linear elasticity

$$a(\mathbf{u}, \mathbf{v}) = \int_{\Omega} \boldsymbol{\sigma}(\mathbf{u}) : \boldsymbol{\varepsilon}(\mathbf{v}) \, dx = \int_{\Omega} a_{ijhk} \frac{\partial u_i}{\partial x_j} \frac{\partial v_h}{\partial x_k} \, dx,$$

where $\varepsilon = (\varepsilon_{ij})_{1 \leq i,j \leq n}$ stands for the strain tensor where $\varepsilon_{ij}(\mathbf{u}) = (\partial u_i/\partial x_j + \partial u_j/\partial x_i)/2$. If Γ_D has positive superficial measure it is well known that the bilinear form $a(\cdot,\cdot)$ is **V**-elliptic

$$a(\mathbf{v}, \mathbf{v}) \ge \alpha \int_{\Omega} \sum_{i,j=1}^{n} \left| \frac{\partial v_i}{\partial x_j} \right|^2 dx \ge \beta \|\mathbf{v}\|_{H^1(\Omega)}^2, \quad \forall \mathbf{v} \in \mathbf{V}.$$

The previous boundary value problem leads to the following variational inequality (see, e.g., [15, 16])

$$\mathbf{u} \in \mathbf{V} : a(\mathbf{u}, \mathbf{v} - \mathbf{u}) + i_N(\mathbf{u}, \mathbf{v} - \mathbf{u}) + i_T(\mathbf{u}, \mathbf{v}) - i_T(\mathbf{u}, \mathbf{u}) > L(\mathbf{v} - \mathbf{u}), \ \forall \ \mathbf{v} \in \mathbf{V}$$
 (4)

where the linear form L is given by $L(\mathbf{v}) = \int_{\Omega} \mathbf{f}(x) \cdot \mathbf{v}(x) dx$ and

$$j_N(\mathbf{u}, \mathbf{v}) = \int_{\Gamma_C} c_N(u_N)_+^{m_N} v_N \, d\gamma(x), \quad j_T(\mathbf{u}, \mathbf{v}) = \int_{\Gamma_C} c_T(u_N)_+^{m_T} |\mathbf{v}_T| \, d\gamma(x).$$

We assume that

$$1 \le m_N, m_T \text{ if } n = 2, \ 1 \le m_N, m_T < 3 \text{ if } n = 3.$$
 (5)

The hypothesis (5) allows us to use the imbedding

$$H^1(\Omega) \hookrightarrow L^q(\Gamma_C)$$

for each $q \in [1, +\infty[$ if n = 2 and for each $q \in [1, 4]$ if n = 3. If n = 2 we have for each $\mathbf{u}, \mathbf{v}, \mathbf{v}_1, \mathbf{v}_2 \in (H^1(\Omega))^2$

$$|j_N(\mathbf{u}, \mathbf{v})| \le \|c_N\|_{L^{\infty}(\Gamma_C)} \|\mathbf{u}\|_{L^{2m_N}(\Gamma_C)}^{m_N} \|\mathbf{v}\|_{L^2(\Gamma_C)} \le C \|\mathbf{u}\|_{H^1(\Omega)}^{m_N} \|\mathbf{v}\|_{H^1(\Omega)}$$
 (6)

and

$$|j_{T}(\mathbf{u}, \mathbf{v}_{1}) - j_{T}(\mathbf{u}, \mathbf{v}_{2})| \leq \int_{\Gamma_{C}} c_{T}(u_{N})_{+}^{m_{T}} |(\mathbf{v}_{1} - \mathbf{v}_{2})_{T}| \, d\gamma(x)$$

$$\leq \|c_{T}\|_{L^{\infty}(\Gamma_{C})} \|\mathbf{u}\|_{L^{2m_{T}}(\Gamma_{C})}^{m_{T}} \|\mathbf{v}_{1} - \mathbf{v}_{2}\|_{L^{2}(\Gamma_{C})}$$

$$\leq C \|\mathbf{u}\|_{H^{1}(\Omega)}^{m_{T}} \|\mathbf{v}_{1} - \mathbf{v}_{2}\|_{H^{1}(\Omega)}.$$

$$(7)$$

If n=3 we obtain similar bounds. Indeed, consider $p \in [1,4]$ such that $1/p + m_N/4 = 1$ (which is possible since $1 \le m_N < 3$) and by Hölder inequality one gets

$$|j_N(\mathbf{u}, \mathbf{v})| \le C \|\mathbf{u}\|_{L^4(\Gamma_C)}^{m_N} \|\mathbf{v}\|_{L^p(\Gamma_C)} \le C \|\mathbf{u}\|_{H^1(\Omega)}^{m_N} \|\mathbf{v}\|_{H^1(\Omega)}$$
(8)

and similarly

$$|j_T(\mathbf{u}, \mathbf{v}_1) - j_T(\mathbf{u}, \mathbf{v}_2)| \le C \|\mathbf{u}\|_{L^4(\Gamma_C)}^{m_T} \|\mathbf{v}_1 - \mathbf{v}_2\|_{L^p(\Gamma_C)} \le C \|\mathbf{u}\|_{H^1(\Omega)}^{m_T} \|\mathbf{v}_1 - \mathbf{v}_2\|_{H^1(\Omega)}.$$
(9)

Let us analyze now the continuity of the applications j_N, j_T with respect to the first argument. If $m_N, m_T = 1$ it is easily seen that

$$\max\{|j_{N}(\mathbf{u}_{1},\mathbf{v})-j_{N}(\mathbf{u}_{2},\mathbf{v})|,|j_{T}(\mathbf{u}_{1},\mathbf{v})-j_{T}(\mathbf{u}_{2},\mathbf{v})|\} \leq C\|\mathbf{u}_{1}-\mathbf{u}_{2}\|_{L^{2}(\Gamma_{C})}\|\mathbf{v}\|_{L^{2}(\Gamma_{C})}$$

$$\leq C\|\mathbf{u}_{1}-\mathbf{u}_{2}\|_{H^{1}(\Omega)}\|\mathbf{v}\|_{H^{1}(\Omega)}.$$

When $m_N, m_T > 1$ we use the inequality

$$|(a)_{+}^{m} - (b)_{+}^{m}| \le m|(a)_{+} - (b)_{+}|((a)_{+}^{m-1} + (b)_{+}^{m-1})$$

 $\le m|a - b|(|a|^{m-1} + |b|^{m-1}) \quad a, b \in \mathbb{R}, m \ge 1.$

If n=2 we have for each $\mathbf{u}_1,\mathbf{u}_2,\mathbf{v}\in (H^1(\Omega))^2$

$$|j_{N}(\mathbf{u}_{1}, \mathbf{v}) - j_{N}(\mathbf{u}_{2}, \mathbf{v})|$$

$$\leq C \int_{\Gamma_{C}} |\mathbf{u}_{1} - \mathbf{u}_{2}| \left(|\mathbf{u}_{1}|^{m_{N}-1} + |\mathbf{u}_{2}|^{m_{N}-1} \right) |\mathbf{v}| \, d\gamma(x)$$

$$\leq C \|\mathbf{u}_{1} - \mathbf{u}_{2}\|_{L^{2}(\Gamma_{C})} \left\{ \|\mathbf{u}_{1}\|_{L^{4(m_{N}-1)}(\Gamma_{C})}^{m_{N}-1} + \|\mathbf{u}_{2}\|_{L^{4(m_{N}-1)}(\Gamma_{C})}^{m_{N}-1} \right\} \|\mathbf{v}\|_{L^{4}(\Gamma_{C})}$$

$$\leq C \|\mathbf{u}_{1} - \mathbf{u}_{2}\|_{H^{1}(\Omega)} \left\{ \|\mathbf{u}_{1}\|_{H^{1}(\Omega)}^{m_{N}-1} + \|\mathbf{u}_{2}\|_{H^{1}(\Omega)}^{m_{N}-1} \right\} \|\mathbf{v}\|_{H^{1}(\Omega)}.$$

$$(10)$$

If n = 3 we consider as before $p \in [1, 4]$ such that $1/4 + (m_N - 1)/4 + 1/p = 1$ and therefore one gets

$$|j_{N}(\mathbf{u}_{1}, \mathbf{v}) - j_{N}(\mathbf{u}_{2}, \mathbf{v})| \leq C \|\mathbf{u}_{1} - \mathbf{u}_{2}\|_{L^{4}(\Gamma_{C})} \left\{ \|\mathbf{u}_{1}\|_{L^{4}(\Gamma_{C})}^{m_{N}-1} + \|\mathbf{u}_{2}\|_{L^{4}(\Gamma_{C})}^{m_{N}-1} \right\} \|\mathbf{v}\|_{L^{p}(\Gamma_{C})}$$

$$\leq C \|\mathbf{u}_{1} - \mathbf{u}_{2}\|_{H^{1}(\Omega)} \left\{ \|\mathbf{u}_{1}\|_{H^{1}(\Omega)}^{m_{N}-1} + \|\mathbf{u}_{2}\|_{H^{1}(\Omega)}^{m_{N}-1} \right\} \|\mathbf{v}\|_{H^{1}(\Omega)}. (11)$$

Similar inequalities hold for the application j_T .

Remark 2.1 The above computations show that for each $\mathbf{u}, \mathbf{v} \in \mathbf{V}$ the functions $j_N(\mathbf{u}, \cdot)$, $j_T(\mathbf{u}, \cdot)$ and $j_N(\cdot, \mathbf{v})$, $j_T(\cdot, \mathbf{v})$ are continuous with respect to the strong topology of \mathbf{V} .

Remark 2.2 It is easily seen that the functions j_N, j_T are strongly continuous on $\mathbf{V} \times \mathbf{V}$

$$\lim_{k \to +\infty} (j_N, j_T)(\mathbf{u}_k, \mathbf{v}_k) = (j_N, j_T)(\mathbf{u}, \mathbf{v})$$

for each sequences $(\mathbf{u}_k, \mathbf{v}_k)_k$ converging strongly in $\mathbf{V} \times \mathbf{V}$ towards (\mathbf{u}, \mathbf{v}) .

Under the above conditions it has been proved in [15] that the problem (4) has solutions. Moreover, uniqueness results are available for sufficiently small coefficients c_N, c_T . Therefore, for $\varepsilon > 0$ small enough, the problems associated to the coefficients $\varepsilon c_N, \varepsilon c_T$ are well posed:

$$\mathbf{u}_{\varepsilon} \in \mathbf{V} : a(\mathbf{u}_{\varepsilon}, \mathbf{v} - \mathbf{u}_{\varepsilon}) + \varepsilon j_{N}(\mathbf{u}_{\varepsilon}, \mathbf{v} - \mathbf{u}_{\varepsilon}) + \varepsilon j_{T}(\mathbf{u}_{\varepsilon}, \mathbf{v}) - \varepsilon j_{T}(\mathbf{u}_{\varepsilon}, \mathbf{u}_{\varepsilon}) \ge L(\mathbf{v} - \mathbf{u}_{\varepsilon}), (12)$$

for all $\mathbf{v} \in \mathbf{V}$. It is easily seen by taking $\mathbf{v} = \mathbf{0}$ and $\mathbf{v} = 2\mathbf{u}_{\varepsilon}$ in (12) that the previous problem reduces to

$$a(\mathbf{u}_{\varepsilon}, \mathbf{u}_{\varepsilon}) + \varepsilon j_{N}(\mathbf{u}_{\varepsilon}, \mathbf{u}_{\varepsilon}) + \varepsilon j_{T}(\mathbf{u}_{\varepsilon}, \mathbf{u}_{\varepsilon}) = L(\mathbf{u}_{\varepsilon})$$
(13)

$$a(\mathbf{u}_{\varepsilon}, \mathbf{v}) + \varepsilon j_N(\mathbf{u}_{\varepsilon}, \mathbf{v}) + \varepsilon j_T(\mathbf{u}_{\varepsilon}, \mathbf{v}) \ge L(\mathbf{v}), \ \forall \ \mathbf{v} \in \mathbf{V}.$$
 (14)

Observing that $j_N(\mathbf{v}, \mathbf{v}) \geq 0$ for each $\mathbf{v} \in \mathbf{V}$ and $j_T \geq 0$ we deduce that $\sup_{\varepsilon > 0} \|\mathbf{u}_{\varepsilon}\|_{H^1(\Omega)} \leq \beta^{-1} \|\mathbf{f}\|_{L^2(\Omega)}$ since for each $\varepsilon > 0$ we have

$$\beta \|\mathbf{u}_{\varepsilon}\|_{H^{1}(\Omega)}^{2} \leq a(\mathbf{u}_{\varepsilon}, \mathbf{u}_{\varepsilon}) \leq L(\mathbf{u}_{\varepsilon}) \leq \|\mathbf{f}\|_{L^{2}(\Omega)} \|\mathbf{u}_{\varepsilon}\|_{L^{2}(\Omega)}.$$

Let us consider \mathbf{u}_0 the solution of the variational equality

$$\mathbf{u}_0 \in \mathbf{V} : a(\mathbf{u}_0, \mathbf{v}) = L(\mathbf{v}), \ \forall \ \mathbf{v} \in \mathbf{V}.$$
 (15)

Combining (12) written with $\mathbf{v} = \mathbf{u}_0$ and (15) written with $\mathbf{v} = \mathbf{u}_{\varepsilon} - \mathbf{u}_0$ implies

$$\beta \|\mathbf{u}_{\varepsilon} - \mathbf{u}_{0}\|_{H^{1}(\Omega)}^{2} \leq \varepsilon j_{N}(\mathbf{u}_{\varepsilon}, \mathbf{u}_{0} - \mathbf{u}_{\varepsilon}) + \varepsilon j_{T}(\mathbf{u}_{\varepsilon}, \mathbf{u}_{0}) - \varepsilon j_{T}(\mathbf{u}_{\varepsilon}, \mathbf{u}_{\varepsilon}).$$
 (16)

By the inequalities (6), (8) we have

$$|j_N(\mathbf{u}_{\varepsilon}, \mathbf{u}_0 - \mathbf{u}_{\varepsilon})| \le C \|\mathbf{u}_{\varepsilon}\|_{H^1(\Omega)}^{m_N} \|\mathbf{u}_{\varepsilon} - \mathbf{u}_0\|_{H^1(\Omega)} \le C \|\mathbf{u}_{\varepsilon} - \mathbf{u}_0\|_{H^1(\Omega)}. \tag{17}$$

Similarly (7), (9) imply

$$|j_T(\mathbf{u}_{\varepsilon}, \mathbf{u}_0) - j_T(\mathbf{u}_{\varepsilon}, \mathbf{u}_{\varepsilon})| \le C \|\mathbf{u}_{\varepsilon}\|_{H^1(\Omega)}^{m_T} \|\mathbf{u}_{\varepsilon} - \mathbf{u}_0\|_{H^1(\Omega)} \le C \|\mathbf{u}_{\varepsilon} - \mathbf{u}_0\|_{H^1(\Omega)}.$$
(18)

Therefore we deduce by (16), (17), (18) that

$$\sup_{\varepsilon > 0} \frac{\|\mathbf{u}_{\varepsilon} - \mathbf{u}_{0}\|_{H^{1}(\Omega)}}{\varepsilon} < +\infty. \tag{19}$$

In particular one obtains the strong convergence $\lim_{\varepsilon \searrow 0} \mathbf{u}_{\varepsilon} = \mathbf{u}_{0}$ in \mathbf{V} , saying that $\mathbf{u}_{\varepsilon} = \mathbf{u}_{0} + \mathcal{O}(\varepsilon)$. But in some applications the parameter $\varepsilon > 0$ is not too small and in such cases the above approximation is not satisfactory: we need to determine the first order correction \mathbf{r} such that $\mathbf{u}_{\varepsilon} = \mathbf{u}_{0} + \varepsilon \mathbf{r} + \mathcal{O}(\varepsilon^{2})$. In what follows we establish the strong convergence in \mathbf{V} of the fluctuations $\varepsilon^{-1}(\mathbf{u}_{\varepsilon} - \mathbf{u}_{0})$ towards \mathbf{r} and we characterize the first order correction \mathbf{r} as the unique solution of a variational inequality on some convex set. We denote by \mathcal{K} the set

$$\mathcal{K} = \{ \mathbf{q} \in \mathbf{V} : a(\mathbf{q}, \mathbf{v} - \mathbf{u}_0) + j_N(\mathbf{u}_0, \mathbf{v} - \mathbf{u}_0) + j_T(\mathbf{u}_0, \mathbf{v}) - j_T(\mathbf{u}_0, \mathbf{u}_0) \ge 0, \ \forall \ \mathbf{v} \in \mathbf{V} \}.$$

Observe that $j_N(\mathbf{u}_0,\cdot), j_T(\mathbf{u}_0,\cdot)$ are convex and therefore the set $-\mathcal{K}$ coincides with the subgradient of $j_N(\mathbf{u}_0,\cdot) + j_T(\mathbf{u}_0,\cdot)$ at \mathbf{u}_0 , with respect to the inner product $a(\cdot,\cdot)$. Therefore \mathcal{K} is a non empty closed convex set in \mathbf{V} .

Theorem 2.1 For each $\varepsilon > 0$ let $\mathbf{r}_{\varepsilon} = \varepsilon^{-1}(\mathbf{u}_{\varepsilon} - \mathbf{u}_{0})$ where \mathbf{u}_{ε} solves (12). Then $\lim_{\varepsilon \searrow 0} \mathbf{r}_{\varepsilon} = \mathbf{r}$ strongly in \mathbf{V} where \mathbf{r} is the unique solution of the variational inequality

$$\mathbf{r} \in \mathcal{K} : a(\mathbf{r}, \mathbf{q} - \mathbf{r}) \ge 0, \ \forall \ \mathbf{q} \in \mathcal{K}.$$
 (20)

Proof. By (19) we know that $(\mathbf{r}_{\varepsilon})_{\varepsilon>0}$ is bounded in **V** and thus there exists $\mathbf{r} \in \mathbf{V}$ and a sequence $(\varepsilon_k)_k$ converging towards zero such that $(\mathbf{r}_k)_k := (\mathbf{r}_{\varepsilon_k})_k$ converges weakly towards \mathbf{r} in **V**. Combining (14) and (15) one gets for each $\mathbf{v} \in \mathbf{V}$

$$a(\mathbf{u}_{\varepsilon}, \mathbf{v}) + \varepsilon j_N(\mathbf{u}_{\varepsilon}, \mathbf{v}) + \varepsilon j_T(\mathbf{u}_{\varepsilon}, \mathbf{v}) > L(\mathbf{v}) = a(\mathbf{u}_0, \mathbf{v})$$

and therefore we obtain for each k, with the notation $\mathbf{u}_k = \mathbf{u}_{\varepsilon_k}$

$$a(\mathbf{r}_k, \mathbf{v}) + j_N(\mathbf{u}_k, \mathbf{v}) + j_T(\mathbf{u}_k, \mathbf{v}) \ge 0.$$

Passing to the limit for $k \to +\infty$ we deduce by Remark 2.1 and by using the weak convergence of $(\mathbf{r}_k)_k$ that

$$a(\mathbf{r}, \mathbf{v}) + j_N(\mathbf{u}_0, \mathbf{v}) + j_T(\mathbf{u}_0, \mathbf{v}) \ge 0, \quad \mathbf{v} \in \mathbf{V}.$$
 (21)

Combining now (13) and (15) we obtain

$$a(\mathbf{u}_{\varepsilon}, \mathbf{u}_{\varepsilon}) + \varepsilon j_N(\mathbf{u}_{\varepsilon}, \mathbf{u}_{\varepsilon}) + \varepsilon j_T(\mathbf{u}_{\varepsilon}, \mathbf{u}_{\varepsilon}) = L(\mathbf{u}_{\varepsilon}) = a(\mathbf{u}_0, \mathbf{u}_{\varepsilon})$$

and therefore

$$a(\mathbf{r}_k, \mathbf{u}_k) + j_N(\mathbf{u}_k, \mathbf{u}_k) + j_T(\mathbf{u}_k, \mathbf{u}_k) = 0, \ \forall \ k.$$
 (22)

Since $(\mathbf{r}_k)_k$ converges weakly in \mathbf{V} and $(\mathbf{u}_k)_k$ converges strongly in \mathbf{V} we deduce that $\lim_{k\to+\infty} a(\mathbf{r}_k,\mathbf{u}_k) = a(\mathbf{r},\mathbf{u}_0)$. Notice also that by Remark 2.2 we have

$$\lim_{k \to +\infty} (j_N, j_T)(\mathbf{u}_k, \mathbf{u}_k) = (j_N, j_T)(\mathbf{u}_0, \mathbf{u}_0).$$

Passing to the limit for $k \to +\infty$ in (22) yields

$$a(\mathbf{r}, \mathbf{u}_0) + j_N(\mathbf{u}_0, \mathbf{u}_0) + j_T(\mathbf{u}_0, \mathbf{u}_0) = 0.$$
(23)

We deduce by (21), (23) that

$$a(\mathbf{r}, \mathbf{v} - \mathbf{u}_0) + j_N(\mathbf{u}_0, \mathbf{v} - \mathbf{u}_0) + j_T(\mathbf{u}_0, \mathbf{v}) - j_T(\mathbf{u}_0, \mathbf{u}_0) \ge 0, \ \forall \mathbf{v} \in \mathbf{V}$$
 (24)

saying that \mathbf{r} belongs to the set \mathcal{K} . Using now (12) with $\mathbf{v} = \mathbf{u}_0$ and (15) implies

$$a(\mathbf{u}_{\varepsilon}, \mathbf{u}_{0} - \mathbf{u}_{\varepsilon}) + \varepsilon j_{N}(\mathbf{u}_{\varepsilon}, \mathbf{u}_{0} - \mathbf{u}_{\varepsilon}) + \varepsilon j_{T}(\mathbf{u}_{\varepsilon}, \mathbf{u}_{0}) - \varepsilon j_{T}(\mathbf{u}_{\varepsilon}, \mathbf{u}_{\varepsilon}) \ge L(\mathbf{u}_{0} - \mathbf{u}_{\varepsilon}) = a(\mathbf{u}_{0}, \mathbf{u}_{0} - \mathbf{u}_{\varepsilon})$$

and therefore, replacing ε by ε_k one gets after division by ε_k^2

$$a(\mathbf{r}_k, \mathbf{r}_k) \le -j_N(\mathbf{u}_k, \mathbf{r}_k) + \frac{1}{\varepsilon_k} \left(j_T(\mathbf{u}_k, \mathbf{u}_0) - j_T(\mathbf{u}_k, \mathbf{u}_k) \right), \ \forall \ k.$$
 (25)

Taking now $\mathbf{v} = \mathbf{u}_k$ in (24) yields

$$a(\mathbf{r}, \mathbf{r}_k) \ge -j_N(\mathbf{u}_0, \mathbf{r}_k) + \frac{1}{\varepsilon_k} \left(j_T(\mathbf{u}_0, \mathbf{u}_0) - j_T(\mathbf{u}_0, \mathbf{u}_k) \right), \ \forall \ k.$$
 (26)

Since the positive part function is non decreasing we have for each k

$$J_N^k := j_N(\mathbf{u}_k, \mathbf{r}_k) - j_N(\mathbf{u}_0, \mathbf{r}_k)$$

$$= \frac{1}{\varepsilon_k} \int_{\Gamma_C} c_N \left((u_{kN})_+^{m_N} - (u_{0N})_+^{m_N} \right) (u_{kN} - u_{0N}) \, d\gamma(x) \ge 0.$$
(27)

We compute now the limit $l = \lim_{k \to +\infty} J_T^k$ where

$$J_T^k := \frac{1}{\varepsilon_k} \left(j_T(\mathbf{u}_k, \mathbf{u}_k) - j_T(\mathbf{u}_k, \mathbf{u}_0) + j_T(\mathbf{u}_0, \mathbf{u}_0) - j_T(\mathbf{u}_0, \mathbf{u}_k) \right).$$

We have for each k

$$J_T^k = \int_{\Gamma_C} c_T(u_{kN})_+^{m_T} \frac{|\mathbf{u}_{kT}| - |\mathbf{u}_{0T}|}{\varepsilon_k} \, \mathrm{d}\gamma(x)$$

$$- \int_{\Gamma_C} c_T(u_{0N})_+^{m_T} \frac{|\mathbf{u}_{kT}| - |\mathbf{u}_{0T}|}{\varepsilon_k} \, \mathrm{d}\gamma(x)$$

$$= \int_{\Gamma_C} c_T \left((u_{kN})_+^{m_T} - (u_{0N})_+^{m_T} \right) \frac{|\mathbf{u}_{kT}| - |\mathbf{u}_{0T}|}{\varepsilon_k} \, \mathrm{d}\gamma(x)$$

and therefore we obtain as before (see for example (10), (11)) that

$$|J_T^k| \leq C \int_{\Gamma_C} \left| (u_{kN})_+^{m_T} - (u_{0N})_+^{m_T} \right| |\mathbf{r}_k| \, \mathrm{d}\gamma(x)$$

$$\leq C \|\mathbf{u}_k - \mathbf{u}_0\|_{H^1(\Omega)} \|\mathbf{r}_k\|_{H^1(\Omega)}$$

$$= C\varepsilon_k \|\mathbf{r}_k\|_{H^1(\Omega)}^2.$$

Since $(\mathbf{r}_k)_k$ is bounded in **V** we deduce that $\lim_{k\to+\infty} J_T^k = 0$. Combining (25), (26), (27) yields

$$a(\mathbf{r}, \mathbf{r}_k) - a(\mathbf{r}_k, \mathbf{r}_k) \ge J_N^k + J_T^k \ge J_T^k$$

and after passing to the limit for $k \to +\infty$ one gets

$$a(\mathbf{r}, \mathbf{r}) \ge \limsup_{k \to +\infty} (a(\mathbf{r}_k, \mathbf{r}_k) + J_T^k) = \limsup_{k \to +\infty} a(\mathbf{r}_k, \mathbf{r}_k).$$

Since $(\mathbf{r}_k)_k$ converges weakly towards \mathbf{r} in \mathbf{V} we have also

$$a(\mathbf{r}, \mathbf{r}) \leq \liminf_{k \to +\infty} a(\mathbf{r}_k, \mathbf{r}_k)$$

implying that $\lim_{k\to+\infty} a(\mathbf{r}_k, \mathbf{r}_k) = a(\mathbf{r}, \mathbf{r})$, which says that $(\mathbf{r}_k)_k$ converges strongly towards \mathbf{r} in \mathbf{V} . Take now any element $\mathbf{q} \in \mathcal{K}$. By the definition of \mathcal{K} (with $\mathbf{v} = \mathbf{u}_k$) we have

$$a(\mathbf{q}, \mathbf{r}_k) + j_N(\mathbf{u}_0, \mathbf{r}_k) + \frac{1}{\varepsilon_k} j_T(\mathbf{u}_0, \mathbf{u}_k) - \frac{1}{\varepsilon_k} j_T(\mathbf{u}_0, \mathbf{u}_0) \ge 0.$$

Using one more time (25) we deduce that

$$a(\mathbf{q}, \mathbf{r}_k) - a(\mathbf{r}_k, \mathbf{r}_k) \geq j_N(\mathbf{u}_k, \mathbf{r}_k) - j_N(\mathbf{u}_0, \mathbf{r}_k)$$

$$+ \frac{1}{\varepsilon_k} (j_T(\mathbf{u}_k, \mathbf{u}_k) - j_T(\mathbf{u}_k, \mathbf{u}_0) - j_T(\mathbf{u}_0, \mathbf{u}_k) + j_T(\mathbf{u}_0, \mathbf{u}_0))$$

$$= J_N^k + J_T^k \geq J_T^k, \ \forall \ k.$$

Passing to the limit for $k \to +\infty$ yields

$$a(\mathbf{q} - \mathbf{r}, \mathbf{r}) \ge 0, \ \mathbf{q} \in \mathcal{K}.$$

Therefore the sequence $(\mathbf{r}_k)_k$ converges strongly in **V** towards the solution of the variational inequality (20). Since this solution is unique, all the family $(\varepsilon^{-1}(\mathbf{u}_{\varepsilon} - \mathbf{u}_0))_{\varepsilon>0}$ converges strongly in **V** towards this solution, as ε goes to zero.

2.2 A simplified friction problem

Now we consider a simplified friction model introduced in [7] (see also [4, 8]) involving a variational inequality of the second kind. Let $\Omega \subset \mathbb{R}^n$, $n \geq 1$, be a bounded domain with a smooth boundary denoted Γ . Let $V = H^1(\Omega)$ and set:

$$u_{\varepsilon} \in V : \qquad a(u_{\varepsilon}, v - u_{\varepsilon}) + \varepsilon j(v) - \varepsilon j(u_{\varepsilon}) \ge L(v - u_{\varepsilon}), \qquad \forall \ v \in V,$$
 (28)

where ε is a nonnegative constant representing in some sense the friction coefficient (see subsection 2.3 hereafter) and

$$a(u,v) = \int_{\Omega} \nabla u \cdot \nabla v \, dx + \int_{\Omega} uv \, dx, \ \forall \ u, v \in V,$$

$$L \in V',$$

$$j(v) = \int_{\Gamma} |v(x)| \, d\gamma(x), \ \forall \ v \in V.$$

It is easy to check (when $L(v) = \int_{\Omega} fv$) that the strong formulation corresponding to (28) is $-\Delta u_{\varepsilon} + u_{\varepsilon} = f$ in Ω , $|\partial u_{\varepsilon}/\partial \mathbf{n}| \leq \varepsilon$ if $u_{\varepsilon} = 0$ on $\partial \Omega$ and $\partial u_{\varepsilon}/\partial \mathbf{n} = -\varepsilon u_{\varepsilon}/|u_{\varepsilon}|$ if $u_{\varepsilon} \neq 0$ on $\partial \Omega$. The inequality (28) is equivalent to the two following conditions:

$$u_{\varepsilon} \in V : \quad a(u_{\varepsilon}, v) + \varepsilon j(v) \ge L(v), \quad \forall v \in V,$$
 (29)

$$a(u_{\varepsilon}, u_{\varepsilon}) + \varepsilon j(u_{\varepsilon}) = L(u_{\varepsilon}). \tag{30}$$

It is easily seen that j is convex and continuous on V (thanks to the trace theorem). Thus, Lions-Stampacchia's theorem ensures the existence of a unique solution to (28), see, e.g., [7, 14].

The main aim of the present section is to study the asymptotic behavior with respect to ε for the solutions of (28). Let us first mention a straightforward property dealing with the behavior of the solution when the loads increase:

Proposition 2.1 Let $L_1, L_2 \in V'$ such that:

$$L_1(v) \ge L_2(v), \ \forall \ v \in V, \ v \ge 0.$$

Let $u_{1,\varepsilon}$ be the solution to (28) with $L = L_1$ and let $u_{2,\varepsilon}$ be the solution to (28) with $L = L_2$. Then

$$u_{1,\varepsilon} \geq u_{2,\varepsilon}$$
.

Proof. As here, ε is fixed, for simplicity we denote the solutions of (28) corresponding to L_1, L_2 by u_1, u_2 respectively. Taking $L = L_1$ and $v = u_1 + (u_1 - u_2)_-$ in (28), we get

$$a(u_1, (u_1 - u_2)_-) + \varepsilon j(u_1 + (u_1 - u_2)_-) - \varepsilon j(u_1) \ge L_1((u_1 - u_2)_-).$$
(31)

Similarly, with $L = L_2$ and $v = u_2 - (u_1 - u_2)_-$ in (28) we obtain,

$$a(u_2, -(u_1 - u_2)_-) + \varepsilon j(u_2 - (u_1 - u_2)_-) - \varepsilon j(u_2) \ge L_2(-(u_1 - u_2)_-). \tag{32}$$

Since

$$j(u_1 + (u_1 - u_2)_-) + j(u_2 - (u_1 - u_2)_-) - j(u_1) - j(u_2) = 0,$$

we get by adding (31) and (32):

$$a(u_1 - u_2, (u_1 - u_2)_-) > (L_1 - L_2)((u_1 - u_2)_-) > 0,$$

and thus

$$-a((u_1 - u_2)_-, (u_1 - u_2)_-) \ge 0,$$

which proves the announced result.

2.2.1 Asymptotic behavior for small friction coefficient

When $\varepsilon = 0$, the variational inequality (28) simply reduces to the equality

$$u_0 \in V : \qquad a(u_0, v) = L(v), \qquad \forall \ v \in V.$$
 (33)

Choosing $v = u_0$ in (28) and $v = u_{\varepsilon} - u_0$ in (33) we get

$$a(u_{\varepsilon}, u_0 - u_{\varepsilon}) + \varepsilon j(u_0) - \varepsilon j(u_{\varepsilon}) \ge L(u_0 - u_{\varepsilon}),$$

and

$$a(u_0, u_{\varepsilon} - u_0) = L(u_{\varepsilon} - u_0).$$

By addition, we obtain

$$||u_{\varepsilon} - u_{0}||_{H^{1}(\Omega)}^{2} \leq \varepsilon j(u_{0}) - \varepsilon j(u_{\varepsilon}) \leq \varepsilon j(u_{\varepsilon} - u_{0}) \leq \varepsilon |\Gamma|^{1/2} ||u_{\varepsilon} - u_{0}||_{L^{2}(\Gamma)}$$

$$\leq C(\Omega)\varepsilon ||u_{\varepsilon} - u_{0}||_{H^{1}(\Omega)}, \quad (34)$$

where Cauchy-Schwarz inequality and the trace theorem have been used, and where $|\Gamma|$ denotes the measure of Γ . Hence u_{ε} converges towards u_0 in V as ε vanishes. Now we prove the following theorem.

Theorem 2.2 Let

$$D = \{ z \in V : a(z, v - u_0) + j(v) - j(u_0) \ge 0, \forall v \in V \},$$

(-D is the nonempty closed convex set of subgradients of j at u_0 with respect to the inner product of V). Then, when ε goes to zero, $(u_{\varepsilon} - u_0)/\varepsilon$ converges strongly in V to the unique solution of

$$y \in D$$
: $a(y, z - y) \ge 0$, $\forall z \in D$. (35)

Proof. Let $y_{\varepsilon} = (u_{\varepsilon} - u_0)/\varepsilon$. From (34) we deduce that $\sup_{\varepsilon > 0} ||y_{\varepsilon}||_{H^1(\Omega)} < +\infty$. Hence there exist $y \in V$ and a sequence $(\varepsilon_k)_k \setminus 0$ such that $(y_k)_k := (y_{\varepsilon_k})_k$ weakly converges in V to y. From (29) and (33) we get that

$$a(y_{\varepsilon}, v) + j(v) \ge 0, \quad \forall v \in V.$$

Taking $\varepsilon = \varepsilon_k$ and passing to the limit, we obtain an inequality satisfied by y:

$$a(y,v) + j(v) \ge 0, \qquad \forall \ v \in V.$$
 (36)

Besides, subtracting (33) (with $v = u_{\varepsilon}$) from (30) yields

$$a(y_{\varepsilon}, u_{\varepsilon}) + j(u_{\varepsilon}) = 0.$$

We take $\varepsilon = \varepsilon_k$. From the strong convergence of $(u_{\varepsilon_k})_k$ and the weak convergence of $(y_{\varepsilon_k})_k$ we get:

$$a(y, u_0) + j(u_0) = 0. (37)$$

Combining (36) and (37) we deduce that y satisfies the following inequality:

$$a(y, v - u_0) + j(v) - j(u_0) \ge 0, \quad \forall v \in V,$$
 (38)

saying that $y \in D$. Now, from (33) and (28) we have

$$a(u_0, u_{\varepsilon} - u_0) = L(u_{\varepsilon} - u_0),$$

$$a(u_{\varepsilon}, u_0 - u_{\varepsilon}) + \varepsilon j(u_0) - \varepsilon j(u_{\varepsilon}) \ge L(u_0 - u_{\varepsilon}).$$

By addition and division by ε^2 we obtain

$$a(y_{\varepsilon}, y_{\varepsilon}) \le \varepsilon^{-1} \left(j(u_0) - j(u_{\varepsilon}) \right). \tag{39}$$

Taking now $v = u_{\varepsilon}$ in the limit problem (38) and dividing by ε , one gets

$$a(y, y_{\varepsilon}) \ge \varepsilon^{-1} \left(j(u_0) - j(u_{\varepsilon}) \right). \tag{40}$$

Putting together (39) and (40) written with $\varepsilon = \varepsilon_k$ we deduce

$$\limsup_{k \to +\infty} a(y_{\varepsilon_k}, y_{\varepsilon_k}) \le \limsup_{k \to +\infty} a(y, y_{\varepsilon_k}) = a(y, y) \le \liminf_{k \to +\infty} a(y_{\varepsilon_k}, y_{\varepsilon_k})$$

 $\limsup_{k\to +\infty} a(y_{\varepsilon_k},y_{\varepsilon_k}) \leq \limsup_{k\to +\infty} a(y,y_{\varepsilon_k}) = a(y,y) \leq \liminf_{k\to +\infty} a(y_{\varepsilon_k},y_{\varepsilon_k}).$ Thus $(a(y_{\varepsilon_k},y_{\varepsilon_k}))_k$ converges to a(y,y) and therefore $\lim_{k\to +\infty} \|y_{\varepsilon_k}-y\|_{H^1(\Omega)} = 0$. And finally let $z\in D$. Then, by the definition of D (with $v=u_\varepsilon$) we have

$$a(z, y_{\varepsilon}) \ge \varepsilon^{-1} \left(j(u_0) - j(u_{\varepsilon}) \right).$$

Together with (39) the last inequality implies

$$a(z, y_{\varepsilon}) \ge a(y_{\varepsilon}, y_{\varepsilon}).$$

Taking $\varepsilon = \varepsilon_k$ and passing to the limit for $k \to +\infty$ we obtain that y solves (35). And since (35) admits a unique solution, the proof is complete.

2.2.2 Asymptotic behavior for large friction coefficients

We consider now the behavior for the solutions of (28) when the friction coefficient goes to infinity. Let $f \in L^2(\Omega)$ and denote by u_{∞} the solution of the Dirichlet problem

$$\begin{cases}
-\Delta u_{\infty}(x) + u_{\infty}(x) = f(x), & x \in \Omega, \\
u_{\infty}(x) = 0, & x \in \Gamma.
\end{cases}$$

Proposition 2.2 Assume that f and Ω are such that $\frac{\partial u_{\infty}}{\partial n} \in L^{\infty}(\Gamma)$ where n represents the unitary outward normal to Ω . Consider the linear application $L \in V'$ given by $L(v) = \int_{\Omega} f(x)v(x) dx$ for each $v \in V$. Then for each $\varepsilon \geq \left\| \frac{\partial u_{\infty}}{\partial n} \right\|_{L^{\infty}(\Gamma)}$ the solution of (28) coincides with u_{∞} .

Remark 2.3 A similar blocking phenomenon (in the entire body and not only on the boundary) occurs in the case of a Bingham fluid when the external loads are small enough (which corresponds in a certain sense to a large yield limit, see [3]).

Proof. (of Proposition 2.2) Let $v \in V$. We have $u_{\infty} \in V$ and note that $j(u_{\infty}) = 0$, thus

$$\begin{split} &a(u_{\infty},v-u_{\infty})+\varepsilon j(v)-\varepsilon j(u_{\infty})\\ &=\int_{\Omega}\nabla u_{\infty}\cdot\nabla(v-u_{\infty})dx+\int_{\Omega}u_{\infty}(x)(v(x)-u_{\infty}(x))dx+\varepsilon j(v)\\ &=\int_{\Gamma}\frac{\partial u_{\infty}}{\partial n}(v(x)-u_{\infty}(x))d\gamma(x)+L(v-u_{\infty})+\varepsilon j(v)\\ &=\int_{\Gamma}\left(\frac{\partial u_{\infty}}{\partial n}\;v(x)+\varepsilon |v(x)|\right)d\gamma(x)+L(v-u_{\infty})\\ &\geq L(v-u_{\infty}). \end{split}$$

Hence, by the uniqueness of the solution, $u_{\varepsilon} = u_{\infty}$ for each $\varepsilon \geq \left\| \frac{\partial u_{\infty}}{\partial n} \right\|_{L^{\infty}(\Gamma)}$.

2.2.3 Monotonicity of the solution with respect to friction

In this subsection we prove that the solution of (28) increases when ε decreases, namely, that

$$0 < \varepsilon_1 \le \varepsilon_2 \Rightarrow u_{\varepsilon_1} \ge u_{\varepsilon_2}. \tag{41}$$

We prove this result for nonnegative solutions. Observing that the solution of (28) with L=0 is $u_{\varepsilon}=0$ and applying Proposition 2.1 we deduce that nonnegative solutions are obtained for instance if L is such that $L(v) \geq 0$ for each $v \in V$, $v \geq 0$.

Let us first consider the regularized problem

$$u_{\varepsilon}^{\delta} \in V : \qquad a(u_{\varepsilon}^{\delta}, v - u_{\varepsilon}^{\delta}) + \varepsilon j_{\delta}(v) - \varepsilon j_{\delta}(u_{\varepsilon}^{\delta}) \ge L(v - u_{\varepsilon}^{\delta}), \qquad \forall \ v \in V,$$
 (42)

where for each $v \in V$ and $\delta \geq 0$ the convex function j_{δ} is given by

$$j_{\delta}(v) = \int_{\Gamma} \sqrt{v(x)^2 + \delta^2} \ d\gamma(x).$$

The following proposition shows that it is sufficient to prove (41) for each regularized problem (42) with $\delta > 0$.

Proposition 2.3 For each fixed $\varepsilon > 0$, when $\delta \searrow 0$, $(u_{\varepsilon}^{\delta})_{\delta}$ strongly converges in V to u_{ε} , the unique solution to Problem (28).

Proof. Since ε is fixed in this proposition, we use the simplified notation $u^{\delta} := u^{\delta}_{\varepsilon}$ and $u := u_{\varepsilon}$. From (42) and (28) we have for each $\delta > 0$:

$$a(u^{\delta}, u - u^{\delta}) + \varepsilon j_{\delta}(u) - \varepsilon j_{\delta}(u^{\delta}) \ge L(u - u^{\delta}),$$

and

$$a(u, u^{\delta} - u) + \varepsilon j(u^{\delta}) - \varepsilon j(u) \ge L(u^{\delta} - u).$$

By adding the two previous inequalities we get

$$a(u - u^{\delta}, u - u^{\delta}) \leq \varepsilon j_{\delta}(u) - \varepsilon j_{\delta}(u^{\delta}) + \varepsilon j(u^{\delta}) - \varepsilon j(u)$$

$$\leq \varepsilon j_{\delta}(u) - \varepsilon j(u)$$

$$= \varepsilon \int_{\Gamma} \left(\frac{\delta^{2}}{\sqrt{u^{2} + \delta^{2}} + |u|} \right) d\gamma(x)$$

$$\leq \varepsilon \delta |\Gamma|,$$

where $|\Gamma|$ denotes the measure of Γ . Hence the result.

We prove now the monotonicity property of u_{ε}^{δ} with respect to ε for each fixed $\delta > 0$.

Proposition 2.4 Consider $\varepsilon_1, \varepsilon_2, \delta \in \mathbb{R}$ such that $0 < \varepsilon_1 \le \varepsilon_2, \delta > 0$. Assume that the solutions $u_{\varepsilon_1}^{\delta}, u_{\varepsilon_2}^{\delta}$ are nonnegative. Then we have the inequality

$$u_{\varepsilon_1}^{\delta} \ge u_{\varepsilon_2}^{\delta}$$
.

Proof. We use the notations $u_1 := u_{\varepsilon_1}^{\delta}$, $u_2 := u_{\varepsilon_2}^{\delta}$. Choosing $v = u_i + hw$ in (28), with $h \in \mathbb{R}$ and $w \in V$ we get for $i \in \{1, 2\}$

$$a(u_i, hw) + \varepsilon_i j_{\delta}(u_i + hw) - \varepsilon_i j_{\delta}(u_i) \ge L(hw).$$

For h > 0 we obtain

$$a(u_i, w) + \frac{\varepsilon_i}{h} \int_{\Gamma} \left(\sqrt{(u_i + hw)^2 + \delta^2} - \sqrt{u_i^2 + \delta^2} \right) d\gamma(x) \ge L(w).$$

Letting $h \setminus 0$ we deduce using Lebesgue's theorem

$$a(u_i, w) + \varepsilon_i \int_{\Gamma} \frac{u_i w}{\sqrt{u_i^2 + \delta^2}} d\gamma(x) \ge L(w).$$

Similarly, considering h < 0, one gets

$$a(u_i, w) + \varepsilon_i \int_{\Gamma} \frac{u_i w}{\sqrt{u_i^2 + \delta^2}} d\gamma(x) \le L(w).$$

Then we obtain for each $w \in V$ and $i \in \{1, 2\}$

$$a(u_i, w) + \varepsilon_i \int_{\Gamma} \frac{u_i w}{\sqrt{u_i^2 + \delta^2}} d\gamma(x) = L(w).$$

We introduce the increasing application $t \to \varphi(t) = \frac{t}{\sqrt{t^2 + \delta^2}}, t \in \mathbb{R}$. The above formulation can be written

$$a(u_i, w) + \varepsilon_i \int_{\Gamma} \varphi(u_i) w \ d\gamma(x) = L(w), \ \ w \in V, \ i \in \{1, 2\}.$$

Taking as test function $w = (u_1 - u_2)_-$ we obtain

$$a(u_1 - u_2, (u_1 - u_2)_-) + \int_{\Gamma} (\varepsilon_1 \varphi(u_1) - \varepsilon_2 \varphi(u_2))(u_1 - u_2)_- d\gamma(x) = 0,$$

implying that

$$a((u_{1} - u_{2})_{-}, (u_{1} - u_{2})_{-}) = \int_{\Gamma} (\varepsilon_{1} \varphi(u_{1}) - \varepsilon_{2} \varphi(u_{2}))(u_{1} - u_{2})_{-} d\gamma(x)$$

$$\leq \varepsilon_{2} \int_{\Gamma} (\varphi(u_{1}) - \varphi(u_{2}))(u_{1} - u_{2})_{-} d\gamma(x).$$

By the monotonicity of φ we deduce that $(\varphi(u_1) - \varphi(u_2))(u_1 - u_2) \le 0$ and therefore

$$a((u_1 - u_2)_-, (u_1 - u_2)_-) \le 0,$$

saying that $u_1 \geq u_2$.

Combining now the Propositions 2.3, 2.4 we obtain the following result

Proposition 2.5 Let $\varepsilon_1, \varepsilon_2 \in \mathbb{R}$ such that $0 < \varepsilon_1 \le \varepsilon_2$ and assume that $L \in V'$ verifies $L(v) \ge 0$ for each $v \in V, v \ge 0$. Then $u_{\varepsilon_1} \ge u_{\varepsilon_2}$ where u_{ε_1} and u_{ε_2} are the solutions of (28) corresponding to ε_1 and ε_2 , respectively.

Proof. Applying Proposition 2.4 we have $u_{\varepsilon_1}^{\delta} \geq u_{\varepsilon_2}^{\delta}$ for each $\delta > 0$. Passing to the limit when $\delta \searrow 0$ we obtain by Proposition 2.3 that $u_{\varepsilon_1} \geq u_{\varepsilon_2}$.

2.3 Application to Tresca's friction model

In this subsection we show that the result obtained in Theorem 2.2 can be extended to a more physically relevant case dealing with the elasticity operator and vector valued functions. More precisely, let us consider an elastic body Ω in \mathbb{R}^n with n=2 or n=3 where plane small strain assumptions are made. The regular boundary Γ of Ω consists in two nonoverlapping parts Γ_D and Γ_C of positive measures. The normal unitary outward vector on Γ is denoted by \mathbf{n} . In its initial stage, the body is in contact on Γ_C with a rigid foundation (the extension to two elastic bodies in contact can be easily made, at least for small strain models) and we suppose that the body slips on Γ_C with friction. The body is clamped on Γ_D and it is subjected to volume forces $\mathbf{f} \in (L^2(\Omega))^n$. We assume that the normal constraint $S = S(x) \geq 0$ is given on Γ_C and $\mu \geq 0$ denotes the friction coefficient between the body and the rigid foundation (this corresponds to a so-called bilateral contact model with Tresca friction).

The frictional contact problem in linear elasticity consists in finding the displacement field $\mathbf{u} := \mathbf{u}_{\mu} : \Omega \to \mathbb{R}^n$ verifying the equations and conditions (43),(44):

$$\begin{cases} \operatorname{\mathbf{div}} \boldsymbol{\sigma}(\mathbf{u}) + \mathbf{f} &= \mathbf{0} & \text{in } \Omega, \\ \mathbf{u} &= \mathbf{0} & \text{on } \Gamma_D. \end{cases}$$
(43)

The conditions describing bilateral contact with Tresca friction (using the same notation as in section 2.1) on Γ_C are:

$$\begin{cases} u_N = 0, \\ |\boldsymbol{\sigma}_T(\mathbf{u})| \le \mu S, & if \quad \mathbf{u}_T = \mathbf{0}, \\ \boldsymbol{\sigma}_T(\mathbf{u}) = -\mu S \frac{\mathbf{u}_T}{|\mathbf{u}_T|}, & if \quad \mathbf{u}_T \ne \mathbf{0}. \end{cases}$$
(44)

The variational inequality issued from (43),(44) is: find $u_{\mu} \in \mathbf{V}$ such that

$$\int_{\Omega} \boldsymbol{\sigma}(\mathbf{u}_{\mu}) : \boldsymbol{\varepsilon}(\mathbf{v} - \mathbf{u}_{\mu}) dx + \mu \int_{\Gamma_{C}} S|\mathbf{v}_{T}| d\gamma(x) - \mu \int_{\Gamma_{C}} S|(\mathbf{u}_{\mu})_{T}| d\gamma(x) \ge \int_{\Omega} \mathbf{f} \cdot (\mathbf{v} - \mathbf{u}_{\mu}) dx, \quad (45)$$

for all $\mathbf{v} \in \mathbf{V}$ where

$$\mathbf{V} = \left\{ \mathbf{v} \in (H^1(\Omega))^n : \mathbf{v} = \mathbf{0} \text{ on } \Gamma_D, \ v_N = 0 \text{ on } \Gamma_C \right\}.$$

For each nonnegative μ there is a unique solution \mathbf{u}_{μ} of (45) and the results in Theorem 2.2 hold straightforwardly in this case.

In addition it is easy to check that the results of Proposition 2.2 can be applied to the inequality (45) dealing with the elasticity operator. More precisely, let $\mathbf{f} \in (L^2(\Omega))^n$ and let \mathbf{u}_{∞} denote the solution of the problem

$$\begin{cases}
-\mathbf{div} \ \boldsymbol{\sigma}(\mathbf{u}_{\infty}) = \mathbf{f} & \text{in } \Omega, \\
\mathbf{u}_{\infty} = \mathbf{0} & \text{on } \Gamma.
\end{cases}$$

Then, if $\mu S(x) \geq |\sigma_T(\mathbf{u}_{\infty}(x))|$ a.e. $x \in \Gamma_C$ then $\mathbf{u}_{\mu} = \mathbf{u}_{\infty}$.

2.4 Some open questions

A first open question is concerned with the study of the widespread Coulomb friction model in linear elasticity. It is well known that this problem admits a solution if the friction coefficient is small enough (see e.g., [18, 12, 5, 6] and the references quoted therein). Note that uniqueness does not hold in the general case, at least for large friction coefficients, see [10, 11]. More recently a first uniqueness result has been obtained in [20] with the assumption that a "regular" solution exists (with some additional technical assumptions) and that the friction coefficient is sufficiently small. This uniqueness result could be the starting point of a asymptotic analysis similar to the one achieved in this work.

Another line of research would consist of obtaining the higher order corrections and searching $\mathbf{r}_1, ..., \mathbf{r}_{\alpha}$ such that $\mathbf{u}_{\varepsilon} = \mathbf{u}_0 + \varepsilon \mathbf{r}_1 + ... + \varepsilon^{\alpha} \mathbf{r}_{\alpha} + \mathcal{O}(\varepsilon^{\alpha+1})$ for the previous models.

In the problem with normal compliance, one could also study the asymptotic behavior in the case where only the friction term involving c_T vanishes and the normal contact term remains unchanged: the so-called lubricated contact (see [2] and the references therein). Note that this nontrivial study is not a direct consequence of the results obtained in this paper. The symmetrical asymptotic study (when c_N vanishes and c_T remains unchanged) corresponds to a foundation which becomes completely soft and disappears is also an open problem.

A large amount of open problems are obtained when trying to generalize all these results to the quasistatic case (see [9]) when the forces very slowly in time.

Finally let us mention that we are not able to extend the results dealing with comparison of solutions such as positivity and monotonicity to the vector valued case involving the elasticity operator.

3 The elastoplastic torsion problem

Now we consider a simple variational inequality of the first kind modelling the elastoplastic torsion of a cylinder.

3.1 Problem set-up

Let $\Omega \subset \mathbb{R}^n$, $n \geq 1$, be a bounded domain with smooth boundary Γ . We consider the variational inequality modelling the torsion of an infinitely long elastoplastic cylinder of cross section Ω and plasticity yield r > 0: find the stress potential u such that

$$u \in K_r$$
: $a(u, v - u) \ge L(v - u), \quad \forall v \in K_r,$ (46)

where $a: H_0^1(\Omega) \times H_0^1(\Omega) \to \mathbb{R}$ is the bilinear form given by:

$$a(u,v) = \int_{\Omega} \nabla u \cdot \nabla v \, dx, \ \forall \ u,v \in H_0^1(\Omega),$$

and $L \in H^{-1}(\Omega) = (H_0^1(\Omega))'$. The notation K_r represents the nonempty closed convex set of admissible stress potentials:

$$K_r = \left\{ v \in H_0^1(\Omega) : |\nabla v| \le r, \text{ a.e. in } \Omega \right\}.$$

From Stampacchia's theorem we deduce that Problem (46) admits a unique solution (see also, e.g., [1, 4, 7, 8, 14]). More generally, for $\varepsilon > 0$ let u_{ε} be the unique solution of the variational inequality:

$$u_{\varepsilon} \in K_{\varepsilon} : \quad a(u_{\varepsilon}, v - u_{\varepsilon}) \ge L(v - u_{\varepsilon}), \quad \forall v \in K_{\varepsilon},$$
 (47)

where

$$K_{\varepsilon} = \left\{ v \in H_0^1(\Omega) : |\nabla v| \le \varepsilon, \text{ a.e. in } \Omega \right\}.$$

When the solution of

$$u \in H_0^1(\Omega)$$
: $a(u,v) = L(v), \quad \forall v \in H_0^1(\Omega),$

lies in $W^{1,\infty}(\Omega)$ then it is easy to see that $u_{\varepsilon} = u$ for each $\varepsilon \geq \|\nabla u\|_{L^{\infty}(\Omega)}$. For $\varepsilon > 0$ we set

$$w_{\varepsilon} = \frac{u_{\varepsilon}}{\varepsilon}$$
.

Our aim is to study the behavior of w_{ε} as ε tends to zero. According to (47), we see that for each positive ε , w_{ε} is the unique solution of

$$w_{\varepsilon} \in K_1 : \qquad \varepsilon a(w_{\varepsilon}, v - w_{\varepsilon}) \ge L(v - w_{\varepsilon}), \qquad \forall \ v \in K_1.$$
 (48)

3.2 Asymptotic behavior for small plasticity yield

In this paragraph we establish the convergence of the family $(w_{\varepsilon})_{\varepsilon>0}$ as ε goes to zero. We introduce the problem

$$u \in K_1 : L(v - u) \le 0, \qquad \forall \ v \in K_1, \tag{49}$$

and we denote by C the set of all the solutions of the above problem

$$C = \{u \in K_1 : L(v - u) \le 0, \forall v \in K_1\}.$$

Proposition 3.1 The set C is nonempty, closed in $H_0^1(\Omega)$ and convex.

Proof. By definition K_1 is bounded in $H_0^1(\Omega)$ and therefore $\{L(v) : v \in K_1\}$ is bounded in \mathbb{R} . We denote by M the supremum

$$M = \sup\{L(v) : v \in K_1\},\$$

and let us consider $(v_n)_n \subset K_1$ a maximizing sequence for M

$$M - \frac{1}{n} \le L(v_n) \le M, \ \forall \ n \ge 1.$$
 (50)

After extraction of a subsequence we can assume that $v_{n_k} \rightharpoonup \tilde{v}$ weakly in $H_0^1(\Omega)$ and therefore $\lim_{k \to +\infty} L(v_{n_k}) = L(\tilde{v})$. Since K_1 is closed and convex, it is also weakly closed implying that \tilde{v} belongs to K_1 . Passing to the limit in (50) we obtain that $M = L(\tilde{v})$ and therefore $\tilde{v} \in C$, saying that C is nonempty. It is easily seen that C is closed and convex.

Theorem 3.1 The family $(w_{\varepsilon})_{\varepsilon>0}$ converges strongly in $H_0^1(\Omega)$ as ε goes to zero towards the element of minimal norm of the set C

$$\lim_{\varepsilon \searrow 0} w_{\varepsilon} = w, \text{ strongly in } H_0^1(\Omega),$$

where

$$w \in C : \|\nabla w\|_{L^{2}(\Omega)} = \min_{v \in C} \|\nabla v\|_{L^{2}(\Omega)}.$$
(51)

Remark 3.1 As L is linear and K_1 is not a strictly convex set (a convex set K is said to be strictly convex if for each u_1, u_2 in K, $u_1 \neq u_2$ and $\lambda \in]0, 1[$, $\lambda u_1 + (1 - \lambda)u_2$ lies in the interior of K) uniqueness to Problem (49) is not guaranteed in general. Let for instance $\Omega =]0, 1[$ and $L(v) = \int_0^{1/4} v(x) dx;$ then the set of solutions to (49) is precisely the set of $u \in K_1$ such that u(x) = x in [0, 1/4]. Amongst all these solutions, a calculation shows that the one which solves (51) is given by w(x) = (1-x)/3 in [1/4, 1].

Proof. (of Theorem 3.1) Since $|\nabla w_{\varepsilon}| \leq 1$ a.e. in Ω , $(\nabla w_{\varepsilon})_{\varepsilon}$ is bounded in $L^{2}(\Omega)$ implying that $(w_{\varepsilon})_{\varepsilon}$ is bounded in $H^{1}_{0}(\Omega)$. Hence, there exists a sequence $(\varepsilon_{k})_{k}$ converging to zero and $w \in H^{1}_{0}(\Omega)$, such that $(w_{k})_{k} := (w_{\varepsilon_{k}})_{k}$ weakly converges in $H^{1}_{0}(\Omega)$ to w. As K_{1} is closed and convex, $w \in K_{1}$. Taking $\varepsilon = \varepsilon_{k}$ in (48) and passing to the limit as $k \to +\infty$ we obtain that w solves (49). We claim that the sequence $(w_{k})_{k}$ converges strongly in $H^{1}_{0}(\Omega)$. As $a(\cdot, \cdot)$ is symmetric, an equivalent formulation for (48) is

$$w_{\varepsilon} \in K_1, \qquad \frac{\varepsilon}{2} a(w_{\varepsilon}, w_{\varepsilon}) - L(w_{\varepsilon}) \le \frac{\varepsilon}{2} a(v, v) - L(v), \qquad \forall v \in K_1.$$
 (52)

Taking $\varepsilon = \varepsilon_k$ and v = w in (52), we deduce from (49) that

$$\forall k \geq 0, \qquad \frac{\varepsilon_k}{2} a(w_k, w_k) - \frac{\varepsilon_k}{2} a(w, w) \leq L(w_k - w) \leq 0.$$

Hence, we obtain the inequality

$$\limsup_{k \to +\infty} a(w_k, w_k) \le a(w, w),$$

while the weak convergence of $(w_k)_k$ and the lower semicontinuity of $a(\cdot,\cdot)$ with respect to the weak topology of $H_0^1(\Omega)$ imply

$$a(w, w) \le \liminf_{k \to +\infty} a(w_k, w_k).$$

Therefore we have the convergence $a(w_k, w_k) \to a(w, w)$ as $k \to +\infty$. This in turn implies that $(a(w_k - w, w_k - w))_k$ converges to zero and thus the convergence of $(w_k)_k$ to w holds strongly in $H_0^1(\Omega)$.

Since generally (49) does not admit a unique solution (see Remark 3.1), we need to characterize w amongst the solutions of Problem (49). Let $z \in C$ be a solution of (49). Taking $\varepsilon = \varepsilon_k$ and v = z in (52), we obtain as before

$$\forall k \geq 0, \qquad a(w_k, w_k) - a(z, z) \leq \frac{2}{\varepsilon_k} L(w_k - z) \leq 0.$$

Hence, passing to the limit as $k \to +\infty$ yields

$$\forall z \in C, \quad a(w, w) < a(z, z).$$

In other words w is the unique solution of Problem (51). As w is unique, the whole family $(w_{\varepsilon})_{\varepsilon}$ converges to w as ε goes to zero. The proof is complete.

Let d_{Γ} denote the distance function with respect to Γ

$$d_{\Gamma}(x) = \inf_{y \in \Gamma} |x - y|, \ \forall \ x \in \overline{\Omega},$$

where $|\cdot|$ stands for the euclidian norm in \mathbb{R}^n . It is straightforward to check that $d_{\Gamma} \in K_1$. The following additional result holds:

Proposition 3.2 Let w be the solution of (51) and assume that the linear application L satisfies

$$L(v) \ge 0, \ \forall \ v \in H_0^1(\Omega), \quad v \ge 0. \tag{53}$$

Then we have

$$L(w-d_{\Gamma})=0.$$

In particular, if $L(v) = \int_{\Omega} f(x)v(x)dx$ for each $v \in H_0^1(\Omega)$ with $f \in L^2(\Omega)$ and f > 0, then $\lim_{\varepsilon \searrow 0} w_{\varepsilon} = d_{\Gamma}$ strongly in $H_0^1(\Omega)$. If in addition there is some positive constant f_0 such that $f \ge f_0$, then

$$\lim_{\varepsilon \searrow 0} \varepsilon^{-1} \| d_{\Gamma} - w_{\varepsilon} \|_{L^{1}(\Omega)} = 0.$$

Proof. To prove the first part of the proposition, it suffices to take $v = d_{\Gamma}$ in (49) to obtain

$$L(d_{\Gamma} - w) < 0.$$

As $d_{\Gamma}-v \geq 0$ for each $v \in K_1$, we deduce by (53) that $L(d_{\Gamma}-w) \geq 0$ and thus $L(d_{\Gamma}-w) = 0$. In particular if f > 0 then $w = d_{\Gamma}$. Assume now that $f \geq f_0 > 0$ for some constant f_0 . Choosing v = w in (48) one gets

$$\frac{f_0}{\varepsilon} \|d_{\Gamma} - w_{\varepsilon}\|_{L^1(\Omega)} \le \frac{1}{\varepsilon} L(d_{\Gamma} - w_{\varepsilon}) = \frac{1}{\varepsilon} L(w - w_{\varepsilon}) \le a(w_{\varepsilon}, w - w_{\varepsilon}),$$

and therefore $\lim_{\varepsilon \searrow 0} \varepsilon^{-1} ||d_{\Gamma} - w_{\varepsilon}||_{L^{1}(\Omega)} = 0.$

Remark 3.2 1. Let us mention that calculations similar to those in Proposition 3.2 were achieved in [7] to prove that the solution of (46) converges to d_{Γ} when f(= constant) tends to infinity.

2. The case f = constant corresponds to a physically meaningful model in which $w = d_{\Gamma}$ according to the previous proposition.

We end this section with a monotonicity property of $(w_{\varepsilon})_{{\varepsilon}>0}$ with respect to ${\varepsilon}$.

Proposition 3.3 Let $0 < \varepsilon_1 \le \varepsilon_2$ and $L \in H^{-1}(\Omega)$ satisfy $L(v) \ge 0$ for each $v \in H_0^1(\Omega)$, $v \ge 0$. Then we have the inequality

$$w_{\varepsilon_1} = \frac{u_{\varepsilon_1}}{\varepsilon_1} \ge \frac{u_{\varepsilon_2}}{\varepsilon_2} = w_{\varepsilon_2}.$$

Proof. We use the notations $w_1 := w_{\varepsilon_1}, w_2 := w_{\varepsilon_2}$. We intend to use the variational inequality (48) with $\varepsilon = \varepsilon_1$ and $v_1 = w_1 + (w_1 - w_2)_-$ (the notation (·)_ stands for the negative part). We need to check that v_1 belongs to K_1 . For this observe that

$$v_1 = w_1 \mathbf{1}_{\{w_1 \ge w_2\}} + w_2 \mathbf{1}_{\{w_1 < w_2\}},$$

and therefore $v_1 \in H_0^1(\Omega), |\nabla v_1| \leq \max\{|\nabla w_1|, |\nabla w_2|\} \leq 1$. We obtain

$$\varepsilon_1 a(w_1, (w_1 - w_2)_-) \ge L((w_1 - w_2)_-).$$
 (54)

Taking now $\varepsilon = \varepsilon_2$ and $v = w_2 - (w_1 - w_2)_-$ one has

$$\varepsilon_2 a(w_2, -(w_1 - w_2)_-) \ge -L((w_1 - w_2)_-).$$
 (55)

Combining (54), (55) yields

$$\varepsilon_1 \varepsilon_2 a(w_1 - w_2, (w_1 - w_2)_-) \ge (\varepsilon_2 - \varepsilon_1) L((w_1 - w_2)_-) \ge 0,$$

implying that $a((w_1 - w_2)_-, (w_1 - w_2)_-) \le 0$. Therefore we deduce that $(w_1 - w_2)_- = 0$ saying that

 $\frac{u_{\varepsilon_1}}{\varepsilon_1} = w_1 \ge w_2 = \frac{u_{\varepsilon_2}}{\varepsilon_2}.$

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