# ON SECOND ORDER DIFFERENTIAL EQUATIONS WITH ASYMPTOTICALLY SMALL DISSIPATION 

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Pour Alban, né le 27 mars 2008


#### Abstract

We investigate the asymptotic properties as $t \rightarrow \infty$ of the differential equation $$
\begin{equation*} \ddot{x}(t)+a(t) \dot{x}(t)+\nabla G(x(t))=0, \quad t \geq 0 \tag{S} \end{equation*}
$$ where $x(\cdot)$ is $\mathbb{R}$-valued, the map $a: \mathbb{R}_{+} \rightarrow \mathbb{R}_{+}$is non increasing, and $G: \mathbb{R} \rightarrow \mathbb{R}$ is a potential with locally Lipschitz continuous derivative. We identify conditions on the function $a(\cdot)$ that guarantee or exclude the convergence of solutions of this problem to points in $\operatorname{argmin} G$, in the case where $G$ is convex and $\operatorname{argmin} G$ is an interval. The condition $\int_{0}^{\infty} e^{-\int_{0}^{t} a(s) d s} d t<\infty$ is shown to be necessary for convergence of trajectories, and a slightly stronger condition is shown to be sufficient.


## 1. Introduction

In this note, we study the differential equation

$$
\begin{equation*}
\ddot{x}(t)+a(t) \dot{x}(t)+\nabla G(x(t))=0, \quad t \geq 0 \tag{S}
\end{equation*}
$$

where $x(\cdot)$ is $\mathbb{R}$-valued, the $\operatorname{map} G: \mathbb{R} \rightarrow \mathbb{R}$ is at least of class $\mathcal{C}^{1}$, and $a: \mathbb{R}_{+} \rightarrow \mathbb{R}_{+}$ is a non increasing function. In a previous paper [3], we studied this differential equation in a finite- or infinite-dimensional Hilbert space $\mathcal{H}$. We are interested in the case where $a(t) \rightarrow 0$ as $t \rightarrow \infty$. Broadly speaking, convergence of solutions can be expected if $a(t) \rightarrow 0$ sufficiently slowly. One of the questions left open in that paper was whether solutions converge to a limit if the property

$$
\begin{equation*}
\int_{0}^{\infty} e^{-\int_{0}^{t} a(s) d s} d t=\infty \tag{1}
\end{equation*}
$$

does not hold and if argmin $G$ consists of more than just one point. In this note, we give a positive answer to this question, in the one dimensional case.

## 2. Preliminary Facts

Throughout this paper, we will denote by $G: \mathbb{R} \rightarrow \mathbb{R}$ a $\mathcal{C}^{1}$ function for which the derivative $g=G^{\prime}$ is Lipschitz continuous, uniformly on bounded sets. The function $a: \mathbb{R}_{+} \rightarrow \mathbb{R}_{+}$will always be assumed to be continuous and non-increasing. We also define the energy

$$
\begin{equation*}
\mathcal{E}(t)=G(x(t))+\frac{1}{2}|\dot{x}(t)|^{2} . \tag{2}
\end{equation*}
$$

[^0]Here are some basic results for solutions of $(\mathcal{S})$ from [3].
For any $\left(x_{0}, x_{1}\right) \in \mathbb{R}^{2}$, the problem $(\mathcal{S})$ has a unique solution $x(\cdot) \in \mathcal{C}^{2}([0, T), \mathbb{R})$ satisfying $x(0)=x_{0}, \dot{x}(0)=x_{1}$ on some maximal time interval $[0, T) \subset[0, \infty)$. For every $t \in[0, T)$, the energy identity holds

$$
\begin{equation*}
\frac{d}{d t} \mathcal{E}(t)=-a(t)|\dot{x}(t)|^{2} \tag{3}
\end{equation*}
$$

If in addition $G$ is bounded from below, then

$$
\begin{equation*}
\int_{0}^{T} a(t)|\dot{x}(t)|^{2} d t<\infty \tag{4}
\end{equation*}
$$

and the solution exists for all $T>0$. If also $G(\xi) \rightarrow \infty$ as $|\xi| \rightarrow \infty$ (i.e. if $G$ is coercive), then all solutions to $(\mathcal{S})$ remain bounded together with their first and second derivatives for all $t>0$. The bound depends only on the initial data. If a solution $x$ to $(\mathcal{S})$ converges toward some $\bar{x} \in \mathbb{R}$, then $\lim _{t \rightarrow \infty} \dot{x}(t)=\lim _{t \rightarrow \infty} \ddot{x}(t)=0$ and $\nabla G(\bar{x})=0$. If $\int_{0}^{\infty} a(s) d s<\infty$ and $\operatorname{if} \inf G>-\infty$, then solutions $x(\cdot)$ of $(\mathcal{S})$ for which $(x(0), \dot{x}(0)) \notin \operatorname{argmin} G \times\{0\}$ cannot converge to a point in $\operatorname{argmin} G$.

For the remainder of this note we shall assume that $\operatorname{argmin} G \neq \varnothing$. Without loss of generality, we may assume that $\min _{\mathbb{R}} G=0$ and $G(0)=0$. If for some $\rho \in \mathbb{R}_{+}$

$$
\begin{equation*}
\forall x \in \mathbb{R}, \quad G(x)-G(z) \leq \rho G^{\prime}(x)(x-z) \tag{G}
\end{equation*}
$$

then it is possible to show that any solution $x$ to the differential equation $(\mathcal{S})$ satisfies

$$
\begin{equation*}
\int_{0}^{\infty} a(t) \mathcal{E}(t) d t<\infty \tag{5}
\end{equation*}
$$

Since $t \mapsto \mathcal{E}(t)$ is decreasing, this estimate implies that $\mathcal{E}(t) \rightarrow \min G=0$ as $t \rightarrow \infty$, provided that $\int_{0}^{\infty} a(t) d t=\infty$. If now $\operatorname{argmin} G=\{\bar{x}\}$ is a singleton, then trajectories must converge to $\bar{x}$ under fairly weak additional conditions. The reader is referred to [3] for details.

## 3. CONVEX POTENTIALS WITH NON-UNIQUE MINIMA

In this section, we investigate the convergence of the trajectories of $(\mathcal{S})$ when $\operatorname{argmin} G$ is not a singleton. While the previous discussion shows that $\int_{0}^{\infty} a(s) d s=$ $\infty$ is a necessary condition for trajectories to converge to a point in $\operatorname{argmin} G$, this condition is clearly not sufficient, as the particular case $G \equiv 0$ shows. In this case, the solution is given by

$$
x(t)=x(0)+\dot{x}(0) \int_{0}^{t} e^{-\int_{0}^{s} a(u) d u} d s
$$

and the solution $x$ converges if and only if (1) does not hold. Therefore it is natural to ask whether for a general potential $G$, the trajectory $x$ is convergent if this condition does not hold. The potential $G$ is assumed to have all the properties listed in the previous section. A general result of non-convergence of the trajectories under the condition (1) is shown in [3]. There, we assume that $G$ is coercive, $\inf _{\mathbb{R}} G=0$, $\operatorname{argmin} G=[\alpha, \beta]$ for some $\alpha<\beta$, and that $G$ is non-increasing on $(-\infty, \alpha]$ and non-decreasing on $[\beta, \infty)$. It is also assumed that $a$ satisfies condition (1). Then either a solution satisfies $(x(0), \dot{x}(0)) \in[\alpha, \beta] \times\{0\}$, or else the $\omega$ - limit set $\omega\left(x_{0}, \dot{x}_{0}\right)$ contains $[\alpha, \beta]$ and hence the trajectory $x$ does not converge.

We now ask if the converse assertion is true: do the trajectories $x$ of $(\mathcal{S})$ converge if (1) does not hold? We give a positive answer when the map $a$ satisfies the following stronger condition

$$
\begin{equation*}
\int_{0}^{\infty} e^{-\theta} \int_{0}^{s} a(u) d u d s<\infty \tag{6}
\end{equation*}
$$

for some $\theta \in(0,1)$.
Theorem 3.1. Let $G: \mathbb{R} \rightarrow \mathbb{R}$ be a convex function of class $\mathcal{C}^{1}$ such that $G^{\prime}$ is Lipschitz continuous on the bounded sets of $\mathbb{R}$. Assume that $\operatorname{argmin} G=[\alpha, \beta]$ with $\alpha<\beta$ and that there exists $\delta>0$ such that

$$
\forall \xi \in(-\infty, \alpha], \quad G^{\prime}(\xi) \leq 2 \delta(\xi-\alpha) \quad \text { and } \quad \forall \xi \in[\beta, \infty), \quad G^{\prime}(\xi) \geq 2 \delta(\xi-\beta)
$$

Let $a: \mathbb{R}_{+} \rightarrow \mathbb{R}_{+}$be a differentiable non increasing map such that $\lim _{t \rightarrow \infty} a(t)=0$ and such that condition (6) holds for some positive $\theta<1$. Then, for any solution non constant solution $x$ to the differential equation $(\mathcal{S}), \lim _{t \rightarrow \infty} x(t)$ exists.

Proof. We may assume without loss of generality that $\alpha=0, \beta=1$. The conditions on $G$ imply that it is coercive, hence $\lim _{t \rightarrow \infty} \mathcal{E}(t)=0$ and $|x(t)| \leq M$ for some $M>0$, for all $t \in \mathbb{R}_{+}$.

Define the set $\mathcal{T}=\{t \geq 0 \mid \dot{x}(t)=0\}$ of sign changes of $\dot{x}$. This set must be discrete, for if it had an accumulation point $t^{*}$, then $\dot{x}\left(t^{*}\right)=0$ and also $\ddot{x}\left(t^{*}\right)=0$ by Rolle's Theorem. Since then $\dot{x}\left(t^{*}\right)=\ddot{x}\left(t^{*}\right)=G^{\prime}\left(x\left(t^{*}\right)\right)=0, x$ would have to be the constant solution 0 , which yields a contradiction.

If $\mathcal{T}$ is a finite set, then $\dot{x}$ does not change sign for sufficiently large $t$, and the trajectory $x$ has a limit. Let us therefore assume that $\mathcal{T}=\left\{t_{n} \mid n \in \mathbb{N}\right\}$, where the $t_{n}$ are increasing and tend to $\infty$. We want to show that this is impossible. Observe that at each $t_{n}, \dot{x}$ must change its sign and $G^{\prime}\left(x\left(t_{n}\right)\right) \neq 0$, since otherwise also $\ddot{x}\left(t_{n}\right)=0$ and we would again have a stationary solution. Without loss of generality, we can assume that $\dot{x}(0)<0, x(0)<0$ and therefore $x\left(t_{0}\right)<0$. Since $G^{\prime}\left(x\left(t_{0}\right)\right)<0$, equation $(\mathcal{S})$ shows that $\ddot{x}\left(t_{0}\right)>0$, hence the map $\dot{x}$ is positive on $\left(t_{0}, t_{1}\right), x\left(t_{1}\right)>1, \dot{x}$ is negative on $\left(t_{1}, t_{2}\right)$, and so on.

The argument so far shows that $G^{\prime}(x(t))$ vanishes on a union of infinitely many disjoint closed intervals,

$$
\{t \mid 0 \leq x(t) \leq 1\}=\cup_{k \geq 0}\left[u_{2 k}, u_{2 k+1}\right]
$$

where $0<t_{0}<u_{0}$ and $u_{2 k-1}<t_{k}<u_{2 k}$ for $k=1,2, \ldots$ Let us observe that, for every $k \in \mathbb{N}$,

$$
1=\left|x\left(u_{2 k+1}\right)-x\left(u_{2 k}\right)\right|=\int_{u_{2 k}}^{u_{2 k+1}}|\dot{x}(t)| d t \leq\left|u_{2 k+1}-u_{2 k}\right| \max _{t \geq u_{2 k}}|\dot{x}(t)|
$$

Since $\lim _{t \rightarrow \infty} \dot{x}(t)=0$, we deduce that $\lim _{k \rightarrow \infty}\left|u_{2 k+1}-u_{2 k}\right|=\infty$.
We next observe that for $u_{2 k} \leq t \leq u_{2 k+1}$ the function $v=\dot{x}$ satisfies $\dot{v}(t)+$ $a(t) v(t)=0$ and hence

$$
\begin{equation*}
\forall t \in\left[u_{2 k}, u_{2 k+1}\right], \quad \dot{x}(t)=\dot{x}\left(u_{2 k}\right) e^{-\int_{u_{2 k}}^{t} a(\tau) d \tau} \tag{7}
\end{equation*}
$$

Claim 3.1. There is a constant $\gamma$ such that $u_{2 k+2}-u_{2 k+1} \leq \gamma$ for all $k \in \mathbb{N}$.

To show this claim, fix $k \in \mathbb{N}$ and assume that $t \in\left[u_{2 k+1}, u_{2 k+2}\right]$. Assume for now that $k$ is odd and thus $x(t) \leq 0$. Define the quantity $A(t)=\exp \left(\frac{1}{2} \int_{0}^{t} a(s) d s\right)$ and set $y(t)=A(t) x(t)$. Then $y$ is the solution of the differential equation

$$
\begin{equation*}
\ddot{y}(t)+A(t) G^{\prime}\left(\frac{y(t)}{A(t)}\right)-\left(\frac{a^{2}(t)}{4}+\frac{\dot{a}(t)}{2}\right) y(t)=0 \tag{8}
\end{equation*}
$$

and satisfies $y\left(u_{2 k+1}\right)=y\left(u_{2 k+2}\right)=0$ and $\dot{y}\left(u_{2 k+1}\right)=A\left(u_{2 k+1}\right) \dot{x}\left(u_{2 k+1}\right)<0$. Since the map $a$ converges to 0 , we can choose $k$ large enough so that $a(t)<2 \sqrt{\delta}$ for every $t \in\left[u_{2 k+1}, u_{2 k+2}\right]$. On the other hand, the assumption on $G^{\prime}$ shows that, for every $t \in\left[u_{2 k+1}, u_{2 k+2}\right]$,

$$
A(t) G^{\prime}\left(\frac{y(t)}{A(t)}\right) \leq 2 \delta y(t)
$$

Recalling finally that $\dot{a}(t) \leq 0$ for every $t \geq 0$, we deduce from (8) that

$$
\forall t \in\left[u_{2 k+1}, u_{2 k+2}\right], \quad \ddot{y}(t)+\delta y(t) \geq 0 .
$$

The unique solution $z$ of the differential equation $\ddot{z}(t)+\delta z(t)=0$ with the same initial conditions as $y$ has the first zero larger than $u_{2 k+1}$ at $u_{2 k+1}+\frac{\pi}{\sqrt{\delta}}$. By a standard comparison argument, we deduce that $y$ vanishes before $z$ does, hence

$$
u_{2 k+2} \leq u_{2 k+1}+\gamma, \quad \gamma=\frac{\pi}{\sqrt{\delta}} .
$$

The same argument applies if $k$ is even. This proves the claim.
Claim 3.2. There is a $k_{0} \in \mathbb{N}$ such that for $k \geq k_{0}$

$$
\left|\dot{x}\left(u_{2 k+2}\right)\right| \leq\left|\dot{x}\left(u_{2 k}\right)\right| e^{-\theta \int_{u_{2 k}}^{u_{2 k+2}} a(s) d s}
$$

where $\theta$ is as in (6).
To prove this, pick $k_{0}$ so large that for all $k \geq k_{0}$

$$
(1-\theta)\left(u_{2 k+2}-u_{2 k}\right) \geq \gamma \theta
$$

This is possible since $u_{2 k+2}-u_{2 k} \rightarrow \infty$ as $k \rightarrow \infty$. Since $a$ is non-increasing, this implies that

$$
\begin{aligned}
\theta \int_{u_{2 k+1}}^{u_{2 k+2}} a(\tau) d \tau & \leq \gamma \theta a\left(u_{2 k+1}\right) \leq(1-\theta)\left(u_{2 k+1}-u_{2 k}\right) a\left(u_{2 k+1}\right) \\
& \leq(1-\theta) \int_{u_{2 k}}^{u_{2 k+1}} a(\tau) d \tau
\end{aligned}
$$

and hence

$$
\theta \int_{u_{2 k}}^{u_{2 k+2}} a(\tau) d \tau \leq \int_{u_{2 k}}^{u_{2 k+1}} a(\tau) d \tau
$$

Then for $k \geq k_{0}$,

$$
\begin{aligned}
\left|\dot{x}\left(u_{2 k+2}\right)\right| & \leq\left|\dot{x}\left(u_{2 k+1}\right)\right|=\left|\dot{x}\left(u_{2 k}\right)\right| e^{-\int_{u_{2 k}}^{u_{2 k+1}} a(s) d s} \\
& \leq\left|\dot{x}\left(u_{2 k}\right)\right| e^{-\theta \int_{u_{2 k}}^{u_{2 k+2}} a(s) d s}
\end{aligned}
$$

proving the claim.

Claim 3.3. If the set $\mathcal{T}$ is unbounded, there must exist a constant $C$, depending on $\mathcal{T}$ and on $x(0), \dot{x}(0)$ such that for all $t \geq 0$

$$
\begin{equation*}
|\dot{x}(t)| \leq C e^{-\theta \int_{0}^{t} a(s) d s} \tag{9}
\end{equation*}
$$

By making sure that $C$ is sufficiently large, we only have to prove the estimate for $t \geq u_{2 k_{0}}$. First assume that $u_{2 k} \leq t \leq u_{2 k+1}$ for some $k$. Then from (7)

$$
|\dot{x}(t)| \leq\left|\dot{x}\left(u_{2 k}\right)\right| e^{-\int_{u_{2 k}}^{t} a(s) d s} \leq\left|\dot{x}\left(u_{2 k}\right)\right| e^{-\theta \int_{u_{2 k}}^{t} a(s) d s}
$$

Using induction, we deduce from Claim 3.2 that

$$
|\dot{x}(t)| \leq\left|\dot{x}\left(u_{2 k_{0}}\right)\right| e^{-\theta \int_{u_{2 k_{0}}}^{t} a(s) d s}=C_{1} e^{-\theta \int_{0}^{t} a(s) d s}
$$

with $C_{1}=\left|\dot{x}\left(u_{2 k_{0}}\right)\right| e^{\theta \int_{0}^{u_{2 k}}} a(s) d s$. Next consider the case where $u_{2 k+1}<t \leq u_{2 k+2}$ for some $k$. Then

$$
|\dot{x}(t)| \leq\left|\dot{x}\left(u_{2 k+1}\right)\right| \leq C_{1} e^{-\theta \int_{0}^{u_{2 k+1}} a(s) d s} \leq C_{1} e^{\theta \int_{u_{2 k+1}}^{u_{2 k+2}} a(\tau) d \tau} e^{-\theta \int_{0}^{t} a(s) d s} .
$$

Due to Claim 3.1, $e^{\theta \int_{u_{2 k+1}}^{u_{2 k+2}} a(\tau) d \tau} \leq C_{2}$ for all $k$, for some constant $C_{2}$. Estimate (9) now follows for $t \geq u_{2 k_{0}}$ with $C=C_{1} C_{2}$. By enlarging $C$ further, the estimate follows for all $t \geq 0$.

Let us now conclude the proof of the theorem. From assumption (6) and estimate (9), we derive that $\dot{x} \in L^{1}(0, \infty)$. Hence $\lim _{t \rightarrow \infty} x(t)$ exists, contradicting the initial assumption. Therefore $\lim _{t \rightarrow \infty} x(t)$ exists after all, and the theorem has been proved.

Remark 3.1. Note that the map $t \mapsto \frac{c}{t+1}$ with $c>1$ satisfies condition (6) for every $\theta \in\left(\frac{1}{c}, 1\right)$. In fact, if merely $a(t) \geq \frac{c}{t+1}$ for $t$ large enough for some $c>1$, then condition (6) is satisfied. Consider next the family of maps $a: \mathbb{R}_{+} \rightarrow \mathbb{R}_{+}$defined by

$$
a(t)=\frac{1}{t+1}+\frac{d}{(t+1) \ln (t+2)}
$$

for some $d>0$. It is immediate to check that condition (1) holds if and only if $d \in$ $(0,1]$. Thus non-stationary trajectories of $(\mathcal{S})$ do not converge when $d \in(0,1]$. But condition (6) is never satisfied, for any $\theta \in(0,1)$ and $d>0$, and the convergence of trajectories remains an open question. Thus there remains a "logarithmic" gap between the criteria for existence and non-existence of limits.

We conclude with some remarks on convergence results in dimension $n>1$. It is possible to extend the non-convergence result given at the beginning of this section to the case where the differential equation is given in a Hilbert space $\mathcal{H}$, see [3]. However, it is not clear how to prove that $\lim _{t \rightarrow \infty} x(t)$ exists, in a general Hilbert space $\mathcal{H}$ and for the case where $G$ is convex and $\operatorname{argmin} G$ is not a singleton. Since in this case $|\dot{x}(t)| \leq \sqrt{2 \mathcal{E}(t)}$, it appears natural to derive convergence results from suitable estimates for $\mathcal{E}(t)$. In [3], we give conditions that imply $\mathcal{E}(t) \leq D a(t)$ for all $t$, for some constant $D>0$. However, since we must also assume that $\int_{0}^{\infty} a(s) d s=\infty$, these estimates are not strong enough to guarantee the convergence of trajectories.

One could try to extend the proof of Theorem 3.1. Set $a_{1}(t)=a(t) \cdot \chi_{S}(x(t))$, where $\chi_{S}$ is the characteristic function of $S$, then $\frac{d}{d t} \mathcal{E}(t) \leq-2 a_{1}(t) \mathcal{E}(t)$, and hence $\mathcal{E}(t) \leq \mathcal{E}(0) e^{-2 \int_{0}^{t} a_{1}(s) d s}$. If the function $t \mapsto e^{-\int_{0}^{t} a_{1}(s) d s}$ can be shown to be in $L^{1}(0, \infty)$, it would follow that $|\dot{x}|$ is integrable, implying the convergence of trajectories. This works in the one-dimensional case since the behavior of trajectories is quite simple. However, if $\operatorname{dim} \mathcal{H}>1$, it is difficult to satisfy this property, since trajectories corresponding to $(\mathcal{S})$ can be expected to behave like trajectories of a billiard problem in $S=\operatorname{argmin} G$ for large times.

When the map $a$ is constant and positive, it is established in [1,2] that the trajectories of $(\mathcal{S})$ are weakly convergent if the potential $G: \mathcal{H} \rightarrow \mathbb{R}$ is convex and $\operatorname{argmin} G \neq \varnothing$, in an arbitrary Hilbert space $\mathcal{H}$. The key ingredient of the proof is the Opial lemma [4], which allows the authors of these papers to prove convergence even if $|\dot{x}(\cdot)|$ is only in $L^{2}(0, \infty)$ and not in $L^{1}(0, \infty)$. However, if e.g. $a(t)=\frac{c}{t+1}$, then Opial's lemma requires that we show $\int_{0}^{\infty}(t+1)|\dot{x}(t)|^{2} d t<\infty$, while (4) implies only $\int_{0}^{\infty} \frac{1}{t+1}|\dot{x}(t)|^{2} d t<\infty$. Hence there remains a gap if arguments similar to those in [1] or [2] are to be used. It is unclear how this gap can be closed.

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[^0]:    2000 Mathematics Subject Classification. 34G20, 34A12, 34D05.
    Key words and phrases. Differential equation, dissipative dynamical system, vanishing damping, asymptotic behavior.

