On the scattering theory of asymptotically flat manifolds and Strichartz inequalities

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Conférence en l'honneur de Vladimir Georgescu

Introduction

Purpose of the talk

- ► Take the question of Strichartz inequalities (for the Schrödinger equation) on asymptotically flat manifolds as a case study to review some related scattering estimates (resolvent estimates, time decay, smoothing estimates), either for comparison or because they are crucial inputs in the proofs of Strichartz inequalities
- ► Present some recent results (joint with H. Mizutani) on Strichartz inequalities on asymptotically flat manifolds

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provided (p,q) is admissible (scaling condition)

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$$||e^{it\Delta}u_0||_{L^2(K)}\lesssim_K ||e^{it\Delta}u_0||_{L^q(\mathbb{R}^n)}, \qquad K\in\mathbb{R}^n.$$

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By scaling

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with $\nu>1/2$. By tracking the dependence on λ , one may obtain the non spectrally localized estimate ($n\geq 3$)

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Intuition. More on the next slides.

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$$R_0(\lambda + i0) - R_0(\lambda - i0)$$

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which is the $\frac{1}{2}$ -smoothing effect for the Schrödinger equation. Note that even locally in time (i.e. with $\mathbb R$ replaced by [-T,T]) this is non trivial.

Intuition. More on the next slides. Technically, they follow from resolvent estimates via a Parseval argument, using that $e^{it\Delta}$ is the Fourier transform $(\lambda \to t)$ of the spectral measure

$$R_0(\lambda + i0) - R_0(\lambda - i0).$$

Rem. This correspondence $\lambda \to t$ also allows to convert resolvent estimates into time decay/propagation estimates (smoothness of $R_0(\lambda \pm i0) \leftrightarrow$ decay of e^{itP})



Strichartz inequalities

Strichartz inequalities

Consider the L^2 normalized semiclassical wave packet

$$G_{z,\zeta,h}(x) = (\pi h)^{-\frac{n}{4}} \exp\left(\frac{i}{h}\zeta \cdot (x-z) - \frac{|x-z|^2}{2h}\right).$$

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with $\langle \tau \rangle = (1 + \tau^2)^{\frac{1}{2}}$. This implies easily

$$\left|\left|e^{i\frac{t}{2}\Delta}G_{z,\zeta,h}\right|\right|_{L^{q}}=\left(2/q\right)^{\frac{n}{2q}}\left(\frac{1}{\pi h\langle t/h\rangle^{2}}\right)^{\frac{n}{2}\left(\frac{1}{2}-\frac{1}{q}\right)}$$

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Remark. The translation by $(t/h)\zeta$ is **not** used. Only the spreading/dilation factor $\langle t/h \rangle$ plays a role.

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In particular, for $q = 2^* = 2n/(n-2)$,

$$\int_{-T}^{T} \left| \left| e^{i\frac{t}{2}\Delta} G_{z,\zeta,h} \right| \right|_{L^{2*}}^{2} dt$$

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for all $h \in (0,1]$

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for all $h \in (0,1]$ and $z \in \mathbb{R}^n$.

Strichartz inequalities vs smoothing effect for a wave packet Smoothing effect (local in time)

$$|\langle D \rangle^s e^{i\frac{t}{2}\Delta} G_{z,\zeta,h}(x)|$$

Smoothing effect (local in time)

$$\left|\langle D\rangle^s e^{i\frac{t}{2}\Delta} G_{z,\zeta,h}(x)\right| \sim \left|\langle \zeta/h\rangle^s \frac{\pi^{-\frac{n}{4}}}{\left(h\langle t/h\rangle^2\right)^{\frac{n}{4}}} \exp\left(-\frac{\left|x-z-(t/h)\zeta\right|^2}{2h\langle t/h\rangle^2}\right) \qquad h \to 0,$$

Smoothing effect (local in time)

$$\begin{split} \left| \langle D \rangle^s e^{i \frac{t}{2} \Delta} G_{z,\zeta,h}(x) \right| &\sim & \langle \zeta/h \rangle^s \frac{\pi^{-\frac{n}{4}}}{\left(h \langle t/h \rangle^2 \right)^{\frac{n}{4}}} \exp \left(-\frac{\left| x - z - (t/h) \zeta \right|^2}{2h \langle t/h \rangle^2} \right) & h \to 0, \\ &= & \langle \zeta/h \rangle^s G_{z,\zeta,h}^t(x). \end{split}$$

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We assume that $\zeta \neq 0$,

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We assume that $\zeta \neq 0$, say $|\zeta| = 1$

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If we further integrate in time on $[-T, T]_t$,

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$$c_n {\color{red}h} \langle \zeta/h \rangle^{2s} \int_{-T/h}^{T/h} \langle \tau \rangle^{-n} \int \langle h^{\frac{1}{2}} y + z + \tau \zeta \rangle^{-2\nu} \exp \left(-\frac{y^2}{\langle \tau \rangle^2} \right) dy d\tau$$

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which is bounded by

$$c_n h \langle 1/h \rangle^{2s} \int_{-T/h}^{T/h} \int \langle h^{\frac{1}{2}} Y_1 \langle \tau \rangle + z_1 + \tau \rangle^{-2\nu} \exp\left(-Y^2\right) dY d\tau$$

Smoothing effect (local in time)

$$\begin{aligned} \left| \langle D \rangle^s e^{i\frac{t}{2}\Delta} G_{z,\zeta,h}(x) \right| &\sim & \langle \zeta/h \rangle^s \frac{\pi^{-\frac{n}{4}}}{\left(h \langle t/h \rangle^2 \right)^{\frac{n}{4}}} \exp \left(-\frac{\left| x - z - (t/h)\zeta \right|^2}{2h \langle t/h \rangle^2} \right) & h \to 0, \\ &= & \langle \zeta/h \rangle^s G_{z,\zeta,h}^t(x). \end{aligned}$$

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$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^sG^t_{z,\zeta,h}\right|\right|^2_{L^2_x}=c_n\langle\zeta/h\rangle^{2s}\langle t/h\rangle^{-n}\int\langle h^{\frac{1}{2}}y+z+t\zeta/h\rangle^{-2\nu}\exp\left(-\frac{y^2}{\langle t/h\rangle^2}\right)dy$$

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Remark. Up to the term $Y_1\langle \tau \rangle$, there is no more contribution of the spreading $\langle \tau \rangle$.

Smoothing effect (local in time)

$$\begin{split} \left| \langle D \rangle^s e^{i\frac{t}{2}\Delta} G_{z,\zeta,h}(x) \right| &\sim & \langle \zeta/h \rangle^s \frac{\pi^{-\frac{n}{4}}}{\left(h \langle t/h \rangle^2 \right)^{\frac{n}{4}}} \exp \left(-\frac{\left| x - z - (t/h)\zeta \right|^2}{2h \langle t/h \rangle^2} \right) \qquad h \to 0, \\ &= & \langle \zeta/h \rangle^s G_{z,\zeta,h}^t(x). \end{split}$$

We assume that $\zeta \neq 0$, say $|\zeta|=1$ and then, by possibly rotating the axis, that $\zeta=(1,0,\ldots,0).$ Then

$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^sG^t_{z,\zeta,h}\right|\right|^2_{L^2_x}=c_n\langle\zeta/h\rangle^{2s}\langle t/h\rangle^{-n}\int\langle h^{\frac{1}{2}}y+z+t\zeta/h\rangle^{-2\nu}\exp\left(-\frac{y^2}{\langle t/h\rangle^2}\right)\mathrm{d}y$$

If we further integrate in time on $[-T, T]_t$,

$$c_n {\color{red}h} \langle \zeta/h \rangle^{2s} \int_{-T/h}^{T/h} \langle \tau \rangle^{-n} \int \langle h^{\frac{1}{2}} y + z + \tau \zeta \rangle^{-2\nu} \exp \left(-\frac{y^2}{\langle \tau \rangle^2} \right) dy d\tau$$

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Remark. Up to the term $Y_1\langle \tau \rangle$, there is no more contribution of the spreading $\langle \tau \rangle$. Here, the main role will be played the translation by $(t/h)\zeta = \tau \zeta$.

Recall we are estimating

$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^sG^t_{z,\zeta,h}\right|\right|^2_{L^2_{t,x}}\lesssim h^{1-2s}\int_{-T/h}^{T/h}\int\langle h^{\frac{1}{2}}Y_1\langle\tau\rangle+z_1+\tau\rangle^{-2\nu}\exp\left(-Y^2\right)dYd\tau.$$

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▶ In the region $|h^{1/2}Y_1| \le \epsilon$ ($\epsilon \ll 1$ fixed),

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In the region $|h^{1/2}Y_1| \le \epsilon$ ($\epsilon \ll 1$ fixed), we integrate in time by using the variable

$$ilde{ au} = au + h^{rac{1}{2}} Y_1 \langle au
angle \qquad ext{(Jacobian} = 1 + O(\epsilon))$$

Recall we are estimating

$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^sG^t_{z,\zeta,h}\right|\right|^2_{L^2_{t,x}}\lesssim h^{1-2s}\int_{-T/h}^{T/h}\int\langle h^{\frac{1}{2}}Y_1\langle\tau\rangle+z_1+\tau\rangle^{-2\nu}\exp\left(-Y^2\right)dYd\tau.$$

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Scattering inequalities turn out to play a crucial role in this problem.

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▶ More general model: asymptotically conical manifolds.

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Good models of scattering theory

More general model: asymptotically conical manifolds. In polar coordinates, $\mathbb{R}^n \setminus 0$ equipped with the Euclidean metric is isometric to

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Question: behavior of $R(\lambda \pm i0)$ and (2) as $\lambda \to \infty$ (high energy) and $\lambda \to 0$ (low energy) ?

Scattering estimates on asymptotically flat manifolds High energy estimates $(\lambda \to +\infty)$

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▶ Partial converse for trapping manifolds: if there are trapped geodesics

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[Bony-Burg-Ramond]



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Robust estimates for powers

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consequence on time decay

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[Burq-Guillarmou-Hassell]

Intuition (non trapping case):

▶ Inside a compact set *K*, combine

$$||\mathbf{1}_K e^{i \cdot P} u_0||_{L^2([-T,T],L^{2^*})} \lesssim_T ||u_0||_{H^{1/2}(M)} \text{ and } ||\mathbf{1}_K e^{i \cdot P} v_0||_{L^2([-T,T],H^{1/2})} \lesssim_T ||v_0||_{L^2([-T,T],H^{1/2})} ||v_0||_{L^2([-T,T],H^{1/2})} \lesssim_T ||v_0||_{L^2([-T,T],H^{1/2})} ||v_$$

▶ Outside a compact set: use that the geometry is close to a nice model (...)

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Results (joint with H. Mizutani)

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Theorem 3 (global space-time estimates without loss of derivatives) If $n \ge 3$ and the trapping is hyperbolic with negative pressure, then for (p,q) admissible

$$||e^{-i\cdot P}u_0||_{L^p(\mathbb{R};L^q(M))}\leq C||u_0||_{L^2(M)}.$$



Let $f_0 \in C_0^{\infty}(\mathbb{R})$ be such that $f_0 = 1$ near 0.

Theorem I (low frequency) If $n \ge 3$ and (p, q) is admissible

$$||f_0(P)e^{-i\cdot P}u_0||_{L^p(\mathbb{R};L^q(M))}\leq C||u_0||_{L^2(M)}.$$

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Theorem 4 (nonlinear scattering) Under the assumptions of Theorem 3, the L^2 critical equation

$$i\partial_t u - Pu = \sigma |u|^{\frac{4}{n}}u, \qquad u_{|t=0} = u_0, \qquad \sigma = \pm 1,$$

with $||u_0||_{L^2}\ll 1$, has a unique solution in (a subspace of) $C(\mathbb{R},L^2)\cap L^{2+\frac{4}{n}}(\mathbb{R}\times M)$ and

$$||u(t)-e^{-itP}u_{\pm}||_{L^2(M)}\to 0, \qquad t\to \pm\infty.$$

Low frequency localization in the uncertainty region:

Low frequency localization in the uncertainty region:in the regime $\lambda=\epsilon^2\to 0,$ how to prove

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Rem. For the localization, $(1 - \chi(\epsilon r))f(P/\epsilon^2)$, one has " $|\xi| \sim \epsilon$ " and " $|x| \gtrsim \epsilon^{-1}$ " \Rightarrow no problem of uncertainty principle to use microlocal techniques

Rest of the proof

At infinity: split $f(P/\lambda)e^{itP}$ into sums of

$$T_{\lambda}(t) = L_{\lambda} f(P/\lambda) e^{itP}$$

with suitable localization operators L_{λ} , and show

$$||T_{\lambda}(t)||_{L^2 \to L^2} \lesssim 1, \qquad ||T_{\lambda}(t)T_{\lambda}(s)||_{L^1 \to L^{\infty}} \lesssim |t-s|^{-\frac{n}{2}}$$

by writing

$$T_{\lambda}(t)T_{\lambda}(s)=\mathsf{approximation}+\mathsf{remainder}$$

- ▶ the "approximation" is explicit enough operator to bound sharply its integral kernel by $|t-s|^{-\frac{n}{2}}$ (dispersion bound)
- ▶ the remainder is a remainder term in a Duhamel formula in which we combine L^2 time decay/propagation estimates (for the time decay) and Sobolev estimates (to replace $L^2 \to L^2$ by $L^1 \to L^\infty$) to derive dispersion bounds.