

A natural quantization for chaotic maps

General objective: relate spectral properties of classical chaos and quantum chaos.

F. Faure (Grenoble) with M. Tsujii (Kyushu univ.)

February 8, 2012

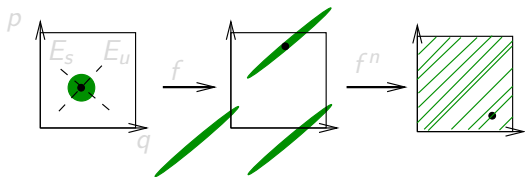
Anosov map

Let (M, ω) a **symplectic compact** manifold, $\dim M = 2d$,
 $f : M \rightarrow M$ a smooth **symplectic Anosov** diffeom.: there exists a Df -invariant decomposition $TM = E_s \oplus E_u$, $\lambda > 1$,

$$\|Df|_{E_s}\| \leq \frac{1}{\lambda}, \quad \|Df|_{E_u}^{-1}\| \leq \frac{1}{\lambda},$$

Examples

Arnold Cat map: $f_0 : \begin{pmatrix} q \\ p \end{pmatrix} \rightarrow \begin{pmatrix} 2 & 1 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} q \\ p \end{pmatrix}$ on $\mathbb{T}^2 = (\mathbb{R}/\mathbb{Z})^2$
and its perturbations $f = g \circ f_0$, $\|g - \text{Id}\|_{C^1} \ll 1$, preserving $\omega = dq \wedge dp$.



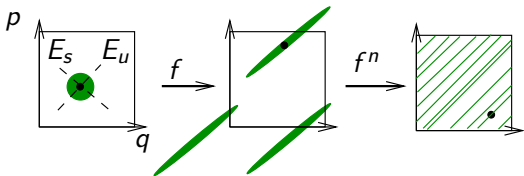
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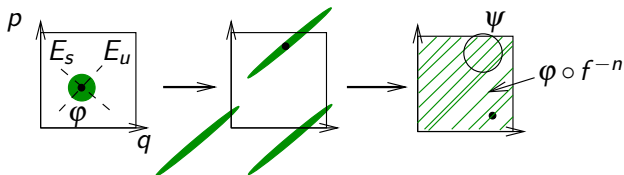
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Mixing



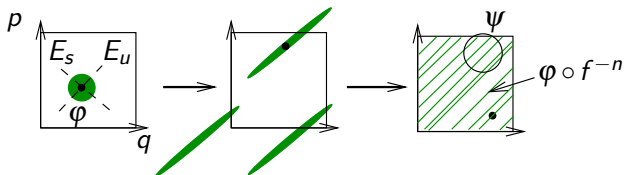
Theorem (Anosov, Ruelle,..)

f is **mixing** with exponential decay of correlations:

$$\exists \alpha < 1, \forall \psi, \varphi \in C^\infty(M), \quad \langle \psi, \varphi \circ f^{-n} \rangle_{L^2(M)} = \left(\int \psi \right) \left(\int \varphi \right) + O(\alpha^n)$$

- This implies Central limit theorem, random aspects, irreversible convergence towards equilibrium, ...
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Question of a natural quantization

Question (*): Existence of a sequence $\hbar = \hbar_j \xrightarrow{j \rightarrow \infty} 0$ and operators and spaces

$\hat{F}_\hbar: \mathcal{H}_\hbar \rightarrow \mathcal{H}_\hbar$, with $\dim \mathcal{H}_\hbar < \infty$, s.t.

① \hat{F}_\hbar is “almost unitary”:

$$\exists \varepsilon_\hbar \rightarrow 0, \forall \hbar, \forall u \in \mathcal{H}_\hbar, (1 - \varepsilon_\hbar) \|u\| \leq \|\hat{F}_\hbar u\| \leq (1 + \varepsilon_\hbar) \|u\|, ,$$

② “Gutzwiller formula”: , $\exists \theta < 1, \forall n, \forall \hbar$,

$$\text{Tr} \left(\hat{F}_\hbar^n \right) = \sum_{x=f^n(x)} \frac{e^{iS_{n,x}/\hbar}}{\sqrt{|\text{Det}(1 - Df_x^n)|}} + C_\hbar \theta^n$$

with action $S_{n,x} = \int pdq - Hdt$ defined later.

Proposition (Unicity)

If \hat{F}_\hbar exists, then its spectrum (with multiplicity) is **unique**. In particular $\dim \mathcal{H}_\hbar$ is unique.

- **proof**: use: if F, F' are matrices with $\forall n, |\text{Tr}(F^n) - \text{Tr}(F'^n)| < C\theta^n$ then F, F' have same spectrum on $|z| > \theta$. So above, $\theta < 1$ is important.

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Existence

What is known for Anosov non linear map $f : \mathbb{T}^2 \rightarrow \mathbb{T}^2$ (Keating, S. Debievre, 90'):

- **Weyl quantization** (or **geometric quantization**) gives $\hat{F}_\hbar : \mathcal{H}_\hbar \rightarrow \mathcal{H}_\hbar$ unitary, $\dim \mathcal{H}_\hbar = \frac{1}{2\pi\hbar}$,
- Gutzwiller formula has a remainder $\hbar\theta^n$, but $\theta = e^{h_0} > 1$, with **"topological entropy"** $h_0 > 0$.
- Exception: the linear cat map map f_0 , for which $\theta = 0$ (then answer to (*) is yes).

Theorem ((*) M.Tsujii, F.F.)

With assumptions:

① $[\omega] \in H^2(M, \mathbb{Z})$,

② 1 is not eigenvalue of the linear map $f_* : H_1(M, \mathbb{R}) \rightarrow H_1(M, \mathbb{R})$

then **answer to Question (*) is yes**, for $N = \frac{1}{2\pi\hbar} \in \mathbb{N}$ large enough. The construction uses **"geometric pre-quantization"** of f . We have

$$\dim \mathcal{H}_\hbar \sim \frac{1}{(2\pi\hbar)^d} \text{Vol}_\omega(M) \quad \left(= \int_M e^{\frac{\omega}{2\pi\hbar}} \text{Todd}(TM) \right)$$

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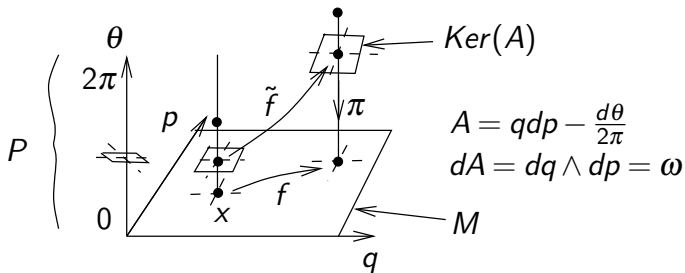
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Geometric pre-quantization of f

Theorem (Kostant, Souriau, Kirillov 70', Zelditch 05)

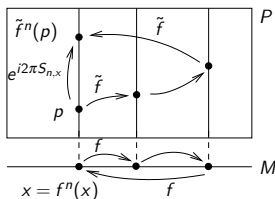
With assumptions 1,2 above, there exists (almost unique):

- 1 $U(1)$ -principal bundle $P \xrightarrow{\pi} M$ with connection 1-form A , curvature $dA = (\pi^* \omega)$
- 2 a “prequantum map” $\tilde{f} : P \rightarrow P$
with $f \circ \pi = \pi \circ \tilde{f}$, $\tilde{f}^* A = A$, $\tilde{f}(e^{i\theta} p) = e^{i\theta} \tilde{f}(p)$

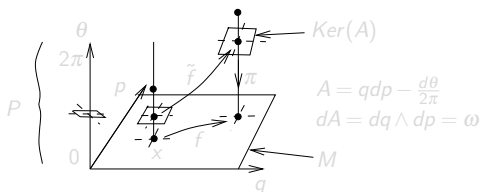


Remarks

- The **action** $e^{i2\pi S_{n,x}}$ of a periodic point is:



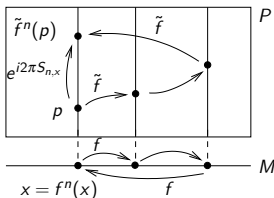
- \tilde{f} is a **partially hyperbolic map** with neutral direction θ preserving the contact 1-form A . It is not obvious that \tilde{f} is mixing with exp. decay of correlations.



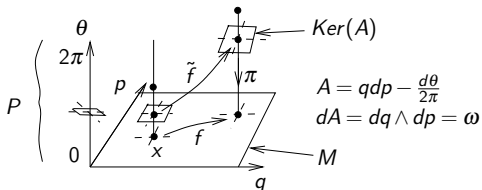
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Prequantum transfer operator

Definition

The **prequantum transfer operator** with **potential** $V \in C^\infty(M)$ is

$$\hat{F} : \begin{cases} C^\infty(P) & \rightarrow C^\infty(P) \\ u & \rightarrow e^{V \circ \pi} (u \circ \tilde{f}^{-1}) \end{cases}$$

- For every $N \in \mathbb{Z}$, \hat{F} preserves the space:

$$C_N^\infty(P) := \left\{ u \in C^\infty(P), \quad u(e^{i\theta} p) = e^{iN\theta} u(p) \right\},$$

called “**equivariant functions**” or “**Fourier mode N in the fiber**”.

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Spectrum of “Ruelle resonances”

Theorem (V.Baladi, M.Tsujii 07, N.Roy, J.Sjöstrand, F.F., 08)

For every $N \in \mathbb{Z}$, in **anisotropic Sobolev spaces** $H_N^m(P)$ of variable order $m(x, \xi)$, with $m(x, \xi) = \pm M$ along $E_{s/u}^*$, then

$$\hat{F}_N : H_N^m(P) \rightarrow H_N^m(P), \quad C_N^\infty(P) \subset H_N^m(P) \subset \mathcal{D}'_N(P)$$

is **bounded** and $r_{\text{ess.}}(\hat{F}_N) \xrightarrow{M \rightarrow \infty} 0$. The eigenvalues outside r_{ess} do not depend on $H_N^m(P)$ and are called **Ruelle resonances** $\text{Res}(\hat{F}_N)$.

Remark : the only obvious resonance is for the “equilibrium” $u = \text{cste}$, $\lambda = 1$ for $N = 0$, .



Spectrum in annuli

Theorem (M. Tsujii, F.F.)

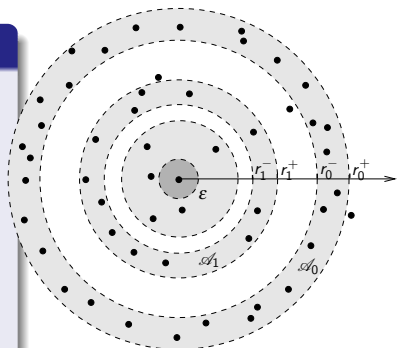
In the semiclassical limit $N = \frac{1}{2\pi\hbar} \rightarrow \infty$,

$$\text{Res}(\hat{F}_N) \subset \left\{ |z| \in \bigcup_{k \geq 0} \underbrace{[r_k^- - \delta_N, r_k^+ + \delta_N]}_{\text{annulus } \mathcal{A}_k} \right\}$$

with $\delta_N = \frac{C}{N^\varepsilon} \rightarrow 0$, $\varepsilon > 0, C > 0$ indept of f ,

$$r_k^- = \lim_{n \rightarrow \infty} \left| \inf_{x \in M} \left(e^{V_n(x)} \left(\underbrace{\|Df_{/E_u}^n(x)\|}_{>1} \right)^{-k} \left| \underbrace{\det Df_{/E_u}^n(x)}_{>1} \right|^{-1/2} \right) \right|^{1/n},$$

$$r_k^+ = \dots \sup_{x \in M} \dots \left(\|Df_{/E_u}^n(x)\|^{-1} \right)^{+k} \dots, \quad V_n(x) = \sum_{j=1}^n V(f^j(x))$$



Rem: $r_{k+1}^- < r_k^-$, $r_{k+1}^+ < r_k^+$. Choice $v = v_0 = \frac{1}{2} \log |\det Df_{/E_u}(x)|$ gives $e^{v_0, n} = |\det Df_{/E_u}^n(x)|^{1/2}$ and $r_0^- = r_0^+ = 1$.

Theorem

Let $\hat{F}_\hbar : \mathcal{H}_\hbar \rightarrow \mathcal{H}_\hbar$: spectral restriction of \hat{F}_N on \mathcal{A}_0 , then

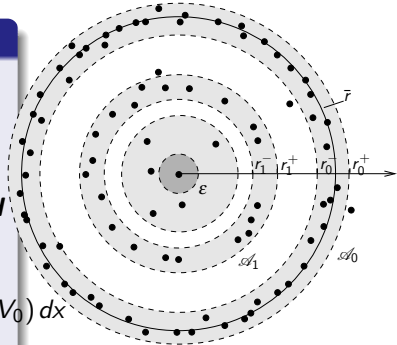
- 1 $\dim \mathcal{H}_\hbar = \frac{1}{(2\pi\hbar)} \text{Vol}(M) + o.$
- 2 Most of eigenvalues of \hat{F}_\hbar **concentrate and equidistribute** on the circle of radius

$$\bar{r} = e^{\langle V - V_0 \rangle}, \quad \langle V - V_0 \rangle := \frac{1}{\text{Vol}(M)} \int (V - V_0) dx$$

- 3 **Gutzwiller trace formula:**

$$\text{Tr}(\hat{F}_\hbar^n) = \sum_{x=f^n(x)} \frac{e^{(V-V_0)_n} e^{iS_{n,x}/\hbar}}{\sqrt{|\text{Det}(1 - Df_x^n)|}} + \hbar^\varepsilon \left(\frac{1}{\lambda}\right)^n$$

→ We get initial Thm (*) if $V = V_0$.



Corollary

The map $\tilde{f} : P \rightarrow P$ is **mixing** with **exponential decay of correlations**.

Heuristic: at large time, $\hat{F}_N^n \simeq \hat{F}_\hbar^n + O\left(\frac{1}{\lambda^n}\right)$, i.e. **“quantum dynamics emerges dynamically”**

Summary. Is it a “natural quantization”?

Construction:

- We start from $f : M \rightarrow M$ **symplectic and Anosov**.
- Then the **prequantum map** $\tilde{f} : P \rightarrow P$ and **prequantum operator** $\hat{F}_N : C_N^\infty(P) \rightarrow C_N^\infty(P)$ are (almost) uniquely defined.
Take the potential $V_0 = \frac{1}{2} \log |\det Df|_{E_u}(x)|$.
- We consider $\Pi_N : H_N^m(P) \rightarrow \mathcal{H}_\hbar$: spectral (finite rank) **projector on the external annulus**. Let $\hat{F}_\hbar = \Pi_N \hat{F}_N \Pi_N$. (set $\hbar = \frac{1}{2\pi N}$).

Properties:

- \hat{F}_\hbar is **almost unitary** (in the limit $\hbar \rightarrow 0$) and satisfies **Gutzwiller formula** with exp. small remainder in large time.
- If $a \in C^\infty(M)$, define $\text{Op}_\hbar(a) := \Pi_N a \Pi_N$. We have (obviously) **exact Egorov theorem**:

$$F_\hbar \text{Op}_\hbar(a) = \text{Op}_\hbar(a \circ f^{-1}) F_\hbar$$

(= $\Pi_N \hat{F}_N a \Pi_N = \Pi_N (a \circ f^{-1}) \hat{F}_N \Pi_N$), and **symbol calculus**:

$$\text{Op}_\hbar(a) \text{Op}_\hbar(b) = \text{Op}_\hbar(ab) + O(\hbar^\varepsilon) \quad (1)$$

$$\text{Op}_\hbar(\bar{a}) = (\text{Op}_\hbar(a))^\dagger + O(\hbar^\varepsilon) \quad (2)$$

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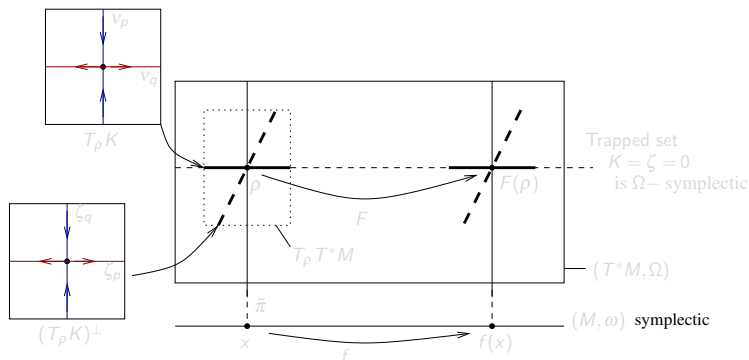
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- Atiyah-Bott, **exact trace formula** (66).

The operator $\hat{F}_N : C_N^\infty(P) \rightarrow C_N^\infty(P)$ is a F.I.O. with **canonical map** on (T^*M, Ω) .

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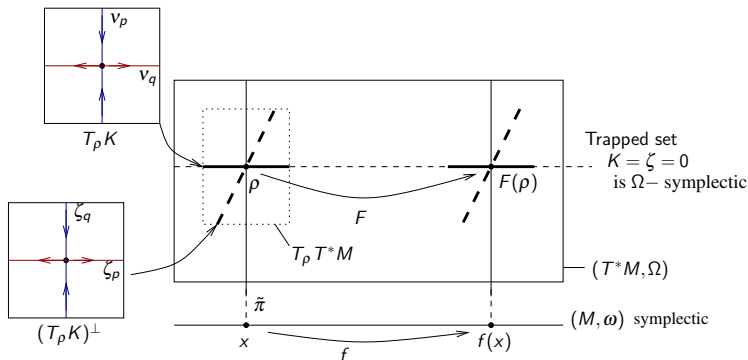


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Some key results for linear maps

- Let $f = \begin{pmatrix} A & 0 \\ 0 & {}_tA^{-1} \end{pmatrix} : \mathbb{R}^{2d} \rightarrow \mathbb{R}^{2d}$ symplectic, hyperbolic with $A : \mathbb{R}^d \rightarrow \mathbb{R}^d$ expanding.
- Let $\hat{F}_f u := u \circ f^{-1}$: unitary in $L^2(\mathbb{R}^{2d})$.
- Then

$$\hat{F}_f = \hat{F}_A \otimes \hat{F}_{{}_tA^{-1}}$$

with $(\hat{F}_A v) := \frac{1}{\sqrt{|\det A|}} v \circ A^{-1}$ unitary.

- In Sobolev space $H^{m(x,\xi)}$, \hat{F}_A has discrete spectrum (in polyn. space) and $r(\hat{F}_A) = \frac{1}{\sqrt{|\det A|}}$.
- Also (Atiyah-Bott): $\text{Tr}^b \hat{F}_f = \int_{\mathbb{R}^{2d}} \langle \delta_x | \hat{F}_f \delta_x \rangle = \int_{\mathbb{R}^{2d}} \delta(x - f^{-1}(x)) = \frac{1}{|\det(1-f)|}$
hence

$$\text{Tr}^b \hat{F}_A = \text{Tr}^b \hat{F}_{{}_tA^{-1}} = \frac{1}{\sqrt{|\det(1-f)|}},$$