# STRICHARTZ ESTIMATES FOR LONG RANGE PERTURBATIONS 

JEAN-MARC BOUCLET AND NIKOLAY TZVETKOV


#### Abstract

We study local in time Strichartz estimates for the Schrödinger equation associated to long range perturbations of the flat Laplacian on the euclidean space. We prove that in such a geometric situation, outside a large ball centered at the origin, the solutions of the Schrödinger equation enjoy the same Strichartz estimates as in the non perturbed situation. The proof is based on the IsozakiKitada parametrix construction. If in addition the metric is non trapping, we prove that the Strichartz estimates hold in the whole space.


## Contents

1. Introduction. ..... 1
2. Functional calculus ..... 5
3. The Isozaki-Kitada parametrix ..... 14
4. Strichartz estimates outside a large ball ..... 21
5. Semi-classical time estimates and applications ..... 26
6. Using the non trapping assumption ..... 31
7. Non homogeneous estimates and nonlinear applications ..... 34
References ..... 39

## 1. Introduction.

Let $(M, g)$ be a $d$-dimensional Riemannian manifold. Denote by $\Delta_{g}$ the LaplaceBeltrami operator associated to the metric $g$. Consider the time dependent Schrödinger equation on $(M, g)$

$$
\begin{equation*}
i u_{t}+\Delta_{g} u=0 \tag{1.1}
\end{equation*}
$$

subject to initial data

$$
\begin{equation*}
\left.u\right|_{t=0}=u_{0} . \tag{1.2}
\end{equation*}
$$

It is well-known (see e.g. [18]) that when $(M, g)$ is the flat Euclidean space (i.e. $\mathbb{R}^{d}$ with the metric $g_{i j}=\delta_{i j}$ the unit $d \times d$ matrix) the solutions of (1.1)-(1.2) enjoy the (local in time) Strichartz estimates

$$
\begin{equation*}
\|u\|_{L^{p}\left([0,1] ; L^{q}\left(\mathbb{R}^{d}\right)\right)} \leq C\left\|u_{0}\right\|_{L^{2}\left(\mathbb{R}^{d}\right)}, \tag{1.3}
\end{equation*}
$$

where

$$
\begin{equation*}
\frac{2}{p}+\frac{d}{q}=\frac{d}{2}, \quad p \geq 2, \quad(p, q) \neq(2, \infty) \tag{1.4}
\end{equation*}
$$

Moreover (1.3) is global in time which means that one can replace [0, 1 ] in the left hand-side of (1.3) by $\mathbb{R}$. In [4, 8] one studies the possible extensions of (1.3) to the situation where $M$ is compact. An important new phenomenon that one has to take into account, when $M$ is compact, is the unavoidable derivative loss in (1.3) for some values of $(p, q)$. By "loss" we mean that $\left\|u_{0}\right\|_{L^{2}}$ in the right hand-side of (1.3) should be replaced by $\left\|u_{0}\right\|_{H^{s}}$ for some positive $s$. Here are two significant examples. If $M$ is the standard sphere $S^{d}, d \geq 3$, then it is proved in [8] that the solutions of (1.1)-(1.2) satisfy

$$
\begin{equation*}
\|u\|_{L^{2}\left([0,1] ; L^{\frac{2 d}{d-2}}\left(S^{d}\right)\right)} \leq C\left\|u_{0}\right\|_{H^{\frac{1}{2}}\left(S^{d}\right)} \tag{1.5}
\end{equation*}
$$

(notice that the couple $\left(2, \frac{2 d}{d-2}\right)$ satisfies (1.4)). Moreover, the $H^{1 / 2}\left(S^{d}\right)$ norm in the right hand-side of (1.5) is sharp in the sense that for every $\varepsilon>0$, the estimate

$$
\|u\|_{L^{2}\left([0,1] ; L^{\frac{2 d}{d-2}}\left(S^{d}\right)\right)} \leq C\left\|u_{0}\right\|_{H^{\frac{1}{2}-\varepsilon}\left(S^{d}\right)}
$$

is false.
A second example where one should encounter losses in (1.3) is the flat torus $\mathbb{T}^{d}=\mathbb{R}^{d} \mid(2 \pi \mathbb{Z})^{d}$. More precisely the estimate

$$
\begin{equation*}
\|u\|_{L^{\frac{2(d+2)}{d}}\left([0,1] ; L^{\frac{2(d+2)}{d}\left(\mathbb{T}^{d}\right)}\right)} \leq C\left\|u_{0}\right\|_{L^{2}\left(\mathbb{T}^{d}\right)} \tag{1.6}
\end{equation*}
$$

is false (notice that again the couple $\left(\frac{2(d+2)}{d}, \frac{2(d+2)}{d}\right)$ satisfies (1.4)). We refer to [4] for a counterexample disproving (1.6) in the case $d=1$. The extension to higher dimensions is straightforward. One may however expect (1.6) to be replaced by

$$
\begin{equation*}
\|u\|_{L^{\frac{2(d+2)}{d}}\left([0,1] ; L^{\frac{2(d+2)}{d}}\left(\mathbb{T}^{d}\right)\right)} \leq C_{s}\left\|u_{0}\right\|_{H^{s}\left(\mathbb{T}^{d}\right)}, \quad s>0 . \tag{1.7}
\end{equation*}
$$

Estimate (1.7) is known for $d=1,2$ (see [4]) (in this case $\frac{2(d+2)}{d}$ is an even integer). For $d \geq 3$, the study of (1.7) leads to an interesting open problem.

When $M$ is not compact, extensions of (1.3) were recently studied by several authors (see [33, 8, 25, 19]). In the works [33, 25, 19 the authors consider non
compact manifolds with metrics which are a "small" perturbation at infinity of a fixed "nice" metric, satisfying a non trapping assumption on the geodesic flow. It turns out that in such a geometric situation, one can prove exactly the same estimates as for the Euclidean space.

In [8], one considers $(M, g)$ to be $\mathbb{R}^{d}$ with a perturbation of the flat metric without the non trapping assumption. In this context one can get the Strichartz estimates with losses, just as in the case of a compact manifold. It is however a priori not clear whether losses of derivatives in the Strichartz estimates may come from the geometry at infinity. The main goal of this paper is to show that one can not have losses in the Strichartz inequalities coming from the geometry at infinity in the case of long range perturbations of the euclidean metric on $\mathbb{R}^{d}$.
Theorem 1. Consider $\mathbb{R}^{d}$ equipped with a smooth Riemannian metric

$$
g(x)=\left(g_{i j}(x)\right)_{i, j=1}^{d}, \quad x \in \mathbb{R}^{d}
$$

satisfying for some $\rho>0$,

$$
\begin{equation*}
\left|\partial^{\alpha}\left(g^{i j}(x)-\delta_{i j}\right)\right| \leq C_{\alpha}\langle x\rangle^{-\rho-|\alpha|}, \quad i, j \in\{1, \ldots d\} \tag{1.8}
\end{equation*}
$$

( $\delta_{i j}$ being the Kronecker symbol) and

$$
\begin{equation*}
c \mathrm{Id} \leq g(x) \leq C \mathrm{Id} \tag{1.9}
\end{equation*}
$$

Then there exists $R>0$ such that for every $T>0$, every $(p, q)$ satisfying (1.4) there exists $C>0$ such that for every $f \in L^{2}\left(\mathbb{R}^{d}\right)$,

$$
\begin{equation*}
\left\|e^{i t \Delta_{g}} f\right\|_{L^{p}\left([-T, T] ; L^{q}(|x| \geq R)\right)} \leq C\|f\|_{L^{2}\left(\mathbb{R}^{d}\right)} \tag{1.10}
\end{equation*}
$$

Moreover, for every $f \in H^{\frac{1}{p}}\left(\mathbb{R}^{d}\right)$,

$$
\begin{equation*}
\left\|e^{i t \Delta_{g}} f\right\|_{L^{p}\left([-T, T] ; L^{q}(|x| \leq R)\right)} \leq C\|f\|_{H^{\frac{1}{p}}\left(\mathbb{R}^{d}\right)} \tag{1.11}
\end{equation*}
$$

Remark 1.1. The result of Theorem $\mathbb{1}$ is stated only for metric perturbations of the flat Laplacian. However, an examination of our proof shows that the statement still holds if we add long range lower order terms. The same remark is valid for Theorem 圆 below.

Remark 1.2. Let us emphasize that estimates in the spirit of (1.10) are known to hold in the context of resolvent estimates for long range perturbations of the Laplacian (cf. [6, 10]). By this we mean the following: the fact that we have no derivative loss in (1.10) as in the free case is somehow similar to the fact that the high energy resolvent estimates with weights supported near infinity are the same as for the free resolvent.

If we suppose that the metric $g$ is non trapping then one can improve (1.11) and get the full family of Strichartz estimates. Recall that $g$ is non trapping if every geodesic (globally defined thanks to (1.9)) leaves every compact set in finite time. Let us now state our second result.
Theorem 2. Under the assumptions of Theorem 1, if in addition we suppose that $g$ is non trapping, then

$$
\begin{equation*}
\left\|e^{i t \Delta_{g}} f\right\|_{L^{p}\left([-T, T] ; L^{q}\left(\mathbb{R}^{d}\right)\right)} \leq C\|f\|_{L^{2}\left(\mathbb{R}^{d}\right)} \tag{1.12}
\end{equation*}
$$

Note that under the short range condition $\rho>1$, estimate (1.12) is proved in [25] by using FBI transform techniques.

Let us now explain the main points in the proof of the above results. The proof of (1.10) is based on the Isozaki-Kitada [21] parametrix construction. Recall that this construction was introduced to build modified scattering operators for long range perturbations of the free Schrödinger group. Let us point out that, since here we are only dealing with finite time estimates, we are not using the Isozaki-Kitada method in its full strength. In particular, we do not need a non trapping assumption on the metric to get (1.10). If we were interested in proving (1.10) with a constant $C$ uniform with respect to $T$, then a non trapping assumption and the full force of the Isozaki-Kitada method would be needed. We will not address this interesting issue here. See [7, 29] for the proof of the global in time estimates in the case of compactly supported perturbations. The proof of (1.11) is essentially contained in [8]. The proof of Theorem 2 is based on ideas introduced in [33, 8]. In fact, it is fair to say that, as far as the spatial regularity is concerned, the estimates established in [8] are gaining $1 / 2$ derivative with respect to the Sobolev embedding. We prove Theorem 2 by showing that the missing $1 / 2$ derivative can be recovered thanks to the local smoothing effect (when it is available). Let us notice that this effect is a consequence of standard resolvent estimates for non trapping perturbations of the Laplacian. It would be interesting to know whether intermediate situations may exist and if so to quantify them in terms of the metric. It is worth mentioning the work (7), where (1.12) with $C_{\varepsilon}\|f\|_{H^{\varepsilon}}, \varepsilon>0$ instead of $C\|f\|_{L^{2}}$ is studied, i.e. an unnecessary $\varepsilon$ derivative loss is accepted. In this context, let us recall that the analysis in [8, 9] has shown that, if one is interested in non linear applications, the losses in term of Sobolev regularity in the right hand-side of (1.12) are more dramatic then the losses in terms of the range of possible values of $(p, q)$ in the left hand-side of (1.12).

The rest of this paper is organized as follows. In the next section, we introduce the functional calculus for $\Delta_{g}$ in terms of pseudo differential calculus. In section 3, we recall the main points of the Isozaki-Kitada parametrix. The analysis of section 3 is then used in section 4 for the proof of (1.10). Section 5 deals with estimates
on time intervals depending on the frequency localization. In section 6, we use the non trapping condition to get the full family of Strichartz estimates in a fixed compact set. Section 7 is devoted to the rather standard non linear applications of the analysis of the previous sections.

Notation. In this paper several numerical constants will be denoted by the same $C$. For $T>0, p \in[1,+\infty]$, and $B$ a Banach space, we denote by $L_{T}^{p} B$ the Banach space of $L^{p}$ functions on $[-T, T]$ with values in $B$ equipped with the natural norm. We denote by $P \geq 0$ the self adjoint realization of $-\Delta_{g}$ on $L^{2}\left(\mathbb{R}^{d}\right)$.

## 2. Functional calculus

In this section $g$ is a metric on $\mathbb{R}^{d}$ such that there exist $C \geq 1 \geq c>0$ such that

$$
\begin{equation*}
c|\xi|^{2} \leq g(x)(\xi, \xi) \leq C|\xi|^{2}, \quad \forall x \in \mathbb{R}^{d}, \forall \xi \in \mathbb{R}^{d}, \tag{2.1}
\end{equation*}
$$

and for every $\alpha \in \mathbb{N}^{d}$,

$$
\begin{equation*}
\left|\partial_{x}^{\alpha} g(x)\right| \leq C_{\alpha} \tag{2.2}
\end{equation*}
$$

Notice that we do not assume the long range condition (1.8).
Further, we denote by

$$
p_{0}(x, \xi):=g(x)(\xi, \xi) \equiv \sum_{i, j=1}^{d} g^{i j}(x) \xi_{i} \xi_{j}
$$

the principal symbol of $-\Delta_{g}$. Here we adopt the standard notation for $\left(g^{i j}(x)\right) \equiv$ $\left(g_{i j}(x)\right)^{-1}$.

The goal of this section is to approximate $\left.\left.\varphi\left(-h^{2} \Delta_{g}\right), h \in\right] 0,1\right]$, for a given bump function $\varphi: \mathbb{R} \rightarrow \mathbb{R}$, by a suitable semi-classical pseudo differential operator. Similar considerations are performed in [8], where $\Delta_{g}$ is the Laplace-Beltrami operator on a compact Riemannian manifold. Here we follow a similar scheme. The new feature in our analysis is the $L^{p}$ bound for the remainder of the pseudo differential expansion of $\varphi\left(-h^{2} \Delta_{g}\right)$. In [8] this is done by only invoking $L^{2}$ considerations and the fact that, on a compact manifold $M, L^{\infty}(M)$ is continuously embedded in $L^{2}(M)$. Here, we can not use this fact. We overcome the difficulty by using $L^{p}$ bounds for powers of the resolvent of $P$.

Recall that $P \geq 0$ is the self adjoint realization of $-\Delta_{g}$ on $L^{2}\left(\mathbb{R}^{d}\right)$. We first collect several classical properties of $P$. For every $s \in \mathbb{R}$, there exists a constant $C_{s}$ such that for every $u \in \mathcal{S}\left(\mathbb{R}^{d}\right)$,

$$
\begin{equation*}
C_{s}^{-1}\left\|(P+1)^{s / 2} u\right\|_{L^{2}\left(\mathbb{R}^{d}\right)} \leq\|u\|_{H^{s}\left(\mathbb{R}^{d}\right)} \leq C_{s}\left\|(P+1)^{s / 2} u\right\|_{L^{2}\left(\mathbb{R}^{d}\right)} . \tag{2.3}
\end{equation*}
$$

Next, we recall that the Schwartz class is stable under the action of the resolvent of $P$. More precisely, for every $z \in \mathbb{C} \backslash\left[0, \infty\left[\right.\right.$ the map $(P-z)^{-1}$ is acting continuously
on $\mathcal{S}\left(\mathbb{R}^{d}\right)$. As a consequence, by the standard duality argument it acts continuously on $\mathcal{S}^{\prime}\left(\mathbb{R}^{d}\right)$ too. In particular, for every $w \in L^{p}\left(\mathbb{R}^{d}\right), 1 \leq p \leq+\infty,(P-z)^{-1} w$ is well defined.

The elliptic nature of $P$ also implies that, for every $s \in \mathbb{R}$, there exists $C_{s}$ such that for every $u \in \mathcal{S}\left(\mathbb{R}^{d}\right)$,

$$
\left\|(P+1)^{-1} u\right\|_{H^{s+2}\left(\mathbb{R}^{d}\right)} \leq C_{s}\|u\|_{H^{s}\left(\mathbb{R}^{d}\right)} .
$$

As a consequence for every $z \in \mathbb{C} \backslash[0, \infty[$,

$$
\begin{equation*}
\left\|(P-z)^{-1} u\right\|_{H^{s+2}\left(\mathbb{R}^{d}\right)} \leq \frac{\widetilde{C}_{s}\langle z\rangle}{|\operatorname{Im} z|}\|u\|_{H^{s}\left(\mathbb{R}^{d}\right)} . \tag{2.4}
\end{equation*}
$$

Indeed, we can write

$$
\begin{aligned}
\left\|(P-z)^{-1} u\right\|_{H^{s+2}} & \leq C\left\|(P+1)^{\frac{s+2}{2}}(P-z)^{-1} u\right\|_{L^{2}} \\
& \leq C\left\|(P-z)^{-1}(P+1)(P+1)^{\frac{s+2}{2}}(P+1)^{-1} u\right\|_{L^{2}} \\
& \leq C\left\|(P-z)^{-1}(P+1)\right\|_{L^{2} \rightarrow L^{2}}\left\|(P+1)^{\frac{s}{2}} u\right\|_{L^{2}} \\
& \leq C \sup _{\lambda \in \mathbb{R}^{+}} \frac{|\lambda+1|}{|\lambda-z|}\|u\|_{H^{s}} \\
& \leq \frac{C\langle z\rangle}{|\operatorname{Im} z|}\|u\|_{H^{s}}
\end{aligned}
$$

which proves (2.4).
We next state a bound for sufficiently large powers of the resolvent of $P$.
Proposition 2.1. Let $n_{0}>d / 2$ be an integer. For $z \in \mathbb{C} \backslash[0, \infty[$, we denote by $K_{z}(x, y)$ the Schwartz kernel of the operator $(P-z)^{-n_{0}}$. Then for every $\alpha \in \mathbb{N}^{d}$, there exist $C_{\alpha}>0$ and $n(\alpha) \in \mathbb{N}$ such that for every $(x, y) \in \mathbb{R}^{2 d}$, and every $z \in \mathbb{C} \backslash[0, \infty[$,

$$
\left|(x-y)^{\alpha} K_{z}(x, y)\right| \leq C_{\alpha}\left(\frac{\langle z\rangle}{|\operatorname{Im} z|}\right)^{n(\alpha)}
$$

In particular there exist $N$ and $C>0$ such that

$$
\left|K_{z}(x, y)\right| \leq C\left(\frac{\langle z\rangle}{|\operatorname{Im} z|}\right)^{N} \frac{1}{\left(1+|x-y|^{2}\right)^{d}} .
$$

Proof. Let first $|\alpha|=0$. Observe that for $s>d / 2$ any bounded linear map

$$
A: H^{-s}\left(\mathbb{R}^{d}\right) \longrightarrow H^{s}\left(\mathbb{R}^{d}\right)
$$

has a Schwartz kernel $K_{A}(x, y)$ which is a bounded continuous function on $\mathbb{R}^{2 d}$. Moreover, there exists $C>0$ such that for every bounded map from $H^{-s}$ to $H^{s}$,
every $(x, y) \in \mathbb{R}^{2 d}$,

$$
\begin{equation*}
\left|K_{A}(x, y)\right| \leq C\||A \||, \tag{2.5}
\end{equation*}
$$

where $\||A|\|$ denotes the norm of $A$,

$$
\left\|\left|A\left\|\mid:=\sup _{\|u\|_{H^{-s}=1}}\right\| A u \|_{H^{s}} .\right.\right.
$$

Using (2.4), we get the estimate

$$
\left\|(P-z)^{-n_{0}} u\right\|_{H^{n_{0}}} \leq C\left(\frac{\langle z\rangle}{|\operatorname{Im} z|}\right)^{n_{0}}\|u\|_{H^{-n_{0}}}
$$

Therefore, in view of (2.5), the assertion of Proposition 2.1 holds for $|\alpha|=0$.
Let next $|\alpha|=1$. Clearly, for $k \in\{1, \ldots, d\},\left(x_{k}-y_{k}\right) K_{z}(x, y)$ is the kernel of $\left[x_{k},(P-z)^{-n_{0}}\right]$. Using the identity

$$
\left[x_{k},(P-z)^{-1}\right]=(P-z)^{-1}\left[P, x_{k}\right](P-z)^{-1}
$$

we arrive at the formula

$$
\begin{aligned}
{\left[x_{k},(P-z)^{-n_{0}}\right] } & =\sum_{j=0}^{n_{0}-1}(P-z)^{-(j+1)}\left[P, x_{k}\right](P-z)^{-\left(n_{0}-j\right)} \\
& =\sum_{j_{1}+j_{2}=n_{0}-1}(P-z)^{-\left(j_{1}+1\right)}\left[P, x_{k}\right](P-z)^{-\left(j_{2}+1\right)}
\end{aligned}
$$

On the other hand, by invoking (2.4), we obtain that if $j_{1}+j_{2}=n_{0}-1$, and if $s \in \mathbb{R}$ is such that $n_{0}>s>d / 2$, then the linear map

$$
(P-z)^{-\left(j_{1}+1\right)}\left[P, x_{k}\right](P-z)^{-\left(j_{2}+1\right)}
$$

is bounded from $H^{-s}\left(\mathbb{R}^{d}\right)$ to

$$
H^{-s+2\left(j_{2}+1\right)-1+2\left(j_{1}+1\right)}\left(\mathbb{R}^{d}\right)=H^{-s+2 n_{0}+1}\left(\mathbb{R}^{d}\right) \subset H^{s}\left(\mathbb{R}^{d}\right)
$$

with operator norm bounded by a polynomial of

$$
\begin{equation*}
\frac{\langle z\rangle}{|\operatorname{Im} z|} \tag{2.6}
\end{equation*}
$$

which, using (2.5), yields the assertion of Proposition 2.1 for $|\alpha|=1$.
Let finally $|\alpha|$ be arbitrary. For $k \in\{1, \ldots, d\}, \alpha \in \mathbb{N}^{d}$ and $A$ a function of $P$, we introduce the notations

$$
\zeta_{k} A:=\left[x_{k}, A\right], \quad \zeta^{\alpha} A=\zeta_{1}^{\alpha_{1}} \ldots \zeta_{d}^{\alpha_{d}} A
$$

Notice that $\zeta_{k} \zeta_{j} A=\zeta_{j} \zeta_{k} A$. Using an induction argument, one can check that

$$
\zeta^{\alpha}(P-z)^{-n_{0}}
$$

is a linear combination of terms of type
$(P-z)^{-\left(1+j_{1}\right)}\left(\zeta^{\alpha^{1}} P\right)(P-z)^{-\left(1+j_{2}\right)}\left(\zeta^{\alpha^{2}} P\right) \ldots(P-z)^{-\left(1+j_{n}\right)}\left(\zeta^{\alpha^{n}} P\right)(P-z)^{-\left(1+j_{n+1}\right)}$
with

$$
\alpha^{i} \neq 0, \quad 1 \leq n \leq|\alpha|, \quad \alpha^{1}+\cdots+\alpha^{n}=\alpha, \quad j_{1}+j_{2}+\cdots+j_{n+1}=n_{0}-1 .
$$

Therefore for $n_{0}>s>d / 2$, the map $\zeta^{\alpha}(P-z)^{-n_{0}}$ is bounded map from $H^{-s}\left(\mathbb{R}^{d}\right)$ to $H^{2 n_{0}+|\alpha|-s}\left(\mathbb{R}^{d}\right)$ with operator norm bounded by a polynomial of (2.6). By invoking (2.5), we complete the proof of Proposition 2.1.

We will use the result of Proposition 2.1 to get $L^{p}$ bounds for sufficiently large powers of the resolvent of $P$. For that purpose, we recall the well known Schur lemma.

Proposition 2.2. Let $K(x, y)$ be a continuous function on $\mathbb{R}^{2 d}$ satisfying

$$
\int_{\mathbb{R}^{d}}|K(x, y)| d x \leq C, \quad \int_{\mathbb{R}^{d}}|K(x, y)| d y \leq C
$$

for some positive constant $C$. Then for every $p \in[1,+\infty]$, the linear map with kernel $K$ is bounded on $L^{p}\left(\mathbb{R}^{d}\right)$ with norm $\leq C$.

Proof. The statement is straightforward for $p=\infty$ and $p=1$. The case of an arbitrary $p$ follows then by interpolation (one can also easily give a direct proof avoiding the interpolation).

A direct combination of Proposition 2.1 and Proposition 2.2 gives the following statement.

Proposition 2.3. Let us fix an integer $n_{0}>d / 2$. Then there exist $C>0$ and $n \in \mathbb{N}$ such that for every $p \in[1, \infty]$, $u \in L^{p}\left(\mathbb{R}^{d}\right)$, every $z \in \mathbb{C} \backslash[0, \infty[$,

$$
\left\|(P-z)^{-n_{0}} u\right\|_{L^{p}\left(\mathbb{R}^{d}\right)} \leq C\left(\frac{\langle z\rangle}{|\operatorname{Im} z|}\right)^{n}\|u\|_{L^{p}\left(\mathbb{R}^{d}\right)} .
$$

We next give another consequence of the Schur lemma.
Proposition 2.4. Let $m>d$ and let $a(x, \xi, h)$ be a continuous function on $\mathbb{R}^{d} \times$ $\left.\left.\mathbb{R}^{d} \times\right] 0,1\right]$ smooth with respect to the second variable satisfying

$$
\begin{equation*}
\left|\partial_{\xi}^{\beta} a(x, \xi, h)\right| \leq C_{\beta}\langle\xi\rangle^{-m}, \quad \beta \in \mathbb{N}^{d} . \tag{2.7}
\end{equation*}
$$

Then for $\infty \geq r \geq q \geq 1$,

$$
\begin{equation*}
\left.\left.\|a(x, h D, h)\|_{L^{q}\left(\mathbb{R}^{d}\right) \rightarrow L^{r}\left(\mathbb{R}^{d}\right)} \leq C_{q r} h^{d\left(\frac{1}{r}-\frac{1}{q}\right)}, \quad \forall h \in\right] 0,1\right] \tag{2.8}
\end{equation*}
$$

with

$$
C_{q r} \lesssim \sup _{|\beta| \leq d+1} C_{\beta}
$$

Proof. Write

$$
\begin{equation*}
a(x, h D, h) u(x)=\int_{\mathbb{R}^{d}} K_{h}(x, y) u(y) d y \tag{2.9}
\end{equation*}
$$

with

$$
K_{h}(x, y)=(2 \pi h)^{-d} \int_{\mathbb{R}^{d}} e^{i \frac{(x-y) \cdot \xi}{h}} a(x, \xi, h) d \xi
$$

Notice that, thanks to (2.7) and the assumption $m>d$, the last integral is absolutely convergent. Let us denote by $\hat{a}$ the Fourier transform of $a$ with respect to the second variable. Then we can write

$$
K_{h}(x, y)=(2 \pi h)^{-d} \hat{a}\left(x, \frac{x-y}{h}, h\right)
$$

Thanks to (2.7), for every $\beta \in \mathbb{N}^{d}$,

$$
\left|z^{\beta} \hat{a}(x, z, h)\right| \lesssim C_{\beta}
$$

Hence

$$
\sup _{x \in \mathbb{R}^{d}} \sup _{h \in] 0,1]}|\hat{a}(x, z, h)| \lesssim\langle z\rangle^{-d-1} \sup _{|\beta| \leq d+1} C_{\beta}
$$

Therefore

$$
\int_{\mathbb{R}^{d}}\left|K_{h}(x, y)\right| d x=(2 \pi)^{-d} \int_{\mathbb{R}^{d}}|\hat{a}(y+h z, z, h)| d z \lesssim \sup _{|\beta| \leq d+1} C_{\beta}
$$

and

$$
\int_{\mathbb{R}^{d}}\left|K_{h}(x, y)\right| d y=(2 \pi)^{-d} \int_{\mathbb{R}^{d}}|\hat{a}(x, z, h)| d z \lesssim \sup _{|\beta| \leq d+1} C_{\beta}
$$

Applying Proposition 2.2 completes the proof of (2.8) for $q=r$. Thanks to (2.7),

$$
\left|K_{h}(x, y)\right| \lesssim C_{0} h^{-d}
$$

and coming back to (2.9), this completes the proof of (2.8) for $r=\infty$ and $q=1$. Let us finally fix arbitrary $r, q$ satisfying $\infty \geq r \geq q \geq 1$. Interpolating between the $L^{1} \rightarrow L^{1}$ and $L^{1} \rightarrow L^{\infty}$ bounds, we get the $L^{1} \rightarrow L^{r / q}$ bound. Then, interpolation between the $L^{\infty} \rightarrow L^{\infty}$ and the $L^{1} \rightarrow L^{r / q}$ bounds, we get the $L^{q} \rightarrow L^{r}$ bound. This completes the proof of Proposition 2.4.

The next statement describes functions of $P$ in terms of semi-classical pseudo differential operators. Recall that $p_{0}$ denotes the principal symbol of $P$.

Proposition 2.5. Let $g$ be a metric on $\mathbb{R}^{d}$ satisfying (2.1) and (2.2). Then for every $\varphi \in C_{0}^{\infty}(\mathbb{R})$, there exist symbols $\left(a_{k}\right)_{k \geq 0}$ satisfying

$$
\begin{gather*}
\left|\partial_{x}^{\alpha} \partial_{\xi}^{\beta} a_{k}(x, \xi)\right| \leq C_{k \alpha \beta}, \quad \forall(x, \xi) \in \operatorname{supp} \varphi \circ p_{0}  \tag{2.10}\\
a_{k}(x, \xi)=0, \quad \forall(x, \xi) \notin \operatorname{supp} \varphi \circ p_{0} \tag{2.11}
\end{gather*}
$$

and there exists $n_{1} \in \mathbb{N}$ such that for every $N \geq 1$, every $1 \leq r \leq q \leq \infty$, there exists $C_{N q r}$ such that for every $\left.\left.h \in\right] 0,1\right]$,

$$
\left\|\varphi\left(h^{2} P\right)-\sum_{k=0}^{N} h^{k} a_{k}(x, h D)\right\|_{L^{q}\left(\mathbb{R}^{d}\right) \rightarrow L^{r}\left(\mathbb{R}^{d}\right)} \leq C_{N q r} h^{N-n_{1}+d\left(\frac{1}{r}-\frac{1}{q}\right)}
$$

Moreover for every $s \geq 0$ there exists $N_{s}$ such that for $N \geq N_{s}$,

$$
\left\|\varphi\left(h^{2} P\right)-\sum_{k=0}^{N} h^{k} a_{k}(x, h D)\right\|_{H^{-s}\left(\mathbb{R}^{d}\right) \rightarrow H^{s}\left(\mathbb{R}^{d}\right)} \leq C h^{N-N_{s}}
$$

Remark 2.6. If we suppose that the metric $g$ satisfies (1.8) then we can replace the bound (2.10) in Proposition 2.5 by

$$
\begin{equation*}
\left|\partial_{x}^{\alpha} \partial_{\xi}^{\beta} a_{k}(x, \xi)\right| \leq C_{k \alpha \beta}\langle x\rangle^{-k-|\alpha|} \tag{2.12}
\end{equation*}
$$

for every $(x, \xi) \in \mathbb{R}^{2 d}$ such that $p_{0}(x, \xi) \in \operatorname{supp} \varphi$.
Proof of Proposition 2.5. We first describe the classical construction (see [30, 26, 8]) of a parametrix for $\left(h^{2} P-z\right)^{-1}, z \in \mathbb{C} \backslash[0, \infty[, h \in] 0,1]$. There exist symbols

$$
q_{0}(x, \xi, z), q_{1}(x, \xi, z), \ldots, q_{N}(x, \xi, z)
$$

satisfying

$$
\left|\partial_{x}^{\alpha} \partial_{\xi}^{\beta} q_{k}(x, \xi, z)\right| \leq C_{k \alpha \beta}\langle\xi\rangle^{-2-k-|\beta|}\left(\frac{\langle z\rangle}{|\operatorname{Im} z|}\right)^{n(k, \alpha, \beta)}
$$

such that for every $h \in] 0,1]$, and for every $N \geq 1$,

$$
\left(h^{2} P-z\right) \sum_{k=0}^{N} q_{k}(x, h D, z)=\operatorname{Id}+h^{N+1} r_{N+1}(x, h D, z, h)
$$

with

$$
\left|\partial_{x}^{\alpha} \partial_{\xi}^{\beta} r_{N+1}(x, \xi, z, h)\right| \leq C_{N \alpha \beta}\langle\xi\rangle^{-N-1-|\beta|}\left(\frac{\langle z\rangle}{|\operatorname{Im} z|}\right)^{n(N, \alpha, \beta)}
$$

uniformly in $h \in] 0,1]$. Moreover the symbols are analytic with respect to $z \in$ $\mathbb{C} \backslash[0, \infty[$ and we can write

$$
q_{0}(x, \xi, z)=\frac{1}{p_{0}(x, \xi)-z}
$$

where $p_{0}(x, \xi)=g(x)(\xi, \xi)$ is the principal symbol of $-\Delta_{g}$. In addition, for $k \geq 1$, $q_{k}$ takes the form

$$
q_{k}(x, \xi, z)=\sum_{j=1}^{2 k-1} \frac{d_{j, k}(x, \xi)}{\left(p_{0}(x, \xi)-z\right)^{1+j}}
$$

with $d_{j, k} \in S^{2 j-k}\left(\mathbb{R}^{2 d}\right)$ (they are polynomials in $\xi$ with coefficients which are linear combinations of products of derivatives of the coefficients of the inverse of the metric).

Therefore for every $h \in] 0,1]$ every $z \in \mathbb{C} \backslash[0, \infty[$, every $N \geq 1$,

$$
\left(h^{2} P-z\right)^{-1}=\sum_{k=0}^{N} h^{k} q_{k}(x, h D, z)-\left(h^{2} P-z\right)^{-1} h^{N+1} r_{N+1}(x, h D, z, h) .
$$

Then, for every $\varphi \in C_{0}^{\infty}(\mathbb{R})$, we can use the Helffer-Sjöstrand formula (see [20, 14]),

$$
\varphi\left(h^{2} P\right)=-\frac{1}{\pi} \int_{\mathbb{C}} \bar{\partial} \widetilde{\varphi}(z)\left(h^{2} P-z\right)^{-1} d L(z),
$$

where $d L(z)$ denotes the Lebesgue measure on $\mathbb{C}$ and $\widetilde{\varphi}(z) \in C_{0}^{\infty}(\mathbb{C})$ is an almost analytic extension of $\varphi$ which satisfies

$$
\begin{equation*}
\forall \Lambda>0, \quad|\bar{\partial} \widetilde{\varphi}(z)| \leq C_{\Lambda}|\operatorname{Im} z|^{\Lambda} \tag{2.13}
\end{equation*}
$$

This implies that $\varphi\left(h^{2} P\right)$ can be written as

$$
\begin{align*}
& \varphi\left(h^{2} P\right)=\sum_{k=0}^{N} h^{k} a_{k}(x, h D)+  \tag{2.14}\\
&+\frac{h^{N+1}}{\pi} \int_{\mathbb{C}} \bar{\partial} \widetilde{\varphi}(z)\left(h^{2} P-z\right)^{-1} r_{N+1}(x, h D, z, h) d L(z)
\end{align*}
$$

with

$$
a_{0}(x, \xi)=\varphi\left(p_{0}(x, \xi)\right)
$$

and, for $k \geq 1$,

$$
a_{k}(x, \xi)=\sum_{j=1}^{2 k-1} \frac{(-1)^{j}}{j!} d_{j, k}(x, \xi) \varphi^{(j)}\left(p_{0}(x, \xi)\right)
$$

We now state a bound for the action of the map $a_{k}(x, h D)$ on Lebesgue spaces.
Lemma 2.7. For $1 \leq q \leq r \leq \infty$,

$$
\left\|a_{k}(x, h D)\right\|_{L^{q}\left(\mathbb{R}^{d}\right) \rightarrow L^{r}\left(\mathbb{R}^{d}\right)} \leq C_{q r k} h^{d\left(\frac{1}{r}-\frac{1}{q}\right)} .
$$

Proof. Since $a_{k}(x, \xi)$ is smooth and compactly supported with respect to $\xi$, Lemma 2.7 is a direct consequence of Proposition 2.4.

We next state a bound for the second term in the right hand-side of (2.14).
Lemma 2.8. Let $n_{0}>d / 2$ be an integer. Set

$$
R_{N}:=\frac{h^{N+1}}{\pi} \int_{\mathbb{C}} \bar{\partial} \widetilde{\varphi}(z)\left(h^{2} P-z\right)^{-1} r_{N+1}(x, h D, z, h) d L(z) .
$$

Then for every $N>d+2 n_{0}-3$, every $1 \leq q \leq r \leq \infty$, there exists $C>0$ such that for every $u \in L^{r}\left(\mathbb{R}^{d}\right)$, every $\left.\left.h \in\right] 0,1\right]$,

$$
\begin{equation*}
\left\|R_{N} u\right\|_{L^{q}\left(\mathbb{R}^{d}\right)} \leq C h^{N+1-2 n_{0}+d\left(\frac{1}{r}-\frac{1}{q}\right)}\|u\|_{L^{r}\left(\mathbb{R}^{d}\right)} . \tag{2.15}
\end{equation*}
$$

Moreover for every $s \geq 0$ there exists $N_{s}$ such that for every $N \geq N_{s}$, there exists $C>0$ such that for every $h \in] 0,1]$, every $u \in H^{-s}\left(\mathbb{R}^{d}\right)$,

$$
\begin{equation*}
\left\|h^{N+1} R_{N} u\right\|_{H^{s}\left(\mathbb{R}^{d}\right)} \leq C h^{N-N_{s}}\|u\|_{H^{-s}\left(\mathbb{R}^{d}\right)} . \tag{2.16}
\end{equation*}
$$

Proof. Define $\widetilde{r}_{N+1}(x, h D, z, h)$ by setting

$$
\left(h^{2} P-z\right)^{-1} r_{N+1}(x, h D, z, h)=\left(h^{2} P-z\right)^{-n_{0}} \widetilde{r}_{N+1}(x, h D, z, h),
$$

with a fixed $n_{0}>d / 2$. Then $\widetilde{r}_{N+1}$ satisfies

$$
\left|\partial_{x}^{\alpha} \partial_{\xi}^{\beta} \widetilde{r}_{N+1}(x, \xi, z, h)\right| \leq C_{\alpha \beta}\langle\xi\rangle^{-N-1+2\left(n_{0}-1\right)-|\beta|}\left(\frac{\langle z\rangle}{|\operatorname{Im} z|}\right)^{n(\alpha, \beta)}
$$

If $N+1-2\left(n_{0}-1\right)>d$, i.e. $N>d+2 n_{0}-3$, then Proposition 2.4 implies that there exists $n_{1} \in \mathbb{N}$ such that

$$
\left\|\widetilde{r}_{N+1}(x, h D, z, h)\right\|_{L^{q}\left(\mathbb{R}^{d}\right) \rightarrow L^{r}\left(\mathbb{R}^{d}\right)} \leq C h^{d\left(\frac{1}{r}-\frac{1}{q}\right)}\left(\frac{\langle z\rangle}{|\operatorname{Im} z|}\right)^{n_{1}} .
$$

On the other hand, we can write

$$
\left(h^{2} P-z\right)^{-n_{0}}=h^{-2 n_{0}}\left(P-h^{-2} z\right)^{-n_{0}}
$$

thus Proposition 2.3 shows that there exists $n \in \mathbb{N}$ such that for every $r \in[1, \infty]$,

$$
\left\|\left(h^{2} P-z\right)^{-n_{0}}\right\|_{L^{r}\left(\mathbb{R}^{d}\right) \rightarrow L^{r}\left(\mathbb{R}^{d}\right)} \leq C_{r} h^{-2 n_{0}}\left(\frac{\left\langle h^{-2} z\right\rangle}{\left|\operatorname{Im} h^{-2} z\right|}\right)^{n} \leq \tilde{C}_{r} h^{-2 n_{0}}\left(\frac{\langle z\rangle}{|\operatorname{Im} z|}\right)^{n}
$$

where we used that for every $h \in] 0,1]$, every $z \in \mathbb{C} \backslash[0, \infty[$,

$$
\frac{\left\langle h^{-2} z\right\rangle}{\left|\operatorname{Im} h^{-2} z\right|} \leq C \frac{\langle z\rangle}{|\operatorname{Im} z|} .
$$

The proof of (2.15) is completed by taking $\Lambda>n_{1}+n$ in (2.13). Let us next prove (2.16). It suffices to prove the result for $s$ an even integer. It is sufficient to study the action on $L^{2}\left(\mathbb{R}^{d}\right)$ of the map

$$
\int_{\mathbb{C}} \bar{\partial} \widetilde{\varphi}(z)\left(h^{2} P-z\right)^{-1}\left(1+h^{-2} h^{2} P\right)^{s / 2} r_{N+1}(x, h D, z, h)\left(1-h^{-2} h^{2} \Delta\right)^{s / 2} d L(z),
$$

where $\Delta$ is the flat Laplacian on $\mathbb{R}^{d}$. We can then write

$$
\left(1+h^{-2} h^{2} P\right)^{s / 2} r_{N+1}(x, h D, z, h)\left(1-h^{-2} h^{2} \Delta\right)^{s / 2}=h^{-2 s} \widetilde{r}_{N, s}(x, h D, z, h),
$$

where $\widetilde{r}_{N, s}(\cdot, \cdot, z, h) \in S^{-N-1+2 s}$ with semi-norms uniformly bounded with respect to $h \in] 0,1]$ by a by a polynomial of (2.6). Therefore, for $N+1-2 s>d$, we can apply (2.4) and Proposition [2.4 to conclude the proof of (2.16). This completes the proof of Lemma 2.8.

Combining Lemma 2.7 and Lemma 2.8 completes the proof of Proposition 2.5,
We will now give several consequences of Proposition [2.5 that we will use in the sequel. We first quote the following proposition which is a direct consequence of Proposition 2.4 and Proposition 2.5.
Proposition 2.9. Let $\varphi \in C_{0}^{\infty}(\mathbb{R})$. Then for every $\left.\left.h \in\right] 0,1\right]$, every $1 \leq r \leq q \leq \infty$,

$$
\left\|\varphi\left(h^{2} P\right)\right\|_{L^{q}\left(\mathbb{R}^{d}\right) \rightarrow L^{r}\left(\mathbb{R}^{d}\right)} \leq C_{q r} h^{d\left(\frac{1}{r}-\frac{1}{q}\right)}
$$

Next, we state a consequence of the Littlewood-Paley theory in terms of $\varphi\left(h^{2} P\right)$. Consider a Littlewood-Paley partition of the identity

$$
\mathrm{Id}=\varphi_{1}(P)+\sum_{h^{-1}: \text { dyadic }} \varphi\left(h^{2} P\right)
$$

where $\varphi_{1} \in C_{0}^{\infty}(\mathbb{R}), \varphi \in C_{0}^{\infty}(\mathbb{R} \backslash\{0\})$ and " $h^{-1}$ : dyadic" means that in the sum $h^{-1}$ takes all positive powers of 2 as values. The existence of such a partition is standard (see e.g. [1]). We then have the following statement.
Proposition 2.10. Let $T>0$. Then for every $u \in C\left([0, T] ; \mathcal{S}\left(\mathbb{R}^{d}\right)\right)$, every $p \in$ $[2, \infty]$, every $q \in[2,+\infty[$,

$$
\|u\|_{L_{T}^{p} L^{q}} \leq C\|u\|_{L_{T}^{\infty} L^{2}}+C\left(\sum_{h^{-1}: \text { dyadic }}\left\|\varphi\left(h^{2} P\right) u\right\|_{L_{T}^{p} L^{q}}^{2}\right)^{\frac{1}{2}}
$$

Proof. Proceeding as in [8, Corollary 2.3], we obtain that

$$
\|u\|_{L_{T}^{p} L^{q}} \leq C\|u\|_{L_{T}^{\infty} L^{2}}+C\left\|\left(\sum_{h^{-1}: \text { dyadic }}\left\|\varphi\left(h^{2} P\right) u\right\|_{L^{q}}^{2}\right)^{\frac{1}{2}}\right\|_{L^{p}[0, T]} .
$$

Since $p \geq 2$, the Minkowski inequality completes the proof of Proposition 2.10.
In the proof of Theorem 2, we will make use of the following statement.
Proposition 2.11. Let $\varphi \in C_{0}^{\infty}(\mathbb{R} \backslash\{0\})$. Then for every $s \in \mathbb{R}$ there exists $C_{s}$ such that for every $w \in \mathcal{S}\left(\mathbb{R}^{d}\right)$, every $\left.\left.h \in\right] 0,1\right]$,

$$
\begin{equation*}
\left\|\varphi\left(h^{2} P\right) w\right\|_{L^{2}\left(\mathbb{R}^{d}\right)} \leq C_{s} h^{s}\|w\|_{H^{s}\left(\mathbb{R}^{d}\right)} . \tag{2.17}
\end{equation*}
$$

Proof. Using (2.3) and the spectral theorem for $P$, we can write

$$
\begin{aligned}
\left\|\varphi\left(h^{2} P\right) w\right\|_{L^{2}} & =\left\|\varphi\left(h^{2} P\right)(1+P)^{-s / 2}(1+P)^{s / 2} w\right\|_{L^{2}} \\
& \leq C\left\|\varphi\left(h^{2} P\right)(1+P)^{-s / 2}\right\|_{L^{2} \rightarrow L^{2}}\|w\|_{H^{s}} \\
& \leq C \sup _{\lambda \sim h^{-2}}(1+\lambda)^{-s / 2}\|w\|_{H^{s}} \leq C h^{s}\|w\|_{H^{s}}
\end{aligned}
$$

This completes the proof of Proposition 2.11.
In applications to nonlinear problems, one may also need to use $L^{p}$ versions of Proposition 2.11. Here is a precise statement.
Proposition 2.12. Let $\varphi \in C_{0}^{\infty}(\mathbb{R} \backslash\{0\})$ and $p \in[1,+\infty]$. Then there exists $C$ such that for every $w \in \mathcal{S}\left(\mathbb{R}^{d}\right)$, every $\left.\left.h \in\right] 0,1\right]$,

$$
\left\|\varphi\left(h^{2} P\right) w\right\|_{L^{p}\left(\mathbb{R}^{d}\right)} \leq C h\|w\|_{W^{1, p}\left(\mathbb{R}^{d}\right)}
$$

Proof. In view of Proposition 2.5, it suffices to establish the bound

$$
\|a(x, h D) w\|_{L^{p}\left(\mathbb{R}^{d}\right)} \leq C h\|w\|_{W^{1, p}\left(\mathbb{R}^{d}\right)},
$$

where $a(x, \xi)$ is satisfying (2.10) and (2.11). Let $\psi \in C_{0}^{\infty}\left(\mathbb{R}^{d}\right)$ be such that

$$
a(x, \xi) \psi(\xi)=a(x, \xi)
$$

Thanks to Proposition [2.4, the map $a(x, h D)$ is bounded on $L^{p}\left(\mathbb{R}^{d}\right)$ and thus it is sufficient to prove that

$$
\|\psi(h D) w\|_{L^{p}\left(\mathbb{R}^{d}\right)} \leq C h\|w\|_{W^{1, p}\left(\mathbb{R}^{d}\right)}
$$

which is a well-known fact (see e.g. [1, Chapitre 2], [12]). This completes the proof of Proposition 2.12.

Remark 2.13. Proposition 2.12 will be important in the proof of Theorem 5 below. A similar bound in the context of a compact manifold was used in [8]. The additional point here is again the $L^{p}$ boundedness of the remainder in the pseudo differential expansion of $\varphi\left(h^{2} P\right)$ established in Proposition 2.5.

## 3. The Isozaki-Kitada parametrix

This section is devoted to the construction of Isozaki-Kitada. We only give the details for those arguments which are not written explicitly in the papers of the reference list. The reader interested in having all the details for the proofs of the statements in this section can consult [28, Section 4], [2, Appendice], [3, Appendix]. The reader may of course wish to consult the original paper by Isozaki-Kitada [21] which is nevertheless only written for potential perturbations and not in the semiclassical regime.

In this section $g$ is a metric on $\mathbb{R}^{d}$ satisfying (1.8) and (1.9). For $\left.J \Subset\right] 0,+\infty[$ an open interval, $R>0$ and $\sigma \in]-1,1[$, we consider the outgoing and incoming zones $\Gamma^{+}(R, J, \sigma)$ and $\Gamma^{-}(R, J, \sigma)$, defined by

$$
\Gamma^{ \pm}(R, J, \sigma)=\left\{(x, \xi) \in \mathbb{R}^{2 d}:|x|>R,|\xi|^{2} \in J, \pm \frac{\langle x, \xi\rangle}{|x||\xi|}>-\sigma\right\}
$$

The next statement is proved in Robert [28, prop. 4.1] (see also [21]).
Proposition 3.1. For every interval $J \Subset] 0,+\infty[$, every $\sigma \in]-1,1[$ there exist a large number $R$ and phase functions $S_{ \pm} \in C^{\infty}\left(\mathbb{R}^{2 d} ; \mathbb{R}\right)$ satisfying

$$
g(x)\left(\nabla_{x} S_{ \pm}, \nabla_{x} S_{ \pm}\right)=|\xi|^{2}, \quad(x, \xi) \in \Gamma^{ \pm}(R, J, \sigma)
$$

and, for every $(\alpha, \beta) \in \mathbb{N}^{2 d}$, there exists $C_{\alpha \beta}$ such that for every $(x, \xi) \in \mathbb{R}^{2 d}$,

$$
\left|\partial_{x}^{\alpha} \partial_{\xi}^{\beta}\left(S_{ \pm}(x, \xi)-\langle x, \xi\rangle\right)\right| \leq C_{\alpha \beta}\langle x\rangle^{1-\rho-|\alpha|} .
$$

We next state an easy consequence of Proposition 3.1.
Proposition 3.2. For every interval $J \Subset] 0,+\infty[$, every $\sigma \in]-1,1[$ there exist $a$ large number $\tilde{R}$ and two families of phase functions

$$
\left(S_{ \pm, R}\right)_{R \geq \tilde{R}} \in C^{\infty}\left(\mathbb{R}^{2 d} ; \mathbb{R}\right)
$$

such that

$$
\begin{equation*}
g(x)\left(\nabla_{x} S_{ \pm, R}, \nabla_{x} S_{ \pm, R}\right)=|\xi|^{2}, \quad(x, \xi) \in \Gamma^{ \pm}(R, J, \sigma) \tag{3.1}
\end{equation*}
$$

and, such that for every $(\alpha, \beta) \in \mathbb{N}^{2 d}$, there exists $C_{\alpha \beta}$ such that for every $(x, \xi) \in$ $\mathbb{R}^{2 d}$, every $R \geq 2 \tilde{R}$,

$$
\left|\partial_{x}^{\alpha} \partial_{\xi}^{\beta}\left(S_{ \pm, R}(x, \xi)-\langle x, \xi\rangle\right)\right| \leq C_{\alpha \beta} \min \left(\langle x\rangle^{1-\rho-|\alpha|}, R^{1-\rho-|\alpha|}\right) .
$$

Proof. Let $J_{0}$ be an open interval such that $\left.J \Subset J_{0} \Subset\right] 0,+\infty[$. Define a smooth function $\theta_{J, J_{0}}$ such that

$$
\theta_{J, J_{0}}(x)= \begin{cases}1 & \text { when } \\ 0 & x \in J \\ 0 & \text { when } \\ x \notin J_{0}\end{cases}
$$

Let us next fix $\sigma_{0}$ such that $\left.\sigma_{0} \in\right] \sigma, 1\left[\right.$. Let $\kappa_{\sigma, \sigma_{0}}$ be a monotone smooth function such that

$$
\kappa_{\sigma, \sigma_{0}}(x)= \begin{cases}1 & \text { when } \quad x \geq-\sigma, \\ 0 & \text { when } \quad x \leq-\sigma_{0} .\end{cases}
$$

Next, let $\chi \in C^{\infty}$ be monotone and such that

$$
\chi(x)= \begin{cases}1 & \text { when } \quad x \geq 1 \\ 0 & \text { when } \quad x \leq 1 / 2\end{cases}
$$

For $R \geq 1$, we define $\chi_{R}(x)$ as

$$
\chi_{R}(x)=\chi\left(\frac{|x|}{R}\right) .
$$

Then for $\tilde{R} \geq 1$ and $R \geq 2 \tilde{R}$ the function $\psi_{R}(x, \xi)$ defined as

$$
\psi_{R}^{ \pm}(x, \xi)=\chi_{R}(x) \theta_{J, J_{0}}\left(|\xi|^{2}\right) \kappa_{\sigma, \sigma_{0}}\left( \pm \frac{\langle x, \xi\rangle}{|x||\xi|}\right)
$$

satisfies the properties

$$
\operatorname{supp} \psi_{R}^{ \pm} \subset \Gamma^{ \pm}\left(R / 2, J_{0}, \sigma_{0}\right) \subset \Gamma^{ \pm}\left(\tilde{R}, J_{0}, \sigma_{0}\right)
$$

and $\psi_{R}^{ \pm}(x, \xi)=1$ for $(x, \xi) \in \Gamma^{ \pm}(R, J, \sigma)$. The functions $\psi_{R}^{ \pm}(x, \xi)$ also enjoy the bounds

$$
\begin{equation*}
\left|\partial_{x}^{\alpha} \partial_{\xi}^{\beta} \psi_{R}^{ \pm}(x, \xi)\right| \leq C_{\alpha \beta N} \min \left(\langle x\rangle^{-|\alpha|}, R^{-|\alpha|}\right)\langle\xi\rangle^{-N} \tag{3.2}
\end{equation*}
$$

with a constant $C_{\alpha \beta N}$ uniform with respect to $R \geq \tilde{R}$.
For $\tilde{R} \gg 1$, let us denote by $\tilde{S}_{ \pm, \tilde{R}}(x, \xi)$ the phase function given by Proposition 3.1 associated to $J_{0}, \sigma_{0}$ and $\tilde{R}$. Then, by invoking (3.2) and the fact that the derivatives of $\psi_{R}^{ \pm}$with respect to $x$ are supported in a set $\{x \gtrsim R\}$, we observe that the phase functions $S_{ \pm, R}$ defined as

$$
S_{ \pm, R}(x, \xi)=\psi_{R}^{ \pm}(x, \xi) \tilde{S}_{ \pm, \tilde{R}}(x, \xi)+\left(1-\psi_{R}^{ \pm}(x, \xi)\right)\langle x, \xi\rangle
$$

satisfies the claimed properties. This completes the proof of Proposition 3.2,
For a given real number $\mu$, we denote by $S(\mu,-\infty)$ the set of smooth functions $a(x, \xi)$ on $\mathbb{R}^{2 d}$ such that for every $N \in \mathbb{N}$, every $(\alpha, \beta) \in \mathbb{N}^{2 d}$ there exists a constant $C_{N \alpha \beta}$ such that for every $(x, \xi) \in \mathbb{R}^{2 d}$,

$$
\left|\partial_{x}^{\alpha} \partial_{\xi}^{\beta} a(x, \xi)\right| \leq C_{N \alpha \beta}\langle x\rangle^{\mu-|\alpha|}\langle\xi\rangle^{-N} .
$$

We equip $S(\mu,-\infty)$ with the natural Fréchet space topology. The next proposition is devoted to the semi-classical Isozaki-Kitada parametrix.

Proposition 3.3. Let us fix an open interval $J \Subset] 0,+\infty[$ and $\sigma \in]-1,1[$. Consider the open intervals $J_{1}$ and $J_{2}$ so that

$$
\left.J \Subset J_{1} \Subset J_{2} \Subset\right] 0,+\infty[
$$

and real numbers

$$
\sigma<\sigma_{1}<\sigma_{2}<1 .
$$

Then there exists $R_{0} \gg 1$ such that for every $N \in \mathbb{N} \backslash\{0\}$, every $k \in \mathbb{N}$, every $R \geq R_{0}$, every $\chi_{ \pm} \in S(-k,-\infty)$ supported in $\Gamma^{ \pm}(R, J, \sigma)$ we can find:

- a sequence

$$
a_{j}^{ \pm} \in S(-j,-\infty), \quad j=0,1, \ldots, N
$$

of smooth functions on $\mathbb{R}^{2 d}$ supported in $\Gamma^{ \pm}\left(R^{1 / 3}, J_{2}, \sigma_{2}\right)$,

- a sequence

$$
b_{j}^{ \pm} \in S(-k-j,-\infty), \quad j=0,1, \ldots, N
$$

of smooth functions on $\mathbb{R}^{2 d}$ supported in $\Gamma^{ \pm}\left(R^{1 / 2}, J_{1}, \sigma_{1}\right)$
such that for every $h \in] 0,1]$, every $\pm t \geq 0$,

$$
\begin{aligned}
e^{-i \frac{t}{h} h^{2} P} \chi_{ \pm}(x, h D)= & J_{S_{ \pm, R^{1 / 4}}}\left(\sum_{j=0}^{N} h^{j} a_{j}^{ \pm}\right) e^{-i \frac{t}{h} h^{2} P_{0}} J_{S_{ \pm, R^{1 / 4}}}\left(\sum_{j=0}^{N} h^{j} b_{j}^{ \pm}\right)^{\star} \\
& +h^{N+1} R_{N}^{ \pm}(t, h)
\end{aligned}
$$

where the phase functions $S_{ \pm, R^{1 / 4}}$ are defined (with $R^{1 / 4}$ instead of $R$ ) in Proposition 3.2, $P_{0}=-\Delta$ denotes the flat Laplacian on $\mathbb{R}^{d}$ and the maps $J_{S_{ \pm, R^{1 / 4}}}(q)$ are defined by

$$
J_{S_{ \pm, R^{1 / 4}}}(q) u(x)=(2 \pi h)^{-d} \iint e^{\frac{i}{h}\left(S_{ \pm, R^{1 / 4}}(x, \xi)-\langle y, \xi\rangle\right)} q(x, \xi) u(y) d y d \xi
$$

Moreover, for every $T>0$, every $N \geq 1$ and every positive integer $s$, there exists $C$ such that for every $h \in] 0,1]$, every $\pm t \in\left[0, T h^{-1}\right]$, the remainder $R_{N}^{ \pm}(t, h)$ satisfies

$$
\begin{equation*}
\left\|(P+1)^{s} R_{N}^{ \pm}(t, h)(P+1)^{s}\right\|_{L^{2}\left(\mathbb{R}^{d}\right) \rightarrow L^{2}\left(\mathbb{R}^{d}\right)} \leq C h^{N-4 s-2} \tag{3.3}
\end{equation*}
$$

Proof. For the precise construction of $a_{j}^{ \pm}$and $b_{j}^{ \pm}$we refer to [28, Section 4], [2]. However, for the convenience of the reader, we recall the main lines of the method in the outgoing case (the incoming one being similar). We first choose $J_{3}, \sigma_{3}$ such that $J_{2} \Subset J_{3}, \sigma_{2}<\sigma_{3}<1$ and then choose $S_{+, R^{1 / 4}} \equiv S_{+}$(for shortness) as in Proposition 3.2 solving the Hamilton-Jacobi equation (3.1) on $\Gamma^{+}\left(R^{1 / 4}, J_{3}, \sigma_{3}\right)$. We then look for a symbol

$$
a^{+}=a_{0}^{+}+h a_{1}^{+}+\cdots+h^{N} a_{N}^{+}
$$

supported in $\Gamma^{+}\left(R^{1 / 3}, J_{2}, \sigma_{2}\right)$ such that

$$
J_{S_{+}}\left(c_{N}(h)\right):=\left(h^{2} P\right) J_{S_{+}}\left(a^{+}\right)-J_{S_{+}}\left(a^{+}\right)\left(h^{2} P_{0}\right)
$$

has a "small contribution" (see (3.6) below). This leads to a system of equations for $a_{0}^{+}, \cdots, a_{N}^{+}$that take the form of rather standard (time independent) transport equations, but only in the region where the Hamilton-Jacobi equation (3.1) is satisfied. For $R$ large enough, we can solve these equations in a neighborhood
of $\Gamma^{+}\left(R^{1 / 3}, J_{2}, \sigma_{2}\right)$ and by cutting off these solutions by a function supported in $\Gamma^{+}\left(R^{1 / 3}, J_{2}, \sigma_{2}\right)$ which equals 1 near $\Gamma^{+}\left(R^{5 / 12}, J_{1}, \tilde{\sigma}_{1}\right)$ (notice that $1 / 3<5 / 12<$ $1 / 2)$ with $\sigma_{1}<\tilde{\sigma}_{1}<\sigma_{2}$, we can build $a_{0}^{+}, \cdots, a_{N}^{+}$so that

$$
\begin{equation*}
c_{N}(h)=h^{N+1} \tilde{r}_{N+1}(h)+\tilde{c}_{N}(h) \tag{3.4}
\end{equation*}
$$

with $\left(\tilde{r}_{N+1}(h)\right)_{0<h \leq 1}$ in a bounded set of $S(-N,-\infty)$ and $\left(\tilde{c}_{N}(h)\right)_{0<h \leq 1}$ in a bounded set of $S(0,-\infty)$, supported in $\Gamma^{+}\left(R^{1 / 3}, J_{2}, \sigma_{2}\right)$ and such that

$$
\begin{equation*}
\tilde{c}_{N}(h) \equiv 0 \quad \text { near } \quad \Gamma^{+}\left(R^{5 / 12}, J_{1}, \tilde{\sigma}_{1}\right) \tag{3.5}
\end{equation*}
$$

Note that the symbol $c_{N}(h)$ is a priori not small, because of the term $\tilde{c}_{N}(h)$. However, using (3.5), we shall see afterwards that the contribution of $J_{S_{+}}\left(\tilde{c}_{N}(h)\right)$ to the final remainder term of the parametrix is harmless, once multiplied from the right by another FIO with nice support properties. Next, remarking that we can choose $a_{0}^{+}$non vanishing near $\Gamma^{+}\left(R^{1 / 2}, J_{1}, \sigma_{1}\right)$, we can solve another family of (algebraic) equations for $b_{0}^{+}, \cdots, b_{N}^{+}$such that

$$
b^{+}:=b_{0}^{+}+\cdots+h^{N} b_{N}^{+}
$$

is supported in $\Gamma^{+}\left(R^{1 / 2}, J_{1}, \sigma_{1}\right)$ and satisfies

$$
J_{S_{+}}\left(a^{+}\right) J_{S^{+}}\left(b^{+}\right)^{\star}=\chi_{+}(x, h D)+h^{N+1} r_{N}(x, h D, h)
$$

with $\left(r_{N}(h)\right)_{0<h \leq 1}$ in a bounded set of $S(-N,-\infty)$. More precisely, the standard composition rule for the computation of the symbol of $J_{S_{+}}\left(a^{+}\right) J_{S^{+}}\left(b^{+}\right)^{\star}$ (see [26]) show that the above equation leads to a triangular system with unknown $b_{0}^{+}, \ldots, b_{N}^{+}$ and $a_{0}^{+}$on the diagonal.

Combining this last equation and the fact that

$$
\begin{align*}
& e^{-i \frac{t}{h} h^{2} P} J_{S_{+}}\left(a_{+}\right)-J_{S_{+}}\left(a_{+}\right) e^{-i \frac{t}{h} h^{2} P_{0}}  \tag{3.6}\\
&=-\frac{i}{h} \int_{0}^{t} e^{-i \frac{t-\tau}{h} h^{2} P} J_{S_{+}}\left(c_{N}(h)\right) e^{-i \frac{\tau}{h} h^{2} P_{0}} d \tau
\end{align*}
$$

we can then represent the remainder of the Isozaki-Kitada parametrix as

$$
h^{N+1} R_{N}^{+}(t, h)=I+I I+I I I
$$

where

$$
\begin{aligned}
I & =-h^{N+1} e^{-i \frac{t}{h} h^{2} P} r_{N+1}(x, h D, h) \\
I I & =-i h^{N} \int_{0}^{t} e^{-i \frac{t-\tau}{h} h^{2} P} J_{S_{+}}\left(\tilde{r}_{N+1}(h)\right) e^{-i \frac{\tau}{h} h^{2} P_{0}} d \tau \circ J_{S_{+}}\left(b^{+}\right)^{\star} \\
I I I & =-\frac{i}{h} \int_{0}^{t} e^{-i \frac{t-\tau}{h} h^{2} P} K_{h}(\tau) d \tau
\end{aligned}
$$

with $K_{h}(\tau)$ an operator with kernel $\mathcal{K}(x, y, h, \tau)$ satisfying, for all $M \geq 0$ and $\alpha, \beta \in$ $\mathbb{N}^{d}$,

$$
\begin{equation*}
\left|\partial_{x}^{\alpha} \partial_{y}^{\beta} \mathcal{K}(x, y, h, \tau)\right| \leq C_{\alpha \beta M} h^{M}(1+|x|+|y|+\tau)^{-M}, \quad \tau \geq 0 \tag{3.7}
\end{equation*}
$$

Here $K_{h}(\tau)=J_{S_{+}}\left(\tilde{c}_{N}(h)\right) e^{i \frac{\tau}{h} h^{2} P_{0}} J_{S_{+}}\left(b^{+}\right)^{\star}$ contains the contribution of $\tilde{c}_{N}(h)$. The above estimate follows from a non stationary phase argument (see [2, Proposition 2.4.7]), by exploiting the support properties of $\tilde{c}_{N}(h)$ and $b^{+}$. For the sake of completeness, and to somehow prepare the reader to the stationary phase argument used in the next section, we recall how to prove (3.7). This proof can be found in [21] for potential perturbations in the non semi-classical case $(h=1)$. Here we reproduce the proof of [2].

The kernel of $K_{h}(\tau)$ at the point $(x, y)$ is given by an oscillatory integral over a fixed compact set in the $\xi$ variable and the phase function

$$
i h^{-1}\left(\tau|\xi|^{2}+S_{+}(y, \xi)-S_{+}(x, \xi)\right)
$$

More precisely the integration is over those $\xi$ such that

$$
\begin{equation*}
(x, \xi) \in \Gamma^{+}\left(R^{1 / 3}, J_{2}, \sigma_{2}\right) \backslash \Gamma^{+}\left(R^{5 / 12}, J_{1}, \tilde{\sigma}_{1}\right) \text { and }(y, \xi) \in \Gamma^{+}\left(R^{1 / 2}, J_{1}, \sigma_{1}\right) \tag{3.8}
\end{equation*}
$$

Estimate (3.7) follows from standard integrations by parts using the following lemma.
Lemma 3.4. There exist $c>0$ and $R_{0}>1$ such that for every $\tau \geq 0$, every $x, y$, $\xi$ satisfying (3.8), every $R \geq R_{0}$,

$$
\begin{equation*}
\left|\nabla_{\xi}\left(\tau|\xi|^{2}+S_{+}(y, \xi)-S_{+}(x, \xi)\right)\right| \geq c(1+\tau+|x|+|y|) \tag{3.9}
\end{equation*}
$$

Proof. For simplicity, we use the notation $\cos (x, \xi)=\langle x, \xi\rangle /|x||\xi|$. We will also use the following statements that are easily checked:

$$
\begin{aligned}
& \cos \left(\eta, \eta^{\prime}\right) \geq-\sigma>-1 \Rightarrow\left|\eta+\eta^{\prime}\right| \geq(1-|\sigma|)^{1 / 2}\left(|\eta|^{2}+\left|\eta^{\prime}\right|^{2}\right)^{1 / 2} \\
& \left|\eta-\eta^{\prime}\right| \leq \epsilon|\eta| \Rightarrow\left|\cos \left(\eta^{\prime \prime}, \eta\right)-\cos \left(\eta^{\prime \prime}, \eta^{\prime}\right)\right| \leq 2 \epsilon, \quad \forall \eta^{\prime \prime} \in \mathbb{R}^{d} \backslash 0
\end{aligned}
$$

One then remark that if (3.8) holds then either $|x| \leq R^{5 / 12}$ or $-\tilde{\sigma}_{1} \geq \cos (x, \xi)>-\sigma_{2}$.
Assume first that $|x| \leq R^{5 / 12}$. Fix $\epsilon$ such that $\sigma_{2}-\sigma_{1}>2 \epsilon$. By choosing $R$ large enough, we have $\left|\nabla_{\xi} S_{+}(y, \xi)-y\right| \leq \epsilon|y|$ and thus $\cos \left(\nabla_{\xi} S_{+}(y, \xi), \xi\right)>-\left|\sigma_{2}\right|$. This implies that

$$
\left|\nabla_{\xi}\left(\tau|\xi|^{2}+S_{+}(y, \xi)\right)\right| \geq\left(1-\left|\sigma_{2}\right|\right)^{1 / 2}\left(4 \tau^{2}|\xi|^{2}+\left|\nabla_{\xi} S_{+}(y, \xi)\right|^{2}\right)^{1 / 2}
$$

from which (3.9) follows easily since $\left|\nabla_{\xi} S_{+}(x, \xi)\right| \lesssim R^{5 / 12}$ and $\left|\nabla_{\xi} S_{+}(y, \xi)\right| \gtrsim R^{1 / 2}$.
We now assume that $-\tilde{\sigma}_{1} \geq \cos (x, \xi)>-\sigma_{2}$. It suffices to show that, if $R$ is large enough, there exists $\sigma \in(-1,1)$ such that

$$
\begin{equation*}
\cos \left(\nabla_{\xi}\left(S_{+}(y, \xi)-S_{+}(x, \xi)\right), \xi\right) \geq \sigma \tag{3.10}
\end{equation*}
$$

$$
\begin{equation*}
\cos \left(-\nabla_{\xi} S_{+}(x, \xi), \nabla_{\xi} S_{+}(y, \xi)\right) \geq \sigma \tag{3.11}
\end{equation*}
$$

provided (3.8) holds. Indeed, the estimate (3.10) implies that the left hand side of (3.9) is bounded from below by $c\left(4 \tau|\xi|^{2}+\left|\nabla_{\xi}\left(S_{+}(x, \xi)-S_{+}(y, \xi)\right)\right|\right)$ and then (3.11) implies that $\left|\nabla_{\xi}\left(S_{+}(x, \xi)-S_{+}(y, \xi)\right)\right| \geq c(|x|+|y|)$.

It remains to prove (3.10) and (3.11). Let us choose $\epsilon>0$ such that $\tilde{\sigma}_{1}-\sigma_{1}>4 \epsilon$. Observe that $\sigma_{1}+2 \epsilon \in(-1,1)$. By choosing $R$ large enough, we have

$$
\begin{equation*}
\cos \left(\nabla_{\xi} S_{+}(y, \xi), \xi\right) \geq-\sigma_{1}-2 \epsilon \quad \text { and } \quad 2 \epsilon-\tilde{\sigma}_{1}>\cos \left(\nabla_{\xi} S_{+}(x, \xi), \xi\right) \tag{3.12}
\end{equation*}
$$

Therefore $\nabla_{\xi} S_{+}(y, \xi) \neq \nabla_{\xi} S_{+}(x, \xi)$ and the left hand side of (3.10) reads

$$
\frac{\left|\nabla_{\xi} S_{+}(y, \xi)\right| \cos \left(\nabla_{\xi} S_{+}(y, \xi), \xi\right)-\left|\nabla_{\xi} S_{+}(x, \xi)\right| \cos \left(\nabla_{\xi} S_{+}(x, \xi), \xi\right)}{\left|\nabla_{\xi} S_{+}(y, \xi)-\nabla_{\xi} S_{+}(x, \xi)\right|} .
$$

Now, using that for every $\alpha \in \mathbb{R}$, every $X, Y \in \mathbb{R}^{d}, X \neq Y$,

$$
\alpha \frac{|Y|-|X|}{|Y-X|} \geq-|\alpha|
$$

we deduce that (3.10) holds with $\sigma=-\left|\sigma_{1}+2 \epsilon\right|$. Finally, we see that (3.11) must hold for some possibly lower $\sigma$ since otherwise we could find sequences $\left(x_{j}\right),\left(y_{j}\right),\left(\xi_{j}\right)$ satisfying (3.8), $-\tilde{\sigma}_{1} \geq \cos \left(x_{j}, \xi_{j}\right)>-\sigma_{2}$ and such that

$$
\lim _{j \rightarrow \infty} \frac{\nabla_{\xi} S_{+}\left(x_{j}, \xi_{j}\right)}{\left|\nabla_{\xi} S_{+}\left(x_{j}, \xi_{j}\right)\right|}=\lim _{j \rightarrow \infty} \frac{\nabla_{\xi} S_{+}\left(y_{j}, \xi_{j}\right)}{\left|\nabla_{\xi} S_{+}\left(y_{j}, \xi_{j}\right)\right|}
$$

which is forbidden by (3.12) and the fact that $-\sigma_{1}-2 \epsilon>2 \epsilon-\tilde{\sigma}_{1}$. This completes the proof of Lemma 3.4.

Therefore, we get the bounds

$$
\left\|(P+1)^{s} r_{N+1}(x, h D, h)(P+1)^{s}\right\|_{L^{2} \rightarrow L^{2}} \leq C h^{-4 s}
$$

and

$$
\left\|(P+1)^{s} K_{h}(\tau)(P+1)^{s}\right\|_{L^{2} \rightarrow L^{2}} \leq C h^{M}
$$

for $\tau \geq 0$ and $h \in] 0,1]$ (in particular, we see a posteriori that the contribution of $\tilde{c}_{N}(h)$ in the remainder of the parametrix is $\mathcal{O}\left(h^{\infty}\right)$ ). Finally, using the $L^{2}$ boundedness of FIO (see e.g. [26]), we have that, for every $a \in S(0,-\infty)$ and every $k \in \mathbb{N}$, there exists $C_{k, a}$ such that

$$
\left\|\left(h^{2} P\right)^{k} J_{S_{+}}(a)\right\|_{L^{2} \rightarrow L^{2}} \leq C_{k, a},
$$

and hence

$$
\begin{aligned}
& \left\|(P+1)^{s} J_{S_{+}}\left(\tilde{r}_{N+1}(h)\right) e^{-i \frac{\tau}{h} h^{2} P} J_{S_{+}}\left(b^{+}\right)^{\star}(P+1)^{s}\right\|_{L^{2} \rightarrow L^{2}} \\
& =\left\|(P+1)^{s} J_{S_{+}}\left(\tilde{r}_{N+1}(h)\right) e^{-i \frac{\tau}{h} h^{2} P}\left((P+1)^{s} J_{S_{+}}\left(b^{+}\right)\right)^{\star}\right\|_{L^{2} \rightarrow L^{2}} \\
& \leq C h^{-4 s}
\end{aligned}
$$

By integrating the corresponding estimates to $I I$ and $I I I$ over an interval of size $T h^{-1}$ we get the result (3.3).
Remark 3.5. Note that the control of the remainder is easier in our case than in [28, 2] since we only need to integrate on $\left[0, T h^{-1}\right]$ in $\tau$ whereas in [28, 2] one has to integrate over $\mathbb{R}^{+}$. Moreover, we do not use any non trapping assumption on the metric: this is the main point in this paper.

## 4. Strichartz estimates outside a large ball

The goal of this section is to prove (1.10). The main point is to prove the following statement.

Proposition 4.1. Let $\varphi \in C_{0}^{\infty}(\mathbb{R} \backslash\{0\})$ and let $g$ be a metric on $\mathbb{R}^{d}$ satisfying (1.8) and (1.9). Then there exists $R>0$ such that for every $T>0$, every $(p, q)$ satisfying (1.4), every $\chi \in C_{0}^{\infty}\left(\mathbb{R}^{d}\right), \chi \equiv 1$ for $|x|<R$, there exists $C>0$ such that for every $f \in L^{2}\left(\mathbb{R}^{d}\right)$, every $\left.\left.h \in\right] 0,1\right]$,

$$
\begin{equation*}
\left\|(1-\chi) e^{-i t P} \varphi\left(h^{2} P\right) f\right\|_{L^{p}\left([-T, T] ; L^{q}\left(\mathbb{R}^{d}\right)\right)} \leq C\|f\|_{L^{2}\left(\mathbb{R}^{d}\right)} \tag{4.1}
\end{equation*}
$$

Remark 4.2. We could have spectrally localized $f$ in the right hand-side of (4.1), i.e. one could have replaced $\|f\|_{L^{2}\left(\mathbb{R}^{d}\right)}$ by $\left\|\varphi\left(h^{2} P\right) f\right\|_{L^{2}\left(\mathbb{R}^{d}\right)}$. Indeed, let $\tilde{\varphi} \in C_{0}^{\infty}(\mathbb{R})$ be equal to one on the support of $\varphi$. Then $\varphi=\varphi \tilde{\varphi}$ and we can apply (4.1) with $\tilde{\varphi}$ instead of $\varphi$.
Proof. Recall that we denote by $P$ the self adjoint realization of $-\Delta_{g}$ on $L^{2}\left(\mathbb{R}^{d}\right)$. For $\chi \in C_{0}^{\infty}\left(\mathbb{R}^{d}\right), \chi \equiv 1$ for $|x|<R$, a partition of unity argument and Proposition 2.5 allow to write

$$
(1-\chi) \varphi\left(h^{2} P\right)=\sum_{k=0}^{N} h^{k}\left(\theta_{k}^{+}(x, h D)+\theta_{k}^{-}(x, h D)\right)+h^{N+1} R_{N, \chi}(h)
$$

where $\theta_{k}^{ \pm} \in S(-k,-\infty)$ and for some $\left.J \Subset\right] 0,+\infty\left[, \sigma_{ \pm} \in\right]-1,1[$,

$$
\operatorname{supp} \theta_{k}^{ \pm} \subset \Gamma^{ \pm}\left(R, J, \sigma_{ \pm}\right)
$$

More precisely, $\sigma_{ \pm}$and $J$ should be such that

$$
\{(x, \xi):|x|>R, p(x, \xi) \in \operatorname{supp}(\varphi)\} \subset \Gamma^{+}\left(R, J, \sigma_{+}\right) \cup \Gamma^{-}\left(R, J, \sigma_{-}\right)
$$

Let us notice that $\sigma_{ \pm}$can be taken both $1 / 2$.
Furthermore the remainder is such that for every $s \geq 0$ there exists $C>0$ such that for every $h \in] 0,1]$,

$$
\left\|(P+1)^{s / 2} R_{N, \chi}(h)(P+1)^{s / 2}\right\|_{L^{2} \rightarrow L^{2}} \leq C h^{-2 s} .
$$

In addition by the elementary properties of the $h$ pseudo differential calculus (cf. e.g. [26]), we can also write

$$
(1-\chi) \varphi\left(h^{2} P\right)=\sum_{k=0}^{N} h^{k}\left(\chi_{k}^{+}(x, h D)^{\star}+\chi_{k}^{-}(x, h D)^{\star}\right)+h^{N+1} \tilde{R}_{N, \chi}(h)
$$

where $\chi_{k}^{ \pm}$and $\tilde{R}_{N, \chi}(h)$ have similar properties to $\theta_{k}^{ \pm}$and $R_{N, \chi}(h)$ respectively. Using the Sobolev embedding, by taking $N$ large enough, we get the bound

$$
\left\|h^{N+1} \tilde{R}_{N, \chi}(h) e^{-i t P} f\right\|_{L_{T}^{p} L^{q}} \leq C\|f\|_{L^{2}},
$$

provided $(p, q)$ is satisfying (1.4). Therefore it suffices to prove the bound

$$
\begin{equation*}
\left\|\chi_{k}^{ \pm}(x, h D)^{\star} e^{-i t P} f\right\|_{L_{T}^{p} L^{q}} \leq C\|f\|_{L^{2}} . \tag{4.2}
\end{equation*}
$$

Since $\chi_{k}^{ \pm}(x, h D)^{\star} e^{-i t P}$ are clearly $L^{2}$ bounded, uniformly in $h$ and $t$, thanks to the Keel-Tao theorem (see [23], [8, Proposition 2.8]), to get (4.2), it suffices to prove the dispersive inequality

$$
\left\|\chi_{k}^{ \pm}(x, h D)^{\star} e^{-i t P} e^{i s P} \chi_{k}^{ \pm}(x, h D) f\right\|_{L^{\infty}} \leq \frac{C}{|t-s|^{d / 2}}\|f\|_{L^{1}\left(\mathbb{R}^{d}\right)}, \quad t, s \in[-T, T]
$$

uniformly with respect to $h$. By the time rescaling $t \mapsto h t$, and by defining the maps

$$
U_{k, h}^{ \pm}(t)=\chi_{k}^{ \pm}(x, h D)^{\star} e^{-i h t P}, \quad t \in\left[-h^{-1} T, h^{-1} T\right]
$$

it suffices to prove the dispersive inequality

$$
\begin{equation*}
\left\|U_{k, h}^{ \pm}(t)\left(U_{k, h}^{ \pm}(s)\right)^{\star} f\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)} \leq \frac{C}{(h|t-s|)^{d / 2}}\|f\|_{L^{1}\left(\mathbb{R}^{d}\right)}, \quad t, s \in\left[-h^{-1} T, h^{-1} T\right] \tag{4.3}
\end{equation*}
$$

Clearly

$$
\begin{equation*}
U_{k, h}^{ \pm}(t)\left(U_{k, h}^{ \pm}(s)\right)^{\star}=\chi_{k}^{ \pm}(x, h D)^{\star} e^{-i h(t-s) P} \chi_{k}^{ \pm}(x, h D) \tag{4.4}
\end{equation*}
$$

Denote by $K_{ \pm}(t-s, x, y, h)$ (we do not explicit the dependence on $k$ ) the kernel of (4.4). In order to prove (4.3), it is sufficient to show that there exists $C>0$ such that for every $h \in] 0,1]$, every $x, y \in \mathbb{R}^{d}$, every $t, s \in\left[-h^{-1} T, h^{-1} T\right], t \neq s$,

$$
\begin{equation*}
\left|K_{ \pm}(t-s, x, y, h)\right| \leq C|h(t-s)|^{-d / 2} \tag{4.5}
\end{equation*}
$$

The next lemma contains the main trick in our analysis.

Lemma 4.3. For $K_{+}$, it suffices to prove (4.5) for $t-s>0$. Similarly, for $K_{-}$, it suffices to prove (4.5) for $t-s<0$.

Proof. We only consider $K_{+}$, the analysis for $K_{-}$being similar. Suppose that (4.5) holds true for $t-s>0$. Let $t, s \in\left[-h^{-1} T, h^{-1} T\right]$ such that $t-s<0$. Since

$$
U_{k, h}^{+}(t)\left(U_{k, h}^{+}(s)\right)^{\star}=\left(U_{k, h}^{+}(s)\left(U_{k, h}^{+}(t)\right)^{\star}\right)^{\star}
$$

we obtain that

$$
K_{+}(t-s, x, y, h)=\overline{K_{+}(s-t, y, x, h)} .
$$

Since $s-t>0$, our assumption that (4.5) holds for positive values of the first argument of $K_{+}$implies that

$$
\left|K_{+}(t-s, x, y, h)\right|=\left|\overline{K_{+}(s-t, y, x, h)}\right| \leq C(h(s-t))^{-d / 2}
$$

which completes the proof of Lemma 4.3.
It is now clear that the proof of (4.5), and thus of Proposition 4.1, will be finished, once we establish the following lemma.

Lemma 4.4. There exists $R \gg 1$ and $C>0$ such that for every $h \in] 0,1]$, every $\left.\pm t \in] 0, h^{-1} T\right]$, every $x, y \in \mathbb{R}^{d}$,

$$
\left|K_{ \pm}(t, x, y, h)\right| \leq C( \pm h t)^{-d / 2} .
$$

Proof. As before, we only consider the case of $K_{+}$. Denote by $\mathcal{K}_{+}(t, x, y, h)$ the kernel of $\exp (-i t h P) \chi_{k}^{+}(x, h D)$ and by $\tilde{\mathcal{K}}_{+}(x, y, h)$ the kernel of $\chi_{k}^{+}(x, h D)^{\star}$. Clearly $\tilde{\mathcal{K}}_{+}(x, y, h)$ satisfies the assumptions of the Schur lemma uniformly in $h$, and therefore by writing

$$
K_{+}(t, x, y, h)=\int_{\mathbb{R}^{d}} \tilde{\mathcal{K}}_{+}(x, z, h) \mathcal{K}_{+}(t, z, y, h) d z,
$$

we infer that

$$
\left|K_{+}(t, x, y, h)\right| \leq \sup _{z \in \mathbb{R}^{d}}\left|\mathcal{K}_{+}(t, z, y, h)\right| \int_{\mathbb{R}^{d}}\left|\tilde{\mathcal{K}}_{+}(x, z, h)\right| d z \leq C \sup _{z \in \mathbb{R}^{d}}\left|\mathcal{K}_{+}(t, z, y, h)\right| .
$$

Therefore, it suffices to prove that there exists $C>0$ such that for every $h \in] 0,1]$, every $\left.t \in] 0, h^{-1} T\right]$, every $x, y \in \mathbb{R}^{d}$,

$$
\begin{equation*}
\left|\mathcal{K}_{+}(t, x, y, h)\right| \leq C(h t)^{-d / 2} . \tag{4.6}
\end{equation*}
$$

In order to prove (4.6), we will of course use Proposition 3.3. With the notations of Proposition 3.3, since for $N \gg 1$, the map $R_{N}(t, h)$ is bounded, uniformly with respect to $t, h$, from $H^{-s}$ to $H^{s}$ with $s>d / 2$, we deduce that its kernel is bounded
uniformly with respect to $h \in] 0,1], t \in\left[0, h^{-1} T\right]$. Therefore, in view of Proposition 3.3. by expressing the kernel of

$$
J_{S_{+, R^{1 / 4}}}\left(a_{j_{1}}^{+}\right) e^{-i \frac{t}{h} h^{2} P_{0}} J_{S_{+, R^{1 / 4}}}\left(b_{j_{2}}^{+}\right)^{\star}, \quad 0 \leq j_{1}, j_{2} \leq N,
$$

estimates (4.6) will be established once we prove that there exists $C>0$ such that for every $h \in] 0,1]$, every $\left.t \in] 0, h^{-1} T\right]$, every $x, y \in \mathbb{R}^{d}$,

$$
\begin{equation*}
(2 \pi h)^{-d}\left|\int_{\mathbb{R}^{d}} e^{\frac{i}{h} \Phi_{+}(t, R, x, y, \xi)} a(x, \xi) \overline{b(y, \xi)} d \xi\right| \leq C(h t)^{-d / 2} \tag{4.7}
\end{equation*}
$$

where the phase $\Phi_{+}$is defined as

$$
\Phi_{+}(t, R, x, y, \xi)=S_{+, R^{1 / 4}}(x, \xi)-S_{+, R^{1 / 4}}(y, \xi)-t|\xi|^{2}
$$

and $a(x, \xi), b(x, \xi)$ are fixed smooth functions supported in

$$
\left\{(x, \xi) \in \mathbb{R}^{2 d}:|x| \geq R \gg 1,0<c \leq|\xi| \leq C\right\}
$$

for some positive constants $c$ and $C$. In the proof of (4.7), we will consider two different regimes for $t$. If $t \in[0, h]$ then, using the support property with respect to $\xi$, the left hand-side of (4.7) can be estimate by $C h^{-d}$ which, for $t \leq h$, is bounded by $C(h t)^{-d / 2}$. Therefore, we can suppose that $t \geq h$ in (4.7). In this case we will take advantage of the rapid oscillations of $\exp \left(i h^{-1} \Phi_{+}\right)$. When $t \geq h$, the natural big parameter is $t h^{-1}$. We thus set

$$
\frac{1}{h} \Phi_{+}(t, R, x, y, \xi)=\frac{t}{h} \tilde{\Phi}_{+}(t, R, x, y, \xi),
$$

where

$$
\tilde{\Phi}_{+}(t, R, x, y, \xi)=\frac{S_{+, R^{1 / 4}}(x, \xi)-S_{+, R^{1 / 4}}(y, \xi)}{t}-|\xi|^{2}
$$

We can write

$$
\tilde{\Phi}_{+}(t, R, x, y, \xi)=\int_{0}^{1}\left\langle\nabla_{x} S_{+, R^{1 / 4}}(y+\alpha(x-y), \xi), x-y\right\rangle \frac{d \alpha}{t}-|\xi|^{2} .
$$

Therefore, thanks to the properties of the phase function $S_{+, R^{1 / 4}}$ displayed in Proposition 3.2, we obtain that

$$
\nabla_{\xi} \tilde{\Phi}_{+}(t, R, x, y, \xi)=\frac{x-y}{t}-2 \xi+Q(R, x, y, \xi) \cdot \frac{x-y}{t}
$$

where $Q$ is a $d \times d$ matrix satisfying the bound

$$
\left|\partial_{\xi}^{\beta} Q(R, x, y, \xi)\right| \leq C_{\beta} R^{-\rho}
$$

for $R \gg 1, x, y \in \mathbb{R}^{d}, \xi$ on the support of $a(x, \xi) b(y, \xi)$.

Therefore, there exist $R_{0} \gg 1$ and $C_{0}>0$ such that for $R>R_{0}$ and $\left|\frac{x-y}{t}\right| \geq C_{0}$, we have that, for $|\xi| \in[c, C]$,

$$
\begin{equation*}
\left|\nabla_{\xi} \tilde{\Phi}_{+}(t, R, x, y, \xi)\right| \gtrsim\left|\frac{x-y}{t}\right| \tag{4.8}
\end{equation*}
$$

and

$$
\begin{equation*}
\left|\partial_{\xi}^{\beta} \tilde{\Phi}_{+}(t, R, x, y, \xi)\right| \leq C_{\beta}\left|\frac{x-y}{t}\right|, \quad|\beta| \geq 2 . \tag{4.9}
\end{equation*}
$$

Therefore using (4.8), (4.9) and integration by parts with respect to $\xi$, we deduce that for every $\Lambda \geq 1$, the left hand-side of (4.7) is bounded by

$$
C h^{-d} C_{\Lambda}\left(t h^{-1}\right)^{-\Lambda} \leq C(h t)^{-d / 2}
$$

provided $\Lambda$ is taken bigger than $d / 2$.
We can therefore suppose that $t, x, y$ are such that $t \geq h$ and

$$
\begin{equation*}
\left|\frac{x-y}{t}\right| \leq C_{0} . \tag{4.10}
\end{equation*}
$$

In this case, we evaluate the left hand-side of (4.7) by the stationary phase. Under the condition (4.10), we can write

$$
\nabla_{\xi}^{2} \tilde{\Phi}_{+}(t, R, x, y, \xi)=-2 \operatorname{Id}+\nabla_{\xi} Q(R, x, y, \xi) \cdot \frac{x-y}{t}=-2 \operatorname{Id}+\mathcal{O}\left(R^{-\rho}\right)
$$

Therefore for $R \gg 1$, the map

$$
\xi \longmapsto \nabla_{\xi} \tilde{\Phi}_{+}(t, R, x, y, \xi)
$$

is a diffeomorphism from $\mathbb{R}^{d}$ to $\mathbb{R}^{d}$. In particular, for fixed $t \geq h, x, y$ satisfying (4.10), the phase $\tilde{\Phi}_{+}$has a unique non degenerate critical point $\xi_{c r}(t, R, x, y)$. Moreover, thanks to (4.10), for $|\beta| \geq 1$,

$$
\left|\partial_{\xi}^{\beta} \tilde{\Phi}_{+}\left(t, R, x, y, \xi_{c r}(t, R, x, y)\right)\right| \leq C_{\beta} .
$$

We can therefore apply the stationary phase estimate to conclude that for $t \geq h$ and ( $t, x, y$ ) satisfying (4.10), the left hand-side of (4.7) is bounded by

$$
C h^{-d}\left(t h^{-1}\right)^{-d / 2}=C(h t)^{-d / 2} .
$$

This completes the proof of Lemma 4.4,
This completes the proof of Proposition 4.1
It is now clear that (1.10) will be proved, once we establish the following statement.

Proposition 4.5. Let $g$ be a metric on $\mathbb{R}^{d}$ satisfying (1.8) and (1.9). Then there exists $R>0$ such that for every $T>0$, every $(p, q)$ satisfying (1.4), every $\chi \in$ $C_{0}^{\infty}\left(\mathbb{R}^{d}\right), \chi \equiv 1$ for $|x|<R$, there exists $C>0$ such that for every $f \in L^{2}\left(\mathbb{R}^{d}\right)$,

$$
\left\|(1-\chi) e^{-i t P} f\right\|_{L^{p}\left([-T, T] ; L^{q}\left(\mathbb{R}^{d}\right)\right)} \leq C\|f\|_{L^{2}\left(\mathbb{R}^{d}\right)}
$$

Proof. Set

$$
u=e^{-i t P} f
$$

Using Proposition 2.10, we can write

$$
\begin{equation*}
\|(1-\chi) u\|_{L_{T}^{p} L^{q}} \leq C\|f\|_{L^{2}}+C\left(\sum_{h^{-1}: \text { dyadic }}\left\|\varphi\left(h^{2} P\right)(1-\chi) u\right\|_{L_{T}^{p} L^{q}}^{2}\right)^{\frac{1}{2}} . \tag{4.11}
\end{equation*}
$$

Let $\varphi_{1} \in C_{0}^{\infty}(\mathbb{R} \backslash\{0\})$ be equal to one on the support on $\varphi$. We can write

$$
\begin{aligned}
\varphi\left(h^{2} P\right)(1-\chi) & =\varphi\left(h^{2} P\right) \varphi_{1}\left(h^{2} P\right)(1-\chi) \\
& =\varphi\left(h^{2} P\right)(1-\chi) \varphi_{1}\left(h^{2} P\right)+\varphi\left(h^{2} P\right)\left[\chi, \varphi_{1}\left(h^{2} P\right)\right]
\end{aligned}
$$

Using Proposition 2.9 and Proposition 4.1 (see also Remark 4.2), we have

$$
\begin{equation*}
\left\|\varphi\left(h^{2} P\right)(1-\chi) \varphi_{1}\left(h^{2} P\right) u\right\|_{L_{T}^{p} L^{q}} \leq C\left\|\varphi_{1}\left(h^{2} P\right) f\right\|_{L^{2}} . \tag{4.12}
\end{equation*}
$$

Using the Schur lemma, we get

$$
\left\|\left[\chi, \varphi_{1}\left(h^{2} P\right)\right]\right\|_{L^{2} \rightarrow L^{2}} \leq C h, \quad\left\|\left[\left[\chi, \varphi_{1}\left(h^{2} P\right)\right], \varphi\left(h^{2} P\right)\right]\right\|_{L^{2} \rightarrow L^{2}} \leq C h^{2} .
$$

Thus, using Proposition 2.9, we can write

$$
\begin{aligned}
\left\|\varphi\left(h^{2} P\right)\left[\chi, \varphi_{1}\left(h^{2} P\right)\right] u\right\|_{L_{T}^{p} L^{q}} & \leq C h^{-1}\left\|\varphi\left(h^{2} P\right)\left[\chi, \varphi_{1}\left(h^{2} P\right)\right] u\right\|_{L_{T}^{p} L^{2}} \\
& \leq C h\|u\|_{L_{T}^{p} L^{2}}+C\left\|\varphi\left(h^{2} P\right) f\right\|_{L^{2}} \\
& \leq C h\|f\|_{L^{2}}+C\left\|\varphi\left(h^{2} P\right) f\right\|_{L^{2}} .
\end{aligned}
$$

In view of (4.12) and the last estimate, coming back to (4.11) completes the proof of Proposition 4.5

## 5. Semi-classical time estimates and applications

In this section $g$ is a metric satisfying (2.1), (2.2). The next proposition describes the WKB approximation for solutions of the semi-classical Schrödinger equation for times which are small but independent of the semi-classical parameter. This construction is well known (see e.g. [26]). Here is the precise statement.

Proposition 5.1. Let $a(x, \xi)$ be a smooth function on $\mathbb{R}^{2 d}$ satisfying

$$
\begin{equation*}
\left|\partial_{x}^{\alpha} \partial_{\xi}^{\beta} a(x, \xi)\right| \leq C_{\alpha \beta}, \quad \forall(\alpha, \beta) \in \mathbb{N}^{2 d} \tag{5.1}
\end{equation*}
$$

$$
\begin{equation*}
\exists R>0: a(x, \xi)=0, \quad \text { for } \quad|\xi|>R . \tag{5.2}
\end{equation*}
$$

Then there exists $\alpha>0$, there exists

$$
S(t, x, \xi) \in C^{\infty}\left([-\alpha, \alpha] \times \mathbb{R}^{d} \times \mathbb{R}^{d}\right)
$$

and a sequence of smooth functions $a_{j}(t, x, \xi), j \geq 0$ compactly supported with respect to $\xi$ such that for every $u_{0} \in L^{2}\left(\mathbb{R}^{d}\right)$, every $\left.\left.h \in\right] 0,1\right]$, every $N \in \mathbb{N}$, the solutions of the problem

$$
\left(i h \partial_{t}+h^{2} \Delta_{g}\right) u=0,\left.\quad u\right|_{t=0}=a(x, h D) u_{0}
$$

can be represented, for $t \in[-\alpha, \alpha]$ as

$$
u(t, x)=J_{N}(t) u_{0}+R_{N}(t) u_{0}
$$

where $R_{N}(t)$ satisfies

$$
\begin{equation*}
\left\|R_{N}(t)\right\|_{L^{2}\left(\mathbb{R}^{d}\right) \rightarrow H^{k}\left(\mathbb{R}^{d}\right)} \leq C h^{N+1-k}, \quad 0 \leq k \leq N+1 \tag{5.3}
\end{equation*}
$$

uniformly with respect to $t \in[-\alpha, \alpha]$, and

$$
J_{N}(t) u_{0}=(2 \pi h)^{-d} \int_{\mathbb{R}^{d}} e^{i h^{-1} S(t, x, \xi)}\left[\sum_{j=0}^{N} a_{j}(t, x, \xi) h^{j}\right] \widehat{u_{0}}\left(\frac{\xi}{h}\right) d \xi .
$$

Moreover $S(t, x, \xi)$ is the solution of the Hamilton-Jacobi equation

$$
\partial_{t} S+g(x)\left(\nabla_{x} S, \nabla_{x} S\right)=0,\left.\quad S\right|_{t=0}=x \cdot \xi
$$

and $a_{j}(t, x, \xi)$ are solutions of the transport equations

$$
\partial_{t} a_{0}+2 g(x)\left(\nabla_{x} S, \nabla_{x} a_{0}\right)+\Delta_{g} S a_{0}=0, \quad a_{0}(0, x, \xi)=a(x, \xi)
$$

for $j=0$, and

$$
\partial_{t} a_{j}+2 g(x)\left(\nabla_{x} S, \nabla_{x} a_{j}\right)+\Delta_{g} S a_{j}=i \Delta_{g}\left(a_{j-1}\right), \quad a_{j}(0, x, \xi)=0
$$

for $j \geq 1$. Finally, there exists $C_{N}>0$ such that if $u_{0} \in L^{1}\left(\mathbb{R}^{d}\right)$ then for every $t \in[-\alpha, \alpha]$, every $h \in] 0,1]$,

$$
\left\|J_{N}(t) u_{0}\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)} \leq \frac{C_{N}}{(|t| h)^{d / 2}}\left\|u_{0}\right\|_{L^{1}\left(\mathbb{R}^{d}\right)} .
$$

Proof. The proof of Proposition 5.1 is given in [8] when $a(x, \xi)$ is supported in a coordinate patch of the cotangent bundle of a compact manifold. The analysis in the case here is slightly more delicate since the $L^{2}$ bound of the remainder is not straightforward as in [8]. However, using that for $|t| \leq \alpha \ll 1$ one has

$$
\left|\nabla_{x} \nabla_{\xi} S(t, x, \xi)-\mathrm{Id}\right| \ll 1,
$$

we can apply the standard result for $L^{2}$ boundedness of FIO (see e.g. [26]) from which (5.3) follows.

After the time rescaling $t \mapsto h t$, as in [8] an application of the Keel-Tao theorem [23] gives the following Strichartz inequalities (homogeneous and non homogeneous) for the Schrödinger equation on semi-classical time intervals.

Proposition 5.2. Let $\varphi \in C_{0}^{\infty}(\mathbb{R})$. Then there exist $\alpha>0$ and $C>0$ such that for every interval $J \subset \mathbb{R}$ of size $\leq \alpha h, h \in] 0,1]$, if $u$ solves

$$
\left(i \partial_{t}-P\right) u=0,\left.\quad u\right|_{t=0}=\varphi\left(h^{2} P\right) u_{0}, \quad u_{0} \in L^{2}\left(\mathbb{R}^{d}\right)
$$

then

$$
\|u\|_{L^{p}\left(J ; L^{q}\left(\mathbb{R}^{d}\right)\right)} \leq C\left\|\varphi\left(h^{2} P\right) u_{0}\right\|_{L^{2}\left(\mathbb{R}^{d}\right)},
$$

provided ( $p, q$ ) satisfies (1.4).
Moreover, if u solves

$$
\begin{equation*}
\left(i \partial_{t}-P\right) u=\varphi\left(h^{2} P\right) f,\left.\quad u\right|_{t=0}=0 \tag{5.4}
\end{equation*}
$$

then

$$
\begin{equation*}
\|u\|_{L^{p}\left(J ; L^{q}\left(\mathbb{R}^{d}\right)\right)} \leq C\left\|\varphi\left(h^{2} P\right) f\right\|_{L^{p_{1}}\left(J ; L^{q_{1}}\left(\mathbb{R}^{d}\right)\right)}, \tag{5.5}
\end{equation*}
$$

provided $(p, q)$ and $\left(\frac{p_{1}}{p_{1}-1}, \frac{q_{1}}{q_{1}-1}\right)$ satisfy (1.4).
As in [8], Proposition 5.2 yields a Strichartz inequality, with derivative loss in classical Sobolev spaces.

Proposition 5.3. Let $T>0$. Then there exists $C>0$ such that if $u$ solves

$$
\left(i \partial_{t}-P\right) u=0,\left.\quad u\right|_{t=0}=u_{0}, \quad u_{0} \in H^{\frac{1}{p}}\left(\mathbb{R}^{d}\right)
$$

then

$$
\|u\|_{L_{T}^{p} L^{q}\left(\mathbb{R}^{d}\right)} \leq C\left\|u_{0}\right\|_{H^{\frac{1}{p}}\left(\mathbb{R}^{d}\right)},
$$

provided ( $p, q$ ) satisfies (1.4).
Proof. Thanks to Proposition [2.10, it suffices to prove that for every $\varphi \in C_{0}^{\infty}(\mathbb{R})$ there exists $C>0$ such that for every $h \in] 0,1]$, every $f \in L^{2}\left(\mathbb{R}^{d}\right)$,

$$
\left\|\exp (-i t P) \varphi\left(h^{2} P\right) f\right\|_{L_{T}^{p} L^{q}} \leq C h^{-\frac{1}{p}}\|f\|_{L^{2}}
$$

We split the interval $[-T, T]$ into $C h^{-1}$ intervals of size $\alpha h$, where $\alpha$ is the real number involved in the statement of Proposition 5.1. Using the $L^{2}$ boundedness of $\exp (-i t P)$, and applying, according to the above splitting, about $C h^{-1}$ times Proposition 5.2 yields

$$
\left\|\exp (-i t P) \varphi\left(h^{2} P\right) f\right\|_{L_{T}^{p} L^{q}}^{p} \leq C h^{-1}\|f\|_{L^{2}}^{p}
$$

which completes the proof of Proposition 5.3.

The proof of Theorem $\dagger$ is now complete.
Next, we state a non homogeneous extension of Proposition 5.2.
Proposition 5.4. Let $\varphi \in C_{0}^{\infty}(\mathbb{R})$ and $T>0$. Then there exists $C>0$ such that if $u$ is a solution of

$$
\left.\left.i u_{t}-P u=\varphi\left(h^{2} P\right) f, \quad h \in\right] 0,1\right], \quad f \in L^{2}\left([0, T] \times \mathbb{R}^{d}\right)
$$

with initial data

$$
\left.u\right|_{t=0}=\varphi\left(h^{2} P\right) u_{0}, \quad u_{0} \in L^{2}\left(\mathbb{R}^{d}\right)
$$

then

$$
\begin{equation*}
\|u\|_{L_{T}^{p} L^{q}} \leq C\|u\|_{L_{T}^{\infty} L^{2}}+C h^{-1 / 2}\|u\|_{L_{T}^{2} L^{2}}+C h^{1 / 2}\left\|\varphi\left(h^{2} P\right) f\right\|_{L_{T}^{2} L^{2}}, \tag{5.6}
\end{equation*}
$$

provided ( $p, q$ ) satisfies (1.4).
Moreover, if $d=3$, the following estimate holds

$$
\begin{equation*}
\|u\|_{L_{T}^{2} L^{6}} \leq C\|u\|_{L_{T}^{\infty} L^{2}}+C h^{-1 / 2}\|u\|_{L_{T}^{2} L^{2}}+C\left\|\varphi\left(h^{2} P\right) f\right\|_{L_{T}^{2} L^{6 / 5}} . \tag{5.7}
\end{equation*}
$$

Remark 5.5. By taking $f=0$ and $T \sim h$, we observe that Proposition 5.2 is a particular case of (5.6).
Remark 5.6. For $d \geq 4$, an estimate analogous to (5.7) holds. More precisely, one has to replace 6 by $\frac{2 d}{d-2}$ and $6 / 5$ by $\frac{2 d}{d+2}$.
Proof of Proposition 5.4. Clearly we can restrict our considerations to the interval $[0, T]$, the analysis on $[-T, 0]$ being analogous (the sign of $t$ in this discussion is harmless). Write

$$
\begin{equation*}
u(t)=e^{-i t P} \varphi\left(h^{2} P\right) u_{0}-i \int_{0}^{t} e^{-i(t-\tau) P} \varphi\left(h^{2} P\right) f(\tau) d \tau \tag{5.8}
\end{equation*}
$$

Observe that $u \in C\left([-T, T] ; L^{2}\left(\mathbb{R}^{d}\right)\right)$. If $T \leq \alpha h$, then Proposition 5.2 and the triangle inequality give

$$
\begin{aligned}
\|u\|_{L_{T}^{p} L^{q}} & \leq C\left\|\varphi\left(h^{2} P\right) u_{0}\right\|_{L^{2}}+C\left\|\varphi\left(h^{2} P\right) f\right\|_{L_{T}^{1} L^{2}} \\
& \leq C\|u\|_{L_{T}^{\infty} L^{2}}+C h^{1 / 2}\left\|\varphi\left(h^{2} P\right) f\right\|_{L_{T}^{2} L^{2}},
\end{aligned}
$$

provided $(p, q)$ satisfies (1.4). Hence we can suppose that $T \geq \alpha h$. Consider a splitting of $[0, T]$ :

$$
[0, T]=[0, a] \cup J_{1} \cup \cdots \cup J_{k} \cup[b, T],
$$

where for $j=1, \ldots k$, there exists $c_{j}$ such that

$$
J_{j}=\left[c_{j}-\frac{\alpha h}{8}, c_{j}+\frac{\alpha h}{8}\right] \subset\left[\frac{\alpha h}{8}, T-\frac{\alpha h}{8}\right] .
$$

We also suppose that $a \leq \alpha h$ and $T-b \leq \alpha h$. We may also suppose that

$$
c_{1}<c_{2}<\cdots<c_{k} .
$$

Clearly $k \lesssim h^{-1}$. Coming back to the Duhamel formula (5.8), using Proposition 5.2, we obtain that for $(p, q)$ satisfying (1.4),

$$
\begin{equation*}
\|u\|_{L^{p}\left([0, a] ; L^{q}\right)} \leq C\|u\|_{L_{T}^{\infty} L^{2}}+C h^{1 / 2}\left\|\varphi\left(h^{2} P\right) f\right\|_{L^{2}\left([0, a] \times \mathbb{R}^{d}\right)} . \tag{5.9}
\end{equation*}
$$

Similarly, we estimate the contribution of $[b, T]$ by writing,

$$
\begin{equation*}
\|u\|_{L^{p}\left([b, T] ; L^{q}\right)} \leq C\|u\|_{L_{T}^{\infty} L^{2}}+C h^{1 / 2}\left\|\varphi\left(h^{2} P\right) f\right\|_{L^{2}\left([b, T] \times \mathbb{R}^{d}\right)} . \tag{5.10}
\end{equation*}
$$

We next define the intervals

$$
J_{j}^{\prime}:=J_{j}+\left[-\frac{\alpha h}{8}, \frac{\alpha h}{8}\right], \quad j=1, \ldots k .
$$

Observe that $J_{j}^{\prime} \subset[0, T]$. Let us fix $\psi \in C_{0}^{\infty}(\mathbb{R})$ such that $\psi=1$ on $[-1 / 8,1 / 8]$ and $\operatorname{supp} \psi \subset[-1 / 4,1 / 4]$. Set

$$
\psi_{j}(t):=\psi\left(\frac{t-c_{j}}{\alpha h}\right), \quad j=1, \ldots k
$$

and

$$
u_{j}(t):=\psi_{j}(t) u(t) .
$$

Observe that $u_{j}(t)=u(t)$ for $t \in J_{j}$ and that $u_{j}$ solves the equation

$$
\left(i \partial_{t}-P\right) u_{j}=i \psi_{j}^{\prime} u+\psi_{j} \varphi\left(h^{2} P\right) f, \quad u_{j}(0)=0 .
$$

Hence, by writing the Duhamel formula for $u_{j}$, using Proposition 5.2 and the triangle inequality, we get for $(p, q)$ satisfying (1.4),

$$
\begin{aligned}
\|u\|_{L^{p}\left(J_{j}, L^{q}\right)} & \leq\left\|u_{j}\right\|_{L^{p}\left(J_{j}^{\prime} ; L^{q}\right)} \\
& \leq C h^{-1}\|u\|_{L^{1}\left(J_{j}^{\prime} ; L^{2}\right)}+C\left\|\varphi\left(h^{2} P\right) f\right\|_{L^{1}\left(J_{j}^{\prime} ; L^{2}\right)} \\
& \leq C h^{-1 / 2}\|u\|_{L^{2}\left(J_{j}^{\prime} \times \mathbb{R}^{d}\right)}+C h^{1 / 2}\left\|\varphi\left(h^{2} P\right) f\right\|_{L^{2}\left(J_{j}^{\prime} \times \mathbb{R}^{d}\right)} .
\end{aligned}
$$

Summing over $j=1, \ldots k$, since $p \geq 2$, we get

$$
\begin{aligned}
\|u\|_{L^{p}\left([a, b] ; L^{q}\right)}^{p} \leq C h^{-\frac{p}{2}} \sum_{j=1}^{k}\|u\|_{L^{2}\left(J_{j}^{\prime} \times \mathbb{R}^{d}\right)}^{p} & +C h^{\frac{p}{2}} \sum_{j=1}^{k}\left\|\varphi\left(h^{2} P\right) f\right\|_{L^{2}\left(J_{j}^{\prime} \times \mathbb{R}^{d}\right)}^{p} \\
\leq C h^{-\frac{p}{2}}\left(\sum_{j=1}^{k}\|u\|_{L^{2}\left(J_{j}^{\prime} \times \mathbb{R}^{d}\right)}^{2}\right)^{\frac{p}{2}} & \left.+C h^{\frac{p}{2}}\left(\sum_{j=1}^{k}\left\|\varphi\left(h^{2} P\right) f\right\|_{L^{2}\left(J_{j}^{\prime} \times \mathbb{R}^{d}\right)}^{2}\right)\right)^{\frac{p}{2}} \\
& \leq C h^{-\frac{p}{2}}\|u\|_{L_{T}^{2} L^{2}}^{p}+C h^{\frac{p}{2}}\left\|\varphi\left(h^{2} P\right) f\right\|_{L_{T}^{2} L^{2}}^{p} .
\end{aligned}
$$

Hence

$$
\|u\|_{L^{p}\left([a, b] ; L^{q}\right)} \leq C h^{-\frac{1}{2}}\|u\|_{L_{T}^{2} L^{2}}+C h^{\frac{1}{2}}\left\|\varphi\left(h^{2} P\right) f\right\|_{L_{T}^{2} L^{2}} .
$$

Coming back to (5.9) and (5.10) completes the proof of (5.6).
Let us now turn to the proof of (5.7). It follows along the same lines as the proof of (5.6). Indeed, with the above notations, using (5.5) with $d=3$ and $p=p_{1}=2$, we obtain that

$$
\begin{aligned}
\|u\|_{L^{2}\left(J_{j}, L^{6}\right)} & \leq\left\|u_{j}\right\|_{L^{2}\left(J_{j}^{\prime} ; L^{6}\right)} \\
& \leq C h^{-1}\|u\|_{L^{1}\left(J_{j}^{\prime} ; L^{2}\right)}+C\left\|\varphi\left(h^{2} P\right) f\right\|_{L^{2}\left(J_{j}^{\prime} ; L^{6 / 5}\right)} \\
& \leq C h^{-1 / 2}\|u\|_{L^{2}\left(J_{j}^{\prime} \times \mathbb{R}^{d}\right)}+C\left\|\varphi\left(h^{2} P\right) f\right\|_{L^{2}\left(J_{j}^{\prime} ; L^{6 / 5}\right)} .
\end{aligned}
$$

Squaring and summing over $J_{j}$ gives

$$
\|u\|_{L^{2}\left([a, b] ; L^{6}\right)} \leq C h^{-1 / 2}\|u\|_{L_{T}^{2} L^{2}}+C\left\|\varphi\left(h^{2} P\right) f\right\|_{L_{T}^{2} L^{6 / 5}}
$$

Similar estimates holds on $[0, a]$ and $[b, T]$ which ends the proof of (5.7). This completes the proof of Proposition 5.4.

## 6. Using the non trapping assumption

The proof of Theorem 2 will be completed, once we prove the following statement.
Theorem 3. Let the metric $g$ be non trapping and satisfying (1.8) and (1.9). Then for every $T>0$ and $\chi \in C_{0}^{\infty}\left(\mathbb{R}^{d}\right)$ there exists a constant $C$ such that if $u$ solves

$$
\begin{equation*}
i u_{t}+\Delta_{g} u=0,\left.\quad u\right|_{t=0}=u_{0} \in L^{2}\left(\mathbb{R}^{d}\right) \tag{6.1}
\end{equation*}
$$

then

$$
\|\chi u\|_{L_{T}^{p} L^{q}\left(\mathbb{R}^{d}\right)} \leq C\left\|u_{0}\right\|_{L^{2}\left(\mathbb{R}^{d}\right)}
$$

provided the couple $(p, q)$ satisfies the admissibility condition (1.4).
Proof. The non trapping assumption is only needed for the next proposition.
Proposition 6.1. The solution of (6.1) satisfies

$$
\|\chi u\|_{L_{T}^{2} H^{\frac{1}{2}}} \leq C\left\|u_{0}\right\|_{L^{2}\left(\mathbb{R}^{d}\right)}
$$

We refer to [15] for a proof of Proposition 6.1. Such estimates can be seen as a consequence of the smooth perturbation theory of Kato (see [24, Chapter XIII.7]). Let us also recall (see e.g. [9) that, via a quite general argument using the Fourier transform in $t$, one can freeze the time and Proposition 6.1 follows from (the time independent) estimates on the resolvent of $\Delta_{g}$, namely

$$
\begin{equation*}
\left\|\chi(P-\lambda \pm i 0)^{-1} \chi\right\|_{L^{2} \rightarrow L^{2}} \leq C\langle\lambda\rangle^{-1 / 2}, \quad \lambda \gg 1 . \tag{6.2}
\end{equation*}
$$

Recall that such resolvent estimates were extensively studied in particular in connection with the local energy decay for the wave equation $\left(\partial_{t}^{2}-\Delta_{g}\right) u=0$. For a proof of (6.2), we refer for instance to [27, 5].

Notice that Proposition 6.1 is the only place in the proof of Theorem 2 where we use the non trapping assumption.

Let us now come back to the proof of Theorem3. It will be a suitable combination of Proposition 6.1] and Proposition [5.4. Let us fix $\varphi \in C_{0}^{\infty}(\mathbb{R} \backslash\{0\})$. Set

$$
v(t):=\varphi\left(h^{2} P\right) \chi u(t) .
$$

Then $v$ solves

$$
\left(i \partial_{t}-P\right) v=-\varphi\left(h^{2} P\right)[P, \chi] u,\left.\quad v\right|_{t=0}=\varphi\left(h^{2} P\right) \chi u_{0} .
$$

We can now apply Proposition 5.4 to $v$ which gives

$$
\begin{aligned}
\|v\|_{L_{T}^{p} L^{q}} \leq C\left\|\varphi\left(h^{2} P\right) \chi u\right\|_{L_{T}^{\infty} L^{2}}+ & C h^{-\frac{1}{2}}\left\|\varphi\left(h^{2} P\right) \chi u\right\|_{L_{T}^{2} L^{2}} \\
& +C h^{\frac{1}{2}}\left\|\varphi\left(h^{2} P\right)[P, \chi] u\right\|_{L_{T}^{2} L^{2}}:=Q_{1}+Q_{2}+Q_{3},
\end{aligned}
$$

provided $(p, q)$ is satisfying (1.4). We now estimate separately $Q_{1}, Q_{2}$ and $Q_{3}$.
Bound for $Q_{1}$. Using the functional calculus and the Schur lemma, we get

$$
\begin{equation*}
\left\|\left[\varphi\left(h^{2} P\right), \chi\right] w\right\|_{L^{2}\left(\mathbb{R}^{d}\right)} \leq C h\|w\|_{L^{2}\left(\mathbb{R}^{d}\right)} . \tag{6.3}
\end{equation*}
$$

This implies that

$$
\begin{aligned}
Q_{1} & \leq C\left\|\chi \varphi\left(h^{2} P\right) u\right\|_{L_{T}^{\infty} L^{2}}+C h\|u\|_{L_{T}^{\infty} L^{2}} \\
& \leq C\left\|\varphi\left(h^{2} P\right) u\right\|_{L_{T}^{\infty} L^{2}}+C h\|u\|_{L_{T}^{\infty} L^{2}} \\
& =C\left\|\varphi\left(h^{2} P\right) u_{0}\right\|_{L^{2}}+C h\left\|u_{0}\right\|_{L^{2}},
\end{aligned}
$$

where in the last line we used that $\exp (-i t P)$ is an $L^{2}$ isometry.
Bound for $Q_{2}$. Let $\tilde{\varphi} \in C_{0}^{\infty}(\mathbb{R} \backslash\{0\})$ which is equal to one on the support of $\varphi$. Then using (6.3), we get

$$
Q_{2} \leq C h^{-\frac{1}{2}}\left\|\varphi\left(h^{2} P\right) \chi \tilde{\varphi}\left(h^{2} P\right) u\right\|_{L_{T}^{2} L^{2}}+C h^{\frac{1}{2}}\|u\|_{L_{T}^{\infty} L^{2}}:=Q_{21}+Q_{22}
$$

Since the support of $\varphi$ does not meet the origin, we can use (2.17) and thus

$$
Q_{21} \leq C\left\|\chi \tilde{\varphi}\left(h^{2} P\right) u\right\|_{L_{T}^{2} H^{\frac{1}{2}}\left(\mathbb{R}^{d}\right)} .
$$

An application of Proposition 6.1 gives

$$
Q_{21} \leq C\left\|\tilde{\varphi}\left(h^{2} P\right) u_{0}\right\|_{L^{2}}
$$

Clearly

$$
Q_{22}=C h^{\frac{1}{2}}\left\|u_{0}\right\|_{L^{2}} .
$$

Thus

$$
Q_{2} \leq C\left\|\tilde{\varphi}\left(h^{2} P\right) u_{0}\right\|_{L^{2}}+C h^{\frac{1}{2}}\left\|u_{0}\right\|_{L^{2}}
$$

Bound for $Q_{3}$. Let us take again $\tilde{\varphi} \in C_{0}^{\infty}(\mathbb{R} \backslash\{0\})$ which equals one on the support of $\varphi$. An application of Schur lemma yields the bound

$$
\left\|\left[\tilde{\varphi}\left(h^{2} P\right),[\chi, P]\right] w\right\|_{L^{2}\left(\mathbb{R}^{d}\right)} \leq C\|w\|_{L^{2}\left(\mathbb{R}^{d}\right)}
$$

Therefore

$$
Q_{3} \leq C h^{\frac{1}{2}}\left\|\varphi\left(h^{2} P\right)[\chi, P] \tilde{\varphi}\left(h^{2} P\right) u\right\|_{L_{T}^{2} L^{2}}+C h^{\frac{1}{2}}\left\|u_{0}\right\|_{L^{2}}:=Q_{31}+Q_{32}
$$

Let us next fix a $\tilde{\chi} \in C_{0}^{\infty}\left(\mathbb{R}^{d}\right)$ which is equal to one on the support of $\chi$. Then

$$
Q_{31}=C h^{\frac{1}{2}}\left\|\varphi\left(h^{2} P\right)[\chi, P] \tilde{\chi} \tilde{\varphi}\left(h^{2} P\right) u\right\|_{L_{T}^{2} L^{2}}
$$

Using (2.17) with $s=-1 / 2$, we obtain the bound

$$
Q_{31} \leq C\left\|[\chi, P] \tilde{\chi} \tilde{\varphi}\left(h^{2} P\right) u\right\|_{L_{T}^{2} H^{-\frac{1}{2}}} \leq C\left\|\tilde{\chi} \tilde{\varphi}\left(h^{2} P\right) u\right\|_{L_{T}^{2} H^{\frac{1}{2}}}
$$

where in the last line we used that $[\chi, P]$ is a first order differential operator with $C_{0}^{\infty}\left(\mathbb{R}^{d}\right)$ coefficients. A use of Proposition 6.1 yields

$$
Q_{31} \leq C\left\|\tilde{\varphi}\left(h^{2} P\right) u_{0}\right\|_{L^{2}}
$$

and therefore

$$
Q_{3} \leq C\left\|\tilde{\varphi}\left(h^{2} P\right) u_{0}\right\|_{L^{2}}+C h^{\frac{1}{2}}\left\|u_{0}\right\|_{L^{2}}
$$

Using the above bounds for $Q_{1}, Q_{2}, Q_{3}$, we arrive at the bound

$$
\begin{equation*}
\left\|\varphi\left(h^{2} P\right) \chi u\right\|_{L_{T}^{p} L^{q}} \leq C\left\|\tilde{\varphi}\left(h^{2} P\right) u_{0}\right\|_{L^{2}}+C h^{\frac{1}{2}}\left\|u_{0}\right\|_{L^{2}} \tag{6.4}
\end{equation*}
$$

provided $(p, q)$ is satisfying (1.4). With (6.4) in hand the proof of Theorem 3 is reduced to an application of the Littlewood-Paley square function theorem. Indeed, consider a Littlewood-Paley partition of the identity

$$
\begin{equation*}
\mathrm{Id}=\varphi_{1}(P)+\sum_{h^{-1}: \text { dyadic }} \varphi\left(h^{2} P\right) \tag{6.5}
\end{equation*}
$$

where $\varphi_{1} \in C_{0}^{\infty}(\mathbb{R}), \varphi \in C_{0}^{\infty}(\mathbb{R} \backslash\{0\})$. Using Proposition 2.10, we obtain that

$$
\|\chi u\|_{L_{T}^{p} L^{q}} \leq C\left\|u_{0}\right\|_{L^{2}}+C\left(\sum_{h^{-1}: \text { dyadic }}\left\|\varphi\left(h^{2} P\right) \chi u\right\|_{L_{T}^{p} L^{q}}^{2}\right)^{\frac{1}{2}}
$$

Coming to the crucial bound (6.4), using that $\tilde{\varphi} \in C_{0}^{\infty}(\mathbb{R} \backslash\{0\})$ for the first term in the right hand-side of (6.4), and, summing geometric series for the second, we arrive at the bound

$$
\|\chi u\|_{L_{T}^{p} L^{q}} \leq C\left\|u_{0}\right\|_{L^{2}}
$$

This completes the proof of Theorem 3.

Remark 6.2. Let us observe that the constants depending on the time intervals $[0, T]$ in all statements in this paper remain bounded as $T$ varies within a compact set. In other words the only possible blow up of these constants is as $T \rightarrow \infty$.

## 7. Non homogeneous estimates and nonlinear applications

The aim of this section is to give some applications of the estimates established in the previous sections to the Nonlinear Schrödinger equation

$$
\begin{equation*}
\left(i \partial_{t}+\Delta_{g}\right) u=F(u),\left.\quad u\right|_{t=0}=u_{0} \tag{7.1}
\end{equation*}
$$

where $u(t): \mathbb{R}^{d} \rightarrow \mathbb{C}$. The function $F(z), z \in \mathbb{C}$, is assumed to be smooth and vanishing at $z=0$. Moreover, we suppose that $F=\bar{\partial} V$ with a real valued "potential" $V$ satisfying the gauge invariance assumption

$$
V(\omega z)=V(z), \quad \forall \omega \in S^{1}, \forall z \in \mathbb{C} .
$$

In addition, we suppose that for some $\alpha>1$,

$$
\begin{equation*}
\left|\partial^{k_{1}} \bar{\partial}^{k_{2}} V(z)\right| \leq C_{k_{1}, k_{2}}\langle z\rangle^{1+\alpha-k_{1}-k_{2}} . \tag{7.2}
\end{equation*}
$$

The real number $\alpha$ involved in (7.2) corresponds to the "degree" of the nonlinear interaction. The problem (7.1) may, at last formally, be seen as a Hamiltonian PDE in an infinite dimensional phase space, with Hamiltonian

$$
\begin{equation*}
H(u, \bar{u})=\int_{\mathbb{R}^{d}}\left|\nabla_{g} u\right|^{2}+\int_{\mathbb{R}^{d}} V(u) \tag{7.3}
\end{equation*}
$$

and canonical coordinates $(u, \bar{u})$ (in (7.3) we integrate with respect to the volume element associated to $g$ ). Therefore the quantity (7.3) is formally conserved by the flow of (7.1). Another formally conserved quantity by the flow of (7.1) is the $L^{2}$ norm of $u$. In this section we make the defocusing assumption

$$
V(z) \geq 0
$$

on the potential $V$. Under this assumption the $H^{1}\left(\mathbb{R}^{d}\right)$ norm of the solutions of (7.1) may be expected to be controlled uniformly in time under the evolution of (7.1). Therefore the study of (7.1) in the space $H^{1}\left(\mathbb{R}^{d}\right)$ is of particular interest.

In the study of (7.1), $L^{p}$ analogues, $1<p<+\infty$, of (2.3) are useful. More precisely one has the bounds

$$
\begin{equation*}
C_{s, p}^{-1}\left\|(P+1)^{s / 2} u\right\|_{L^{p}\left(\mathbb{R}^{d}\right)} \leq\|u\|_{W^{s, p}\left(\mathbb{R}^{d}\right)} \leq C_{s, p}\left\|(P+1)^{s / 2} u\right\|_{L^{p}\left(\mathbb{R}^{d}\right)} . \tag{7.4}
\end{equation*}
$$

where $1<p<+\infty$. Estimates (7.4) follow from the $L^{p}, 1<p<+\infty$ boundedness of zero order pseudo differential operators (see e.g. [32]).

In this section, we give the rather standard consequences of Theorem $\square$ and Theorem 2 to the $H^{1}$ theory for (7.1). We start with a general result in dimension two.

Theorem 4. Let $\alpha>1$ be an arbitrary real number and let $g$ be a metric on $\mathbb{R}^{2}$ satisfying (2.1), (2.2). Then for every $u_{0} \in H^{1}\left(\mathbb{R}^{2}\right)$ there exists a unique global solution $u \in C\left(\mathbb{R} ; H^{1}\left(\mathbb{R}^{2}\right)\right.$ ) of (7.1).

Proof. Using Theorem 1, we obtain the estimate

$$
\begin{equation*}
\left\|\exp \left(i t \Delta_{g}\right) u_{0}\right\|_{L_{T}^{p} L^{q}} \leq C\left\|u_{0}\right\|_{H^{\frac{1}{p}}} \tag{7.5}
\end{equation*}
$$

provided $(p, q)$ is satisfying (1.4). With (7.5) in hand, the proof of Theorem(4) consists in word by word repetition of the analysis in [8, sec. 3.1 and 3.2].

Without the non trapping assumption, in dimension three, one can only get the following global existence result.

Theorem 5. Consider the cubic defocusing NLS

$$
\begin{equation*}
\left(i \partial_{t}+\Delta_{g}\right) u=|u|^{2} u,\left.\quad u\right|_{t=0}=u_{0} \tag{7.6}
\end{equation*}
$$

where $g$ is a metric on $\mathbb{R}^{3}$ satisfying (2.1), (2.2). Let $s>1$. Then for every $u_{0} \in H^{s}\left(\mathbb{R}^{3}\right)$ there exists a unique global solution $u \in C\left(\mathbb{R} ; H^{s}\left(\mathbb{R}^{3}\right)\right)$ of the Cauchy problem (7.6).

Proof. It relies on Lemmas 7.1 and 7.2 below. The first one is a local existence and uniqueness result which is proved in [8] (using the 3 dimensional analogues of (7.5) which follow from Theorem (1).

Lemma 7.1. Let $s>1$ and $p>2$ be such that $s>\frac{3}{2}-\frac{1}{p}$. Set $\sigma=s-1 / p$ and let $q>2$ be such that $(p, q)$ satisfy (1.4). Then, for all $u_{0} \in H^{s}$ and all $t_{0} \in \mathbb{R}$, there exists $\epsilon>0$ and a unique

$$
u \in C\left(\left[t_{0}-\epsilon, t_{0}+\epsilon\right], H^{s}\right) \cap L^{p}\left(\left[t_{0}-\epsilon, t_{0}+\epsilon\right], W^{\sigma, q}\right)
$$

such that

$$
\begin{equation*}
u(t)=e^{i\left(t-t_{0}\right) \Delta_{g}} u_{0}-i \int_{0}^{t} e^{i(t-\tau) \Delta_{g}}|u(\tau)|^{2} u(\tau) d \tau \tag{7.7}
\end{equation*}
$$

for all $t \in\left[t_{0}-\epsilon, t_{0}+\epsilon\right]$
The key step to get a global existence result is given by the following statement.
Lemma 7.2. With $s, p, \sigma, q$ as in Lemma 7.1, the following holds true: if there exists $T>0$ and

$$
u \in \bigcup_{0 \leq T^{\prime}<T} C\left(\left[0, T^{\prime}\right], H^{s}\right) \cap L^{p}\left(\left[0, T^{\prime}\right], W^{\sigma, q}\right)
$$

solution of (7.7) (with $t_{0}=0$ ) on $\left[0, T^{\prime}\right]$ for all $T^{\prime}$ such that $0 \leq T^{\prime}<T$, then there exists $C$ such that

$$
\sup _{\left[0, T^{\prime}\right]}\|u(t)\|_{H^{s}}+\int_{0}^{T^{\prime}}\|u(t)\|_{L^{\infty}}^{2} d t \leq C
$$

for all $0 \leq T^{\prime}<T$.
Before proving this lemma, let us show how we obtain Theorem 5. We consider $T:=\sup \left\{T^{\prime}>0 \mid u\right.$ solves (7.7) (with $\left.t_{0}=0\right)$ on $\left.\left[0, T^{\prime}\right]\right\}$ and we argue by contradiction, assuming that $T<\infty$. Indeed, using standard non linear estimates and Corollary 2.10 of [8], Lemma 7.2 shows that $u(T):=\lim _{t \rightarrow T} u(t)$ exists in $H^{s}$ and that $u \in L^{p}\left([0, T], W^{\sigma, q}\right)$. Thus, by Lemma 7.1, we can continue the solution on $[T, T+\epsilon]$, for some $\epsilon>0$, with initial data $u(T)$ which yields a contradiction. Of course, we argue similarly for negative times and this proves the existence and the uniqueness of a solution in $C\left([-T, T], H^{s}\right) \cap L^{p}\left([-T, T], W^{\sigma, q}\right)$ for all $T>0$, hence the existence of a solution in $C\left(\mathbb{R}, H^{s}\right)$. We omit the proof of the uniqueness in $C\left(\mathbb{R}, H^{s}\right)$ since it follows as in [8, 3.2] and since, here, the main point is the global existence. Let us finally notice that uniqueness of weak $H^{1}$ solutions can be established as in [8].

We now turn to the proof of Lemma 7.2 .
Proof. The first tool comes from the conservation laws which imply that there is a constant $C$ independent of $t$ (only depending on $\left\|u_{0}\right\|_{H^{1}}$ ) such that

$$
\begin{equation*}
\|u(t)\|_{H^{1}\left(\mathbb{R}^{3}\right)} \leq C \tag{7.8}
\end{equation*}
$$

as far as the solution exists, i.e. on $[0, T)$ here. The rigorous justification of these conservation laws requires a standard approximation argument (see e.g. [16]) The key quantity in this discussion is $\|u\|_{L_{T}^{2} L^{\infty}}$ (the number 2 is reflecting the cubic nature of the nonlinearity). Consider again the Littlewood-Paley partition of the identity (6.5). Then $v:=\varphi\left(h^{2} P\right) u$ solves the problem

$$
\left(i \partial_{t}-P\right) v=\varphi\left(h^{2} P\right)\left(|u|^{2} u\right)
$$

Using Proposition 2.9, Proposition 5.4 and the bound (7.8), we obtain that for all $\theta \leq \inf (1, T / 2)$ (see Remark 6.2)

$$
\begin{aligned}
& \left\|\varphi\left(h^{2} P\right) u\right\|_{L_{\theta}^{2} L^{\infty}} \leq C h^{-1 / 2}\left\|\varphi\left(h^{2} P\right) u\right\|_{L_{\theta}^{2} L^{6}} \\
& \quad \leq C h^{1 / 2}\|u\|_{L_{\theta}^{\infty} H^{1}}+C h^{-1}\left\|\varphi\left(h^{2} P\right) u\right\|_{L_{\theta}^{2} L^{2}}+C h^{-1 / 2}\left\|\varphi\left(h^{2} P\right)\left(|u|^{2} u\right)\right\|_{L_{\theta}^{2} L^{6 / 5}}
\end{aligned}
$$

Next, using Proposition [2.12, the Sobolev inequality and (7.8), we obtain that

$$
\begin{aligned}
h^{-1 / 2}\left\|\varphi\left(h^{2} P\right)\left(|u|^{2} u\right)\right\|_{L_{\theta}^{2} L^{6 / 5}} \leq C h^{1 / 2}\left\||u|^{2} u\right\|_{L_{\theta}^{2} W^{1,6 / 5}} & \leq \\
& \leq C h^{1 / 2}\|u\|_{L_{\theta}^{2} H^{1}} \leq C \sqrt{h \theta} .
\end{aligned}
$$

In summary,

$$
\left\|\varphi\left(h^{2} P\right) u\right\|_{L_{\theta}^{2} L^{\infty}} \leq C h^{1 / 2}+C \sqrt{h \theta}+C h^{-1}\left\|\varphi\left(h^{2} P\right) u\right\|_{L_{\theta}^{2} L^{2}} .
$$

Next, for $N \in \mathbb{N}$, using the Cauchy-Schwarz inequality we get

$$
\sum_{h^{-1} \leq N} h^{-1}\left\|\varphi\left(h^{2} P\right) u\right\|_{L_{\theta}^{2} L^{2}} \leq C \sqrt{\log N}\|u\|_{L_{\theta}^{2} H^{1}} \leq C \sqrt{\theta \log N} .
$$

On the other hand, since $s>1$, we estimate the high frequencies as follows

$$
\sum_{h^{-1} \geq N} h^{-1}\left\|\varphi\left(h^{2} P\right) u\right\|_{L_{\theta}^{2} L^{2}} \leq C N^{-(s-1)}\|u\|_{L_{\theta}^{2} H^{s}}
$$

By taking $N \approx\left(2+\|u\|_{L_{\theta}^{\infty} H^{s}}\right)^{\frac{1}{s-1}}$, we deduce that

$$
\begin{equation*}
\|u\|_{L_{\theta}^{2} L^{\infty}} \leq C+C\left[\theta\left(1+\log \left(2+\|u\|_{L_{\theta}^{\infty} H^{s}}\right)\right)\right]^{\frac{1}{2}} . \tag{7.9}
\end{equation*}
$$

Coming back to the integral equation (7.7) and using the Gronwall lemma, we obtain that

$$
\begin{equation*}
\|u\|_{L_{\theta}^{\infty} H^{s}} \leq\left\|u_{0}\right\|_{H^{s}} e^{C\|u\|_{L_{\theta}^{2} L^{\infty}}^{2}} \leq C\left\|u_{0}\right\|_{H^{s}}\left[2+\|u\|_{L_{\theta}^{\infty} H^{s}}\right]^{\Lambda \theta} \tag{7.10}
\end{equation*}
$$

where $\Lambda$ is a real number depending only on the a priori bound (7.8). Therefore, if we take $\theta$ such that $\Lambda \theta \leq 1 / 2$, we obtain that

$$
\begin{equation*}
\|u\|_{L_{\theta}^{\infty} H^{s}} \leq C\left(\left\|u_{0}\right\|_{H^{s}}+\left\|u_{0}\right\|_{H^{s}}^{2}\right) . \tag{7.11}
\end{equation*}
$$

Iterating finitely many times $(\approx T / \theta$ times) (7.11) and (7.9) yields the result.
Remark 7.3. Once we know that we have a global solution, we can control the growth of $\|u(t)\|_{H^{s}}$ as $t \rightarrow \infty$ since, by iterations of (7.11), one can easily check that

$$
\|u(t)\|_{H^{s}} \leq C \exp (C \exp (C|t|)), \quad t \in \mathbb{R}
$$

Notice that the results of Theorem 4 and Theorem 5 hold without the long range assumption (1.9). Moreover, we do not suppose that the metric is non trapping. If we assume these two conditions, we can improve Theorem to nonlinearities of higher degree and even get global existence results in dimensions $d \geq 4$. For that purpose, we need the following non homogeneous Strichartz estimate.

Theorem 6. Suppose that $g$ is a non trapping metric on $\mathbb{R}^{d}$ satisfying (1.8), (1.9). Then for every $T>0$ there exists $C>0$ such that that if $u$ solves

$$
\begin{equation*}
\left(i \partial_{t}+\Delta_{g}\right) u=f,\left.\quad u\right|_{t=0}=0 \tag{7.12}
\end{equation*}
$$

then

$$
\begin{equation*}
\|u\|_{L_{T}^{p} L^{q}\left(\mathbb{R}^{d}\right)} \leq C\|f\|_{L_{T}^{p_{1}} L^{q_{1}}\left(\mathbb{R}^{d}\right)} \tag{7.13}
\end{equation*}
$$

provided $(p, q),\left(\frac{p_{1}}{p_{1}-1}, \frac{q_{1}}{q_{1}-1}\right)$ are satisfying 1.4) and $p \neq 2, p_{1} \neq 2$.
Proof. The proof is a consequence of the following Christ-Kiselev lemma [13].
Lemma 7.4. Let $T>0$ be a real number. Let $B_{1}$ and $B_{2}$ be two Banach spaces. Let $K(t, s)$ be a locally integrable kernel with values in the bounded operators from $B_{1}$ to $B_{2}$. Let $1 \leq p<q \leq \infty$. For $t \in[0, T]$, we set

$$
A f(t)=\int_{0}^{T} K(t, s) f(s) d s
$$

Assume that

$$
\|A f\|_{L^{q}\left([0, T] ; B_{1}\right)} \leq C\|f\|_{L^{p}\left([0, T] ; B_{2}\right)}
$$

Define the operator $\tilde{A}$ as

$$
\tilde{A} f(t)=\int_{0}^{t} K(t, s) f(s) d s, \quad t \in[0, T]
$$

Then there exists $\tilde{C}>0$ such that

$$
\|\tilde{A} f\|_{L^{q}\left([0, T] ; B_{1}\right)} \leq \tilde{C}\|f\|_{L^{p}\left([0, T] ; B_{2}\right)}
$$

We refer to [31] for a proof of Lemma 7.4, in the form stated here. Let us now return to the proof of Theorem 6, The solution of (7.12) is given by

$$
u(t)=\int_{0}^{t} \exp \left(i(t-\tau) \Delta_{g}\right) f(\tau) d \tau
$$

Consider $v(t)$ defined by

$$
v(t)=\int_{0}^{T} \exp \left(i(t-\tau) \Delta_{g}\right) f(\tau) d \tau
$$

Then

$$
v(t)=\exp \left(i t \Delta_{g}\right) \int_{0}^{T} \exp \left(-i \tau \Delta_{g}\right) f(\tau) d \tau
$$

and by invoking (1.12), we get the bound

$$
\|v\|_{L_{T}^{p} L^{q}} \leq C\left\|\int_{0}^{T} \exp \left(-i \tau \Delta_{g}\right) f(\tau) d \tau\right\|_{L^{2}}
$$

The dual of (1.12) yields

$$
\begin{equation*}
\left\|\int_{0}^{T} \exp \left(-i \tau \Delta_{g}\right) f(\tau) d \tau\right\|_{L^{2}} \leq C\|f\|_{L_{T}^{\frac{p}{p-1}} L^{\frac{q}{q-1}}} . \tag{7.14}
\end{equation*}
$$

Therefore

$$
\|v\|_{L_{T}^{p} L^{q}} \leq C\|f\|_{L_{T}^{\frac{p}{p-1}} L^{\frac{q}{q-1}}} .
$$

An application of Lemma 7.4 gives

$$
\|v\|_{L_{T}^{p} L^{q}} \leq C\|f\|_{L_{T}^{\frac{p}{p-1}} \frac{q}{L^{q-T}}}, \quad p>2 .
$$

This proves (7.13) in the case $\frac{1}{p}+\frac{1}{p_{1}}=\frac{1}{q}+\frac{1}{q_{1}}=1$. Let us now prove (7.13) when $(p, q)$ and $\left(p_{1}, q_{1}\right)$ are decoupled. Since $\exp \left(i t \Delta_{g}\right)$ is an isometry on $L^{2}$, using (7.14) and Remark 6.2, we infer that (7.13) is valid for $(p, q)=(\infty, 2)$ and $\left(p_{1}, q_{1}\right)$ an arbitrary pair satisfying (1.4). Using the homogeneous estimate (1.12) and the Minkowski inequality, we obtain that (7.13) is valid for $(p, q)$ an arbitrary pair satisfying (1.4) and $\left(p_{1}, q_{1}\right)=(\infty, 2)$. Let us now observe that all other cases for $(p, q)$ and $\left(p_{1}, q_{1}\right)$ in (7.13) follow from the considered three particular cases by interpolation. This completes the proof of Theorem 6.

It is now a standard and straightforward consequence of Theorem 6. Theorem 2 and (7.4) (see [17, 22, 11, 16]) that one has the following global well-posedness result for (7.1).
Theorem 7. Let $d \geq 3$. Suppose that $\alpha \in] 1,1+\frac{4}{d-2}[$. Let $g$ be a non trapping metric on $\mathbb{R}^{d}$ satisfying (1.8), (1.9). Then for every $u_{0} \in H^{1}\left(\mathbb{R}^{d}\right)$ there exists a unique global solution $u \in C\left(\mathbb{R} ; H^{1}\left(\mathbb{R}^{d}\right)\right)$ of (7.1).
Remark 7.5. Recall that the endpoint Strichartz estimates are not needed for the standard $H^{1}$ theory of (7.1).
Remark 7.6. Let us emphasize the importance of the non homogeneous Strichartz estimates in the proof of Theorem 7, The lack of such estimates under the very weak hypotheses (2.1), (2.2) makes the study of (7.1) in $H^{1}$ in this case (or in the case of a compact manifold) more difficult and so far restricted only to small dimensions.

## References

[1] S. Alinhac, P. Gérard. Opérateurs pseudo-différentiels et théorème de Nash-Moser. Savoir actuels, Editions du CNRS, 1991.
[2] J.M. Bouclet. Distributions spectrales pour des opérateurs perturbés. PhD Thesis, Nantes University, 2000. http://tel.ccsd.cnrs.fr/document/archives0/00/00/40/25/
[3] J.M. Bouclet. Spectral distributions for long range perturbations. J. Funct. Anal. 212: 431-471, 2004.
[4] J. Bourgain. Fourier transform restriction phenomena for certain lattice subsets and application to nonlinear evolution equations I. Schrödinger equations. Geom. and Funct. Anal., 3: 107-156, 1993.
[5] N. Burq. Semi-classical estimates for the resolvent in nontrapping geometries. International Math. Res. Notices, 221-241, 2002.
[6] N. Burq. Lower bounds for shape resonances widths of long range Schrödinger operators. Amer. J. of Math., 124: 677-735, 2002.
[7] N. Burq. Estimation de Strichartz pour des perturbations à longue portée de l'opérateur de Schrödinger. Seminaire EDP, Ecole polytechnique, 2002.
[8] N. Burq, P. Gérard, N. Tzvetkov. Strichartz inequalities and the non linear Schrödinger equation on compact manifolds. Amer. J. of Math., 126: 569-605, 2004.
[9] N. Burq, P. Gérard, N. Tzvetkov. On non linear Schrödinger equations in exterior domains. Ann. I.H. Poincaré-Analyse non linéaire 21: 295-318, 2004.
[10] F. Cardoso, G. Vodev. Uniform estimates of the resolvent of the Laplace-Beltrami operator on infinite volume Riemannian manifolds. II. Ann. Henri Poincaré 3: 673-691, 2002.
[11] Th. Cazenave. Semilinear Schrödinger equations. Courant lecture notes 10, AMS, 2003.
[12] J.Y. Chemin. Fluides parfait incompressibles. Astérisque 230: 1995.
[13] M. Christ, A. Kiselev. Maximal functions associated to filtrations. J. Funct. Anal. 179: 409-425, 2001.
[14] M. Dimassi, J. Sjöstrand. Spectral asymptotics in the semi-classical limit. London Mathematical Society Lecture Note Series, 268. Cambridge University Press, 1999.
[15] S. Doi. Smoothing effects of Schrödinger evolution groups on Riemannian manifolds. Duke Math. J. 82: 679-706, 1996.
[16] J. Ginibre. Introduction aux équations de Schrödinger non linéaires. Cours de DEA 1994-1995, Edition Paris onze, 1998.
[17] J. Ginibre, G. Velo. On a class of non linear Schrödinger equations. J. Funct. Anal. 32: 1-71, 1979.
[18] J. Ginibre, G. Velo. The global Cauchy problem for the non linear Schrödinger equation. Ann. I.H. Poincaré-Analyse non linéaire, 2: 309-327, 1985.
[19] A. Hassell, T. Tao, J. Wunsch. Sharp Strichartz estimates on non-trapping asymptotically conic manifolds. Amer. J. of Math., to appear.
[20] B. Helffer, J. Sjöstrand. Equation de Schrödinger avec champs magnétiques et équation de Harper. Lecture notes in Phys. 345, Springer-Verlag, Berlin, 1989.
[21] H. Isozaki, H. Kitada. Modified wave operators with time independent modifiers. J. Fac. Sci. University of Tokyo, Section I A 32: 77-104, 1985.
[22] T. Kato. On non linear Schrödinger equation. Ann. I.H. Poincaré, Phys. Théor., 46: 113-129, 1987.
[23] M. Keel, T. Tao. Endpoint Strichartz estimates. Amer. J. of Math., 120: 955-980, 1998.
[24] M. Reed, B. Simon. Methods of modern mathematical physics, vol. 4. Academic Press, 1978.
[25] L. Robbiano, C. Zuily. Strichartz estimates for Schrödinger equations with variable coefficients. Mém. Soc. Math. Fr. (N.S.), No. 101-102, +208 pp, 2005.
[26] D. Robert. Autour de l'approximation semi-classique. Progr. Math. 68, Birkhäuser, Basel, 1987.
[27] D. Robert. Asymptotique de la phase de diffusion à haute énergie pour des perturbation du second ordre du Laplacien. Ann. Scient. ENS 25: 107-134, 1992.
[28] D. Robert. Relative time delay for perturbations of elliptic operators and semi-classical asymptotics. J. Funct. Anal. 126: 36-82, 1994.
[29] I. Rodnianski, T. Tao. Long-time decay estimates for the Schrödinger equation on manifolds. Preprint 2004.
[30] R.T. Seeley, Complex powers of an elliptic operator. Singular integrals (Proc. Sympos. Pure Math. Chicago, III 1966) A.M.S. R.I., 288-307, 1967.
[31] H. Smith, C. Sogge. Global Strichartz estimates for nontrapping perturbations of the Laplacian. Comm. PDE 25: 2171-2183, 2000.
[32] C. Sogge. Fourier integrals in classical analysis. Cambridge tracts in Mathematics, 1993.
[33] G. Staffilani, D. Tataru. Strichartz estimates for a Schrödinger operator with non smooth coefficients. Comm. PDE 27: 1337-1372, 2002.

Département de Mathématiques, Université Lille I, 59655 Villeneuve d’Ascq Cedex, France

E-mail address: jean-marc.bouclet@math.univ-lille1.fr
Département de Mathématiques, Université Lille I, 59655 Villeneuve d’AscQ Cedex, France

E-mail address: nikolay.tzvetkov@math.univ-lille1.fr

